



# Minimising the Effect of Mirror Perturbations on Quantum Decoherence

LIGO SURF 2020

Swadha Pandey

Mentors: Jonathan Richardson, Rana Adhikari, Annalisa Allocca

#### Outline

- 1. Introduction
- 2. Background
- 3. Goals
- 4. Analytic Formalism
- 5. Numerical Optimisation of the aLIGO System
- 6. Conclusion
- 7. Further Work

## Introduction

## Motivation



**Figure 1:** The strain sensitivity of the LIGO detectors is currently limited above approximately 200 Hz by quantum noise.<sup>1</sup>

<sup>&</sup>lt;sup>1</sup>Publicly available image at https://www.ligo.caltech.edu/

Squeezed states of light: our saviour  $\Delta A \Delta \phi \geq \frac{\hbar}{2} \rightarrow$  Unevenly distributed



3 dB  $\uparrow \Longrightarrow$  3  $\times$  Event Rate!

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What is limiting the effective level of squeezing? **Optical losses** What do optical losses result in? Squeezed  $\rightarrow$  squeezed + unsqueezed What is causing these losses? Mode-mismatches! (More than 10% in aLIGO<sup>[1]</sup>) What causes these mode-mismatches? Perturbation in apparatus (R = 5.932418 m?!)

# Can we find an optimal set of design parameters such that the interferometer becomes minimally sensitive to design perturbations?

Perturbations  $\rightarrow$  curvatures and positions of the optical elements

Background

Laser beam  $\rightarrow$  Spatial intensity distribution

$$E(t, x, y, z) = \sum_{j} \sum_{n,m} a_{jnm} u_{nm}(x, y, z) \exp(i(\omega_j t - k_j z))$$
(1)

 $u_{nm}(x, y, z) \rightarrow \text{set of Hermite-Gauss (HG) or Laguerre-Gauss (LG)}$  polynomials



Figure 2: *u<sub>nm</sub>* spatial distributions in the LG basis

#### The Fundamental Gaussian Mode



**Figure 3:** Intensity pattern of a Gaussian beam (left) and the intensity and amplitude distributions of a normalised Gaussian beam (right).<sup>[3]</sup>

### Cavity Eigenmode and Mode-Mismatch



**Figure 4:** Cavity eigenmode: The beam curvature must be equal to the curvature of the mirrors at the mirror positions.<sup>[3]</sup>

## Goals

- Algorithm for numerically optimising any optical setup
- Implement algorithm to optimise the aLIGO design
- Analytic formalism to calculate loss

# Analytic Formalism

To develop an analytic formalism to calculate the total mode-matching loss in a complex optical system as a function of small perturbations of the optic positions and curvatures.

## Small curvature and position perturbations: $LG_{00} \leftrightarrow LG_{10}^{[4, 1]}$ Electric field at any point

$$|\Psi\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \alpha |LG_{00}\rangle + \beta |LG_{10}\rangle$$
(2)

#### **Mode-Mixing Matrices**



$$|\Psi_{\rm out}\rangle = A_2 \times A_1 \times A_0 \times |\Psi_{\rm in}\rangle$$
 (3)

#### Formalism and Simulations: Do They Agree?



Figure 5: Fabry-Perot cavity before the OMC.

### Formalism and Simulations: Do They Agree?



**Figure 6:** Analytic and simulation results for power loss percentage as a function of curvature perturbation percentage at M2.

#### Formalism and Simulations: Do They Agree?



**Figure 7:** Analytic and simulation results for power loss percentage as a function of position perturbation percentage at M2.

# Numerical Optimisation of the aLIGO System

### Optimising the aLIGO System



Figure 8: aLIGO setup

### Optimising the Signal Recycling Cavity



Figure 9: aLIGO X-arm cavity and Signal Recycling Cavity (SRC)



<sup>&</sup>lt;sup>2</sup>Laser and squeezer inputs remain mode-matched to the unperturbed cavity

- Δ*R* ~ *N*(μ = 0, σ = 0.01*R*) and Δ*z* ~ *N*(μ = 0, σ = 3mm), where ~ *N* indicates a normal distribution.
- Repeat 1,000 times
- Plot probability distribution of the squeezing level
- Compute the  $85^{\rm th}$  percentile

#### Expected Squeezing Level for aLIGO SRC

#### Effective squeezing level of unperturbed cavity = 9.8dB



Figure 10: Monte Carlo analysis for perturbed cavity

	$R_{\rm SR3}$	$\mathrm{R}_{\mathrm{SR2}}$	$\mathbf{R}_{\mathrm{SRM}}$	$L_{\rm SR3}$	$L_{\rm SR2}$	$L_{\rm SR1}$
aLIGO (m)	35.97	-6.41	-5.69	19.37	15.44	15.76

20

## Pertrubation Analysis for $R_{\text{SRC}}$

#### Effective squeezing level of unperturbed cavity = 9.8dB



Figure 11: Perturbing only the  $R_{SRC}$ 

	$R_{\rm SR3}$	$\mathrm{R}_{\mathrm{SR2}}$	$\mathbf{R}_{\mathrm{SRM}}$	$L_{\rm SR3}$	$L_{\rm SR2}$	$L_{\rm SR1}$
aLIGO (m)	35.97	-6.41	-5.69	19.37	15.44	15.76

21

### Expected Squeezing Level for Optimised SRC

#### Effective squeezing level of unperturbed cavity = 9.7dB



Figure 12: Monte Carlo analysis for perturbed cavity

	$\mathrm{R}_{\mathrm{SR3}}$	$R_{\mathrm{SR2}}$	$\mathbf{R}_{\mathrm{SRM}}$	$L_{\rm SR3}$	$L_{\mathrm{SR2}}$	${\rm L}_{{ m SR1}}$
Optimised (m)	59.66	-233.60	-3.53	27.69	18.98	9.87

22

Conclusion

- Developed algorithm to minimise sensitivity of any optical setup.
- Optimised the aLIGO Signal Recycling Cavity
- Quantified the improvement.
- Developed analytic formalism.
- Checked agreement between formalism and simulations.

Next Steps

The next steps for this project are:

- Apply to the entire aLIGO setup.
- Take into account the LIGO thermal compensation system.

## Thank You

My sincere thanks and gratitude to my mentors and to LIGO Laboratory for giving me this opportunity.

# **Questions?**

Helmholtz equation:

$$\nabla^2 \mathsf{E} - \frac{\ddot{\mathsf{E}}}{c^2} = 0 \tag{4}$$

The fundamental Gaussian mode at a position z:

$$u_{00}(r,z) = \sqrt{\frac{2}{\pi}} \frac{1}{\omega(z)} \exp(\iota\psi(z)) \exp\left[-r^2\left(\frac{1}{w^2(z)} + \iota\frac{\pi}{\lambda R(z)}\right)\right]$$
(5)

Reflection from a mirror having amplitude reflectivity coefficient r

$$\mathbf{r} = \begin{pmatrix} r & 0\\ 0 & r \end{pmatrix} \tag{6}$$

Modified reflection matrix for a small curvature perturbation  $\delta R$  in a mirror of curvature R

$$\mathbf{r}' = \mathbf{a}.\mathbf{r} = \begin{pmatrix} r\sqrt{1-a^2} & -\iota ra\\ -\iota ra & r\sqrt{1-a^2} \end{pmatrix}, \quad a = \frac{\pi\omega^2(z_m)}{2\lambda R^2}\delta R$$
(7)

where  $\omega(z_m)$  is the beam size at the mirror position  $z_m$ 

The scattering matrix for transmission of a beam through a mirror having amplitude transmissivity coefficient *t* is given by

$$\mathbf{t} = \begin{pmatrix} t & 0\\ 0 & t \end{pmatrix} \tag{8}$$

If one the one-way propagation of a beam across a distance or cavity accumulates a phase  $\phi_0$  in the  $LG_{00}$  mode and  $\phi_1$  in the  $LG_{10}$  mode, this phase accumulation during propagation can be represented by the scattering matrix

$$\phi = \begin{pmatrix} e^{\iota\phi_0} & 0\\ 0 & e^{\iota\phi_1} \end{pmatrix} \tag{9}$$

The modified propagation matrix for a small position perturbation  $\delta z$  in a mirror of radius of curvature *R* 

$$\phi' = \mathbf{b}.\phi = \begin{pmatrix} \sqrt{1 - b^2} e^{\iota \phi_0} & -b e^{\iota \phi_1} \\ -b e^{\iota \phi_0} & \sqrt{1 - b^2} e^{\iota \phi_1} \end{pmatrix}, \quad b = \frac{\delta z}{R}$$
(10)

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