

# Building Blocks of a PT-Symmetric Interferometer

Kagan Yanik

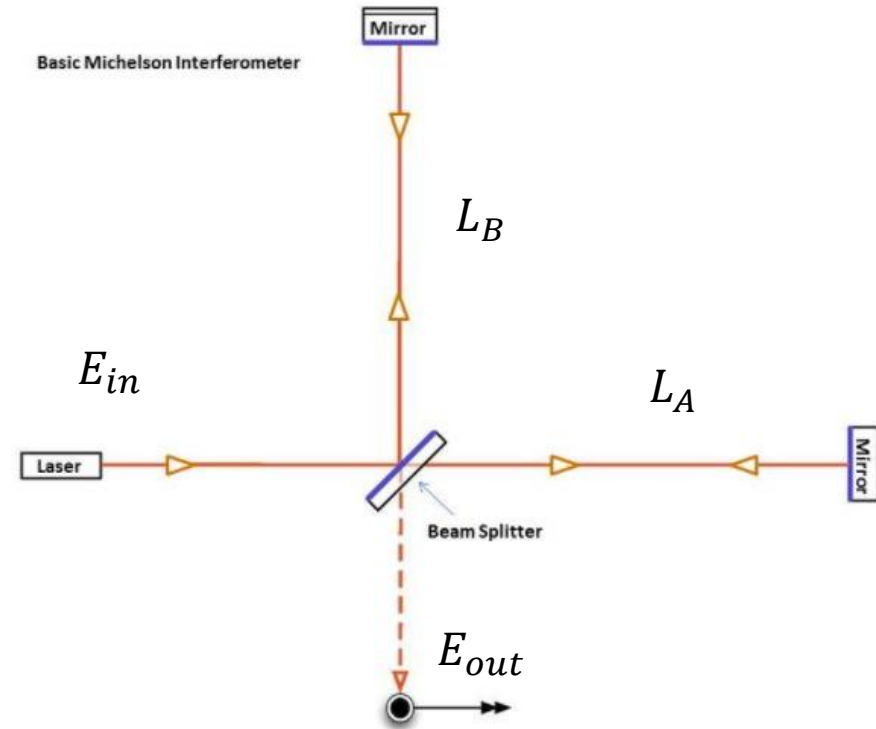
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# Conventional LIGO

- Michelson interferometer
- Differential path length change due to GW
- $\varphi_A = 2kL_A$
- $\varphi_B = 2kL_B$
- $P_{out} = (1 - \cos(\varphi_A - \varphi_B)) \frac{P_{in}}{2}$

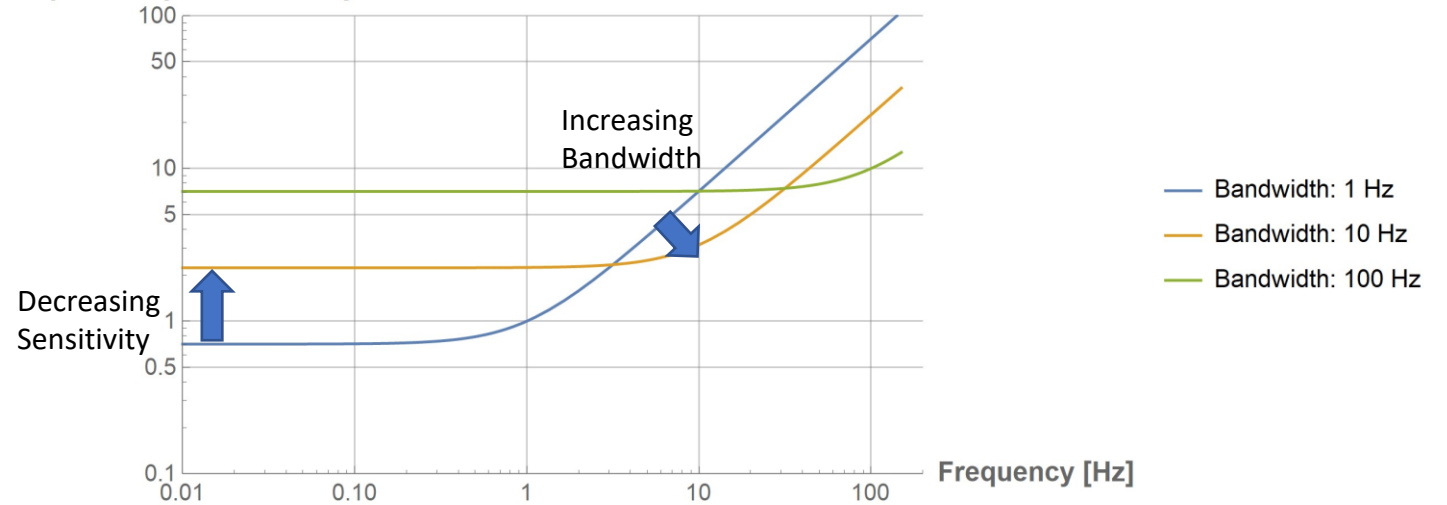


# Bandwidth vs Sensitivity

- Energetic Quantum Limit (EQL)
- We want to increase the bandwidth without sacrificing the sensitivity
- Overcoming the trade-off: Higher sensitivity in the entire frequency range

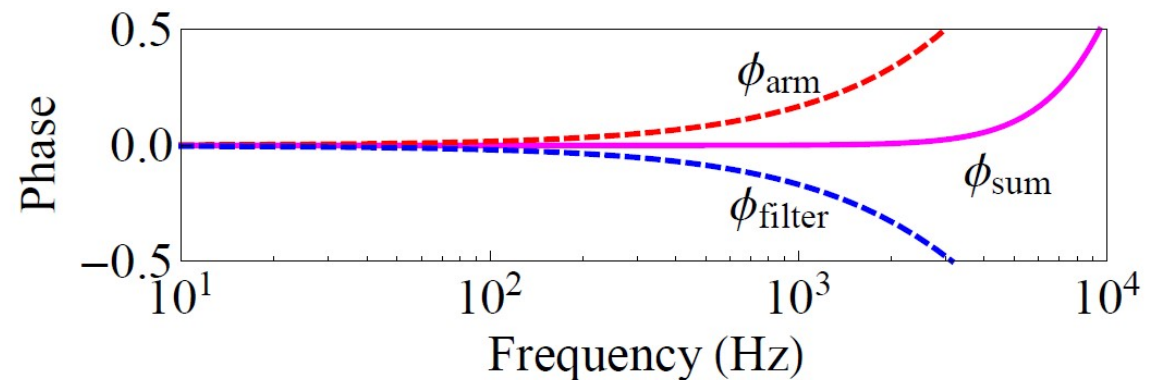
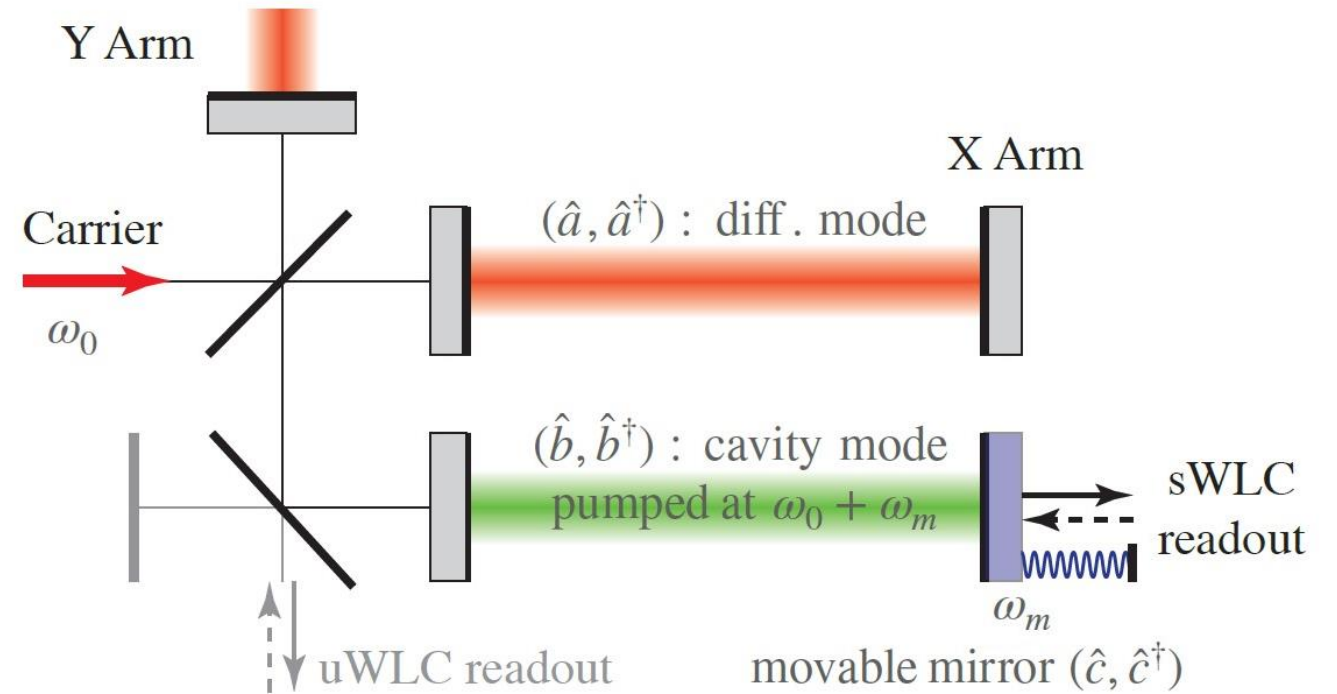
$$\int_0^{+\infty} d\Omega / (2\pi) S_h^{-1}(\Omega) \leq \Delta\mathcal{E}^2 / (4\hbar)^2$$

Amplitude Spectral Density



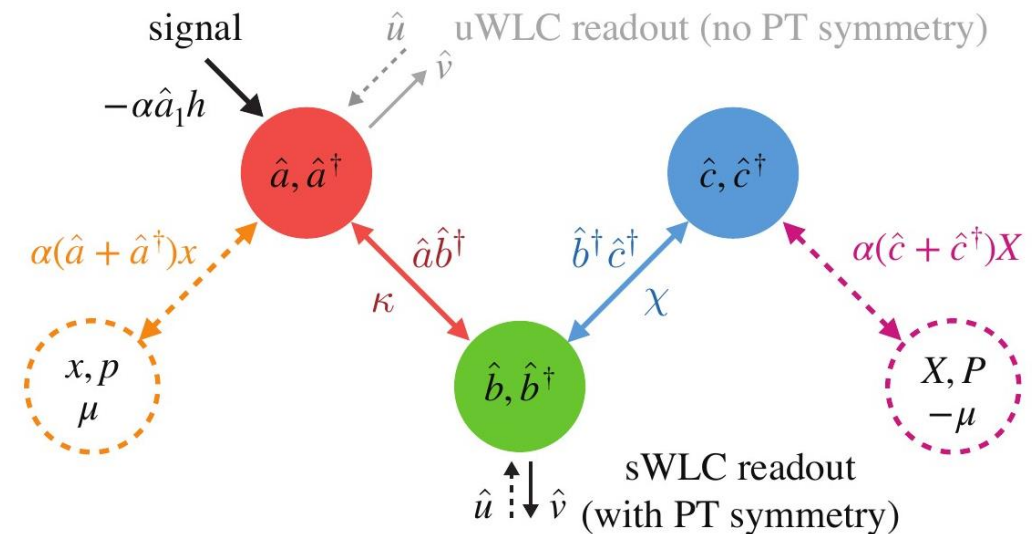
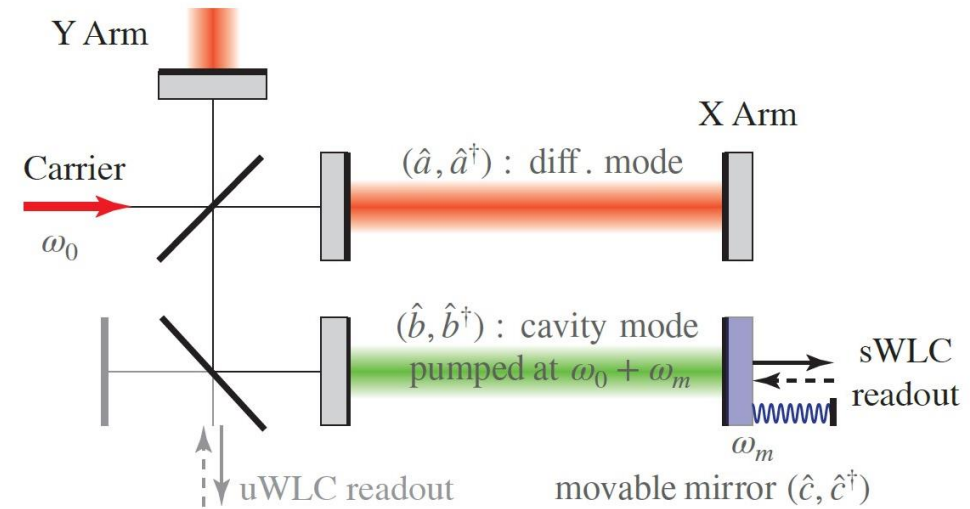
# Setup of the System

- Arm cavity -  $\hat{a}$  (optical mode)
- Filter cavity -  $\hat{b}$  (optical mode)
- Mechanical mode -  $\hat{c}$  (attached to filter cavity)



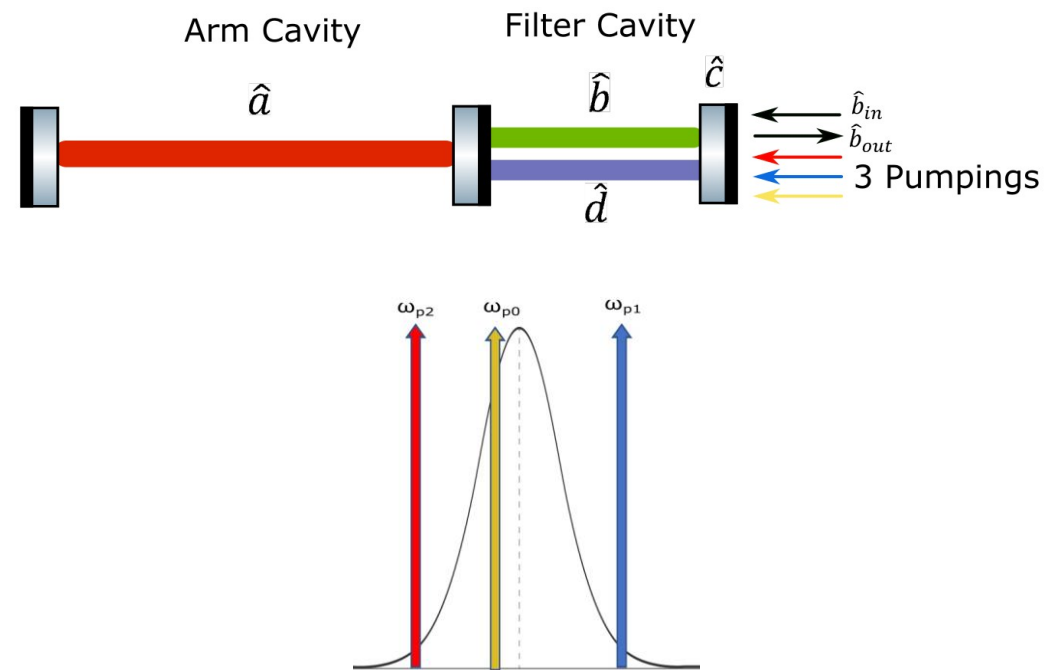
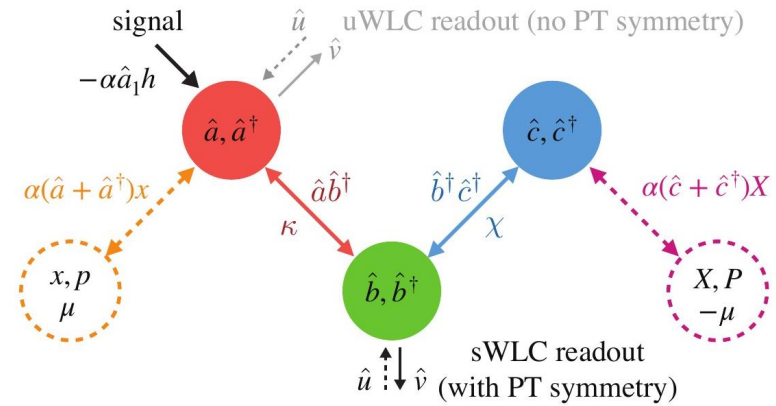
# PT-symmetric Interferometer

- White Light Cavity can improve the sensitivity while sacrificing less of the bandwidth



# Negative Mass

- Optical mode -  $\hat{d}$
- $\hat{V}_{aux} = -\hbar\alpha_d\hat{c}_1\hat{d}_1 - \hbar\omega_d\hat{d}^\dagger\hat{d}$



# Mathematica Simulations

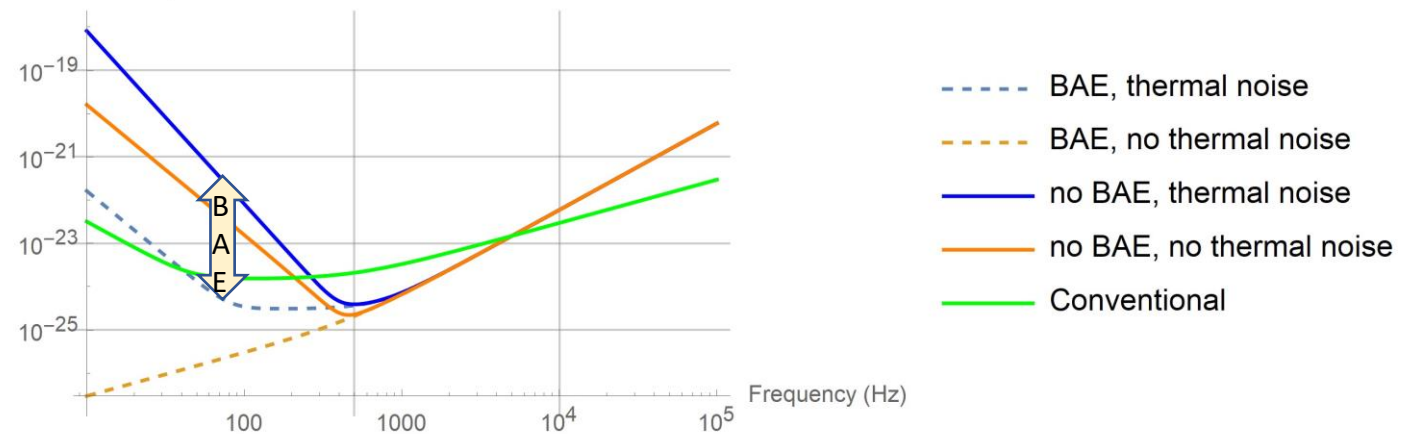


# GW Noise Spectra with backaction

- GW Noise Spectra normalized for signal

Parameter	Value
$\gamma_R$	$(2\pi)500$ Hz
$\gamma_m$	0.0000001 Hz
$\kappa$	$(2\pi)5000$ Hz
$\chi$	$(2\pi)5000$ Hz
$k_B$	$1.38065 \times 10^{-23} \text{ m}^2 \text{ s}^{-2} \text{ kg K}^{-1}$
$T$	4 K
$\mu$	50 kg
$\mu c$	-50 kg
$\omega_m$	$(2\pi)10$ kHz
$\omega_0$	$(2\pi)299792458 \text{ ms}^{-1} / 2\mu\text{m}$
$\hbar$	$1.054571628 \times 10^{-34} \text{ Js}$
$L$	4000 m
$c$	$299792458 \text{ ms}^{-1}$
$P_c$	$3 \times 10^6$ Watt

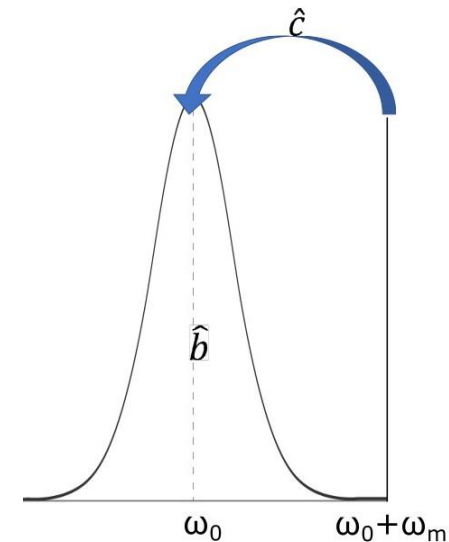
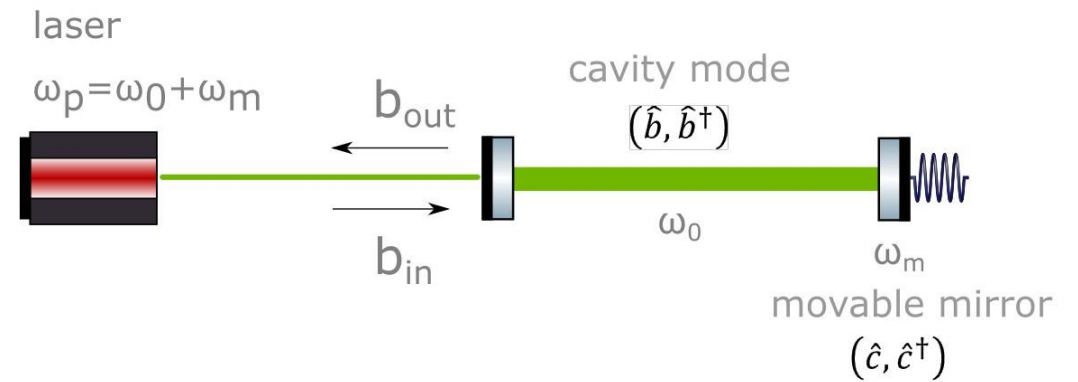
GW Noise Spectral Density  $\sqrt{S_h} (\text{Hz}^{-1/2})$



# Filter Cavity in FINESSE

# Filter Cavity

- Blue detuned pumping
- FP cavity with movable end mirror (mechanical oscillator)

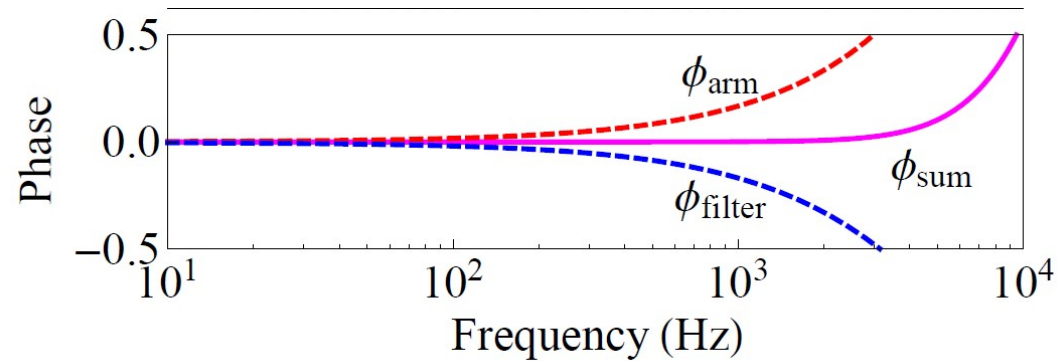


# Filter Cavity

- Anomalous Dispersion

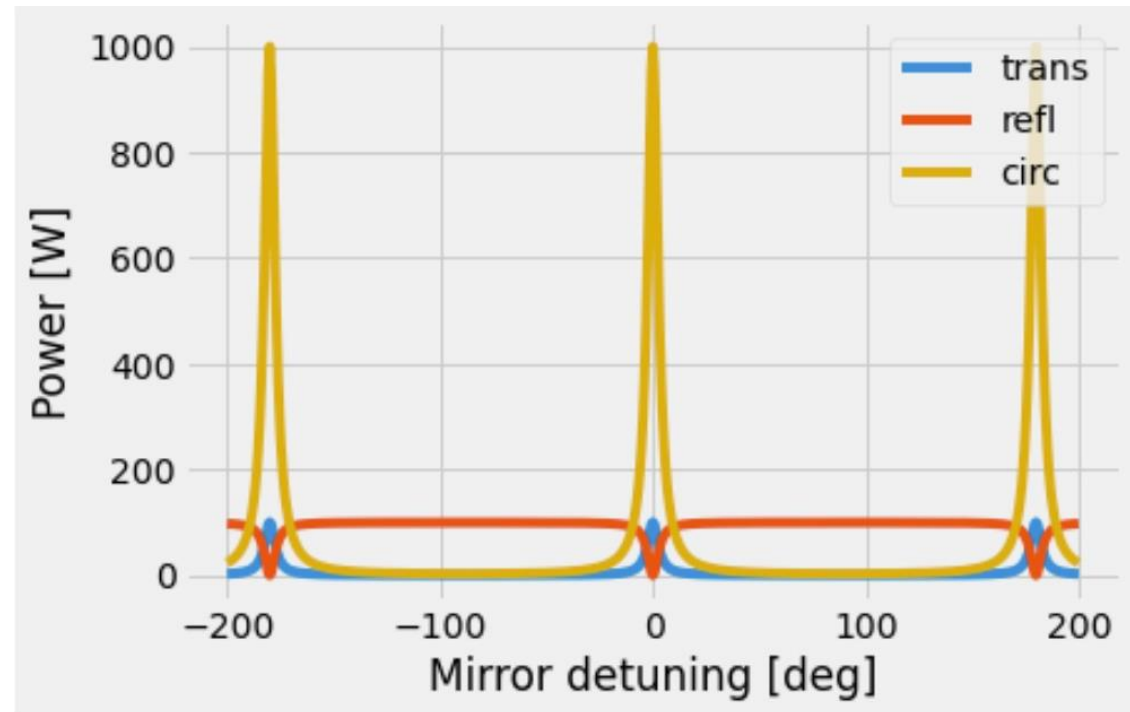
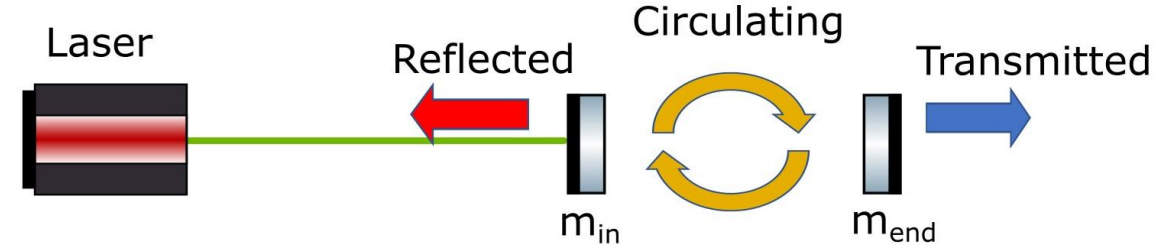
$$\hat{b}_{out}(\Omega) \approx \frac{\Omega + i\gamma_{opt}}{\Omega - i\gamma_{opt}} \hat{b}_{in}(\Omega)$$

$(\gamma_{opt} \gg \gamma_m)$



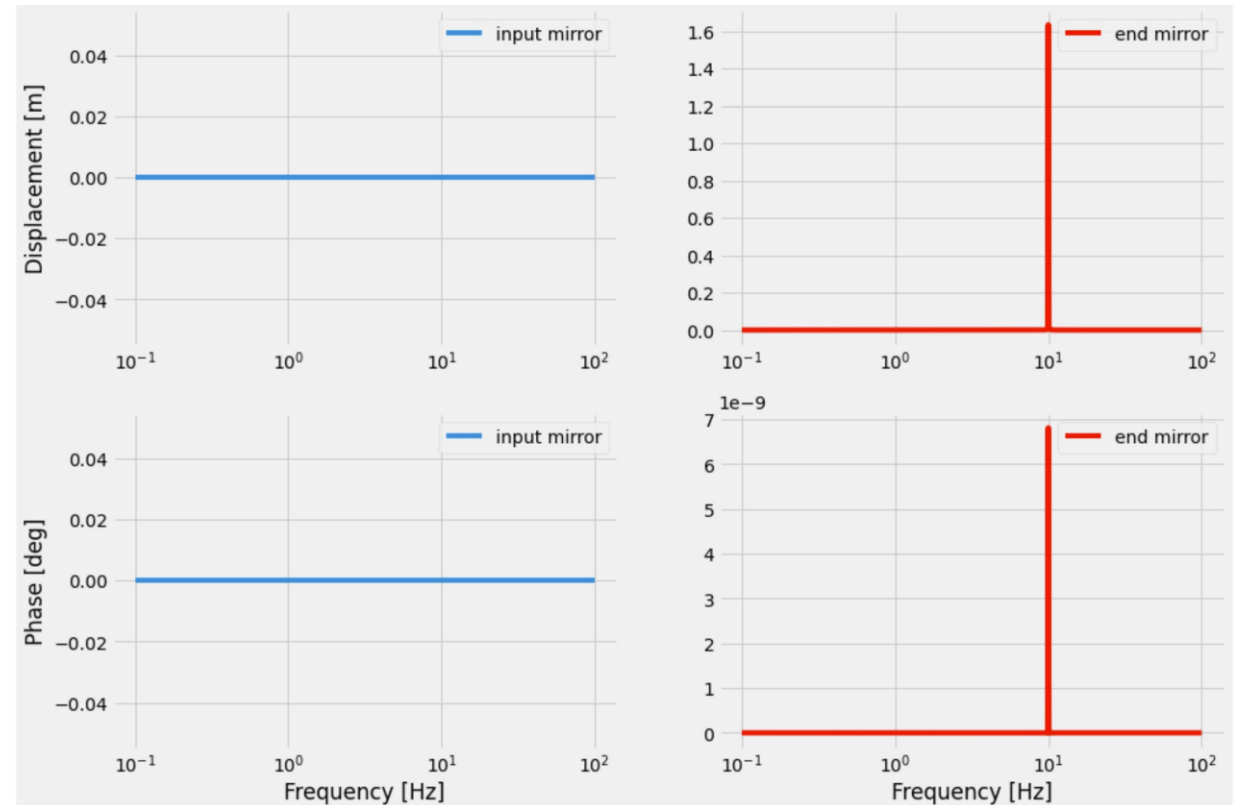
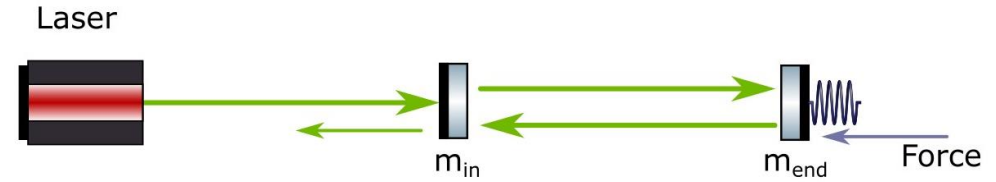
# Fabry Perot Cavity

- Resonance
- $P = 100\text{W}$
- $L = 4\text{m}$
- $T_{in} = T_{end} = 0.1$
- $L_{in} = L_{end} = 0$



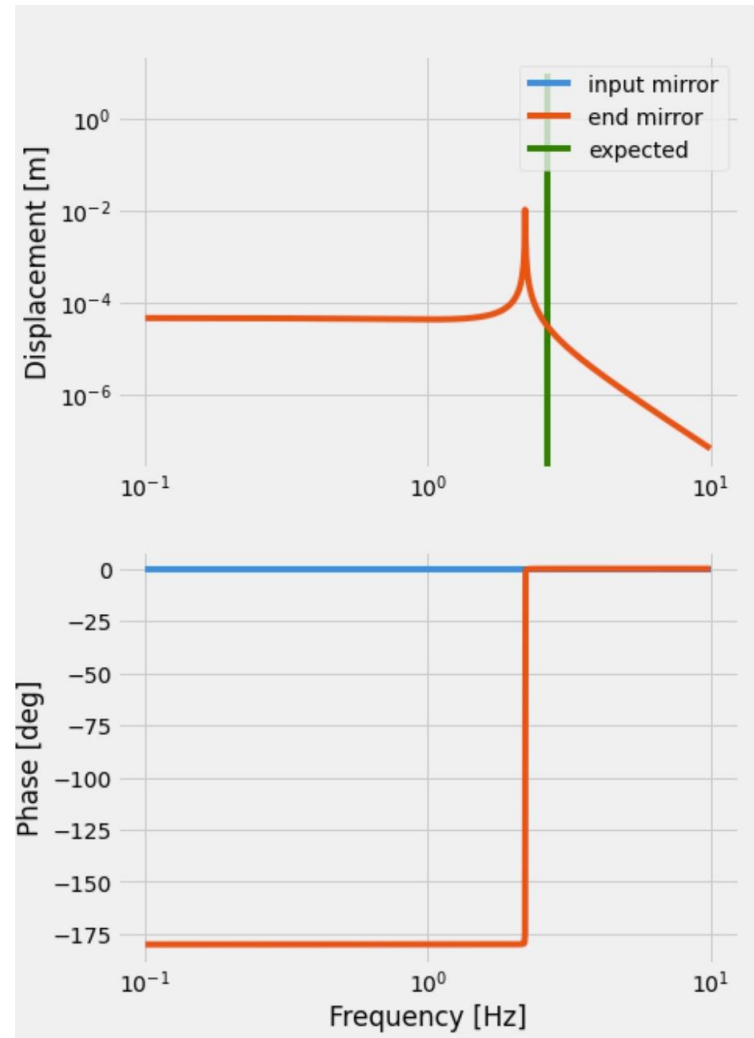
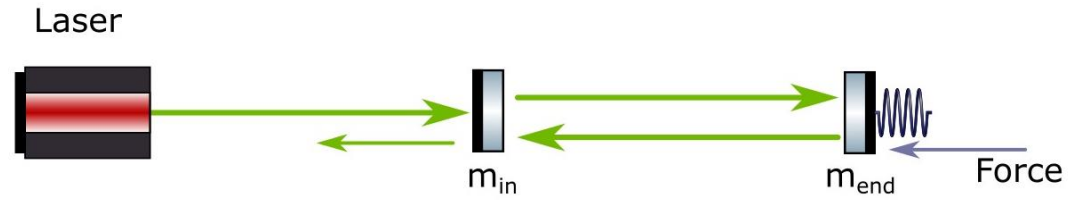
# Fabry Perot Cavity

- Mechanics without detuning
- $P = 3100 \text{ W}$
- $L = 50 \text{ m}$
- $\omega_m = (2\pi)1 \text{ Hz}$
- $\Delta = 0 \text{ Hz}$
- $m = 1 \text{ kg}$
- $Q_m = 8 \times 10^5$



# Fabry Perot Cavity

- Optical Spring Effect
- $P = 3100 \text{ W}$
- $L = 50 \text{ m}$
- $\omega_m = (2\pi)1 \text{ Hz}$
- $\Delta = (2\pi)300 \text{ Hz}$
- $m = 1 \text{ kg}$
- $Q_m = 8 \times 10^5$



$$\omega_{os} = \frac{2P_c}{\lambda_0 L} \left( \frac{\Delta}{\Delta^2 + \gamma^2} \right)$$

$$\omega = \sqrt{\omega_m^2 + \omega_{os}^2}$$

# Future Directions

- Realize anomalous dispersion for filter cavity
- Add the rest of the building blocks of the system (arm cavity, negative mass)
- Evaluate losses at different parts of the system



# References

- [1] Li, Xiang, et al. "Broadband sensitivity improvement via coherent quantum feedback with PT symmetry." *arXiv preprint arXiv:2012.00836* (2020).
- [2] Miao, Haixing, et al. "Enhancing the bandwidth of gravitational-wave detectors with unstable optomechanical filters." *Physical review letters* 115.21 (2015): 211104.