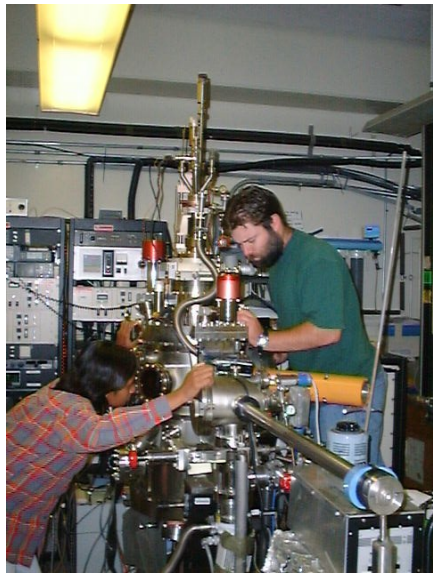


# Two level systems and thermal noise

*why do coatings have this noise and how do we reduce it*



UHV growth chambers for prototyping materials

**Frances Hellman**  
University of California, Berkeley  
New LIGO Coatings group member

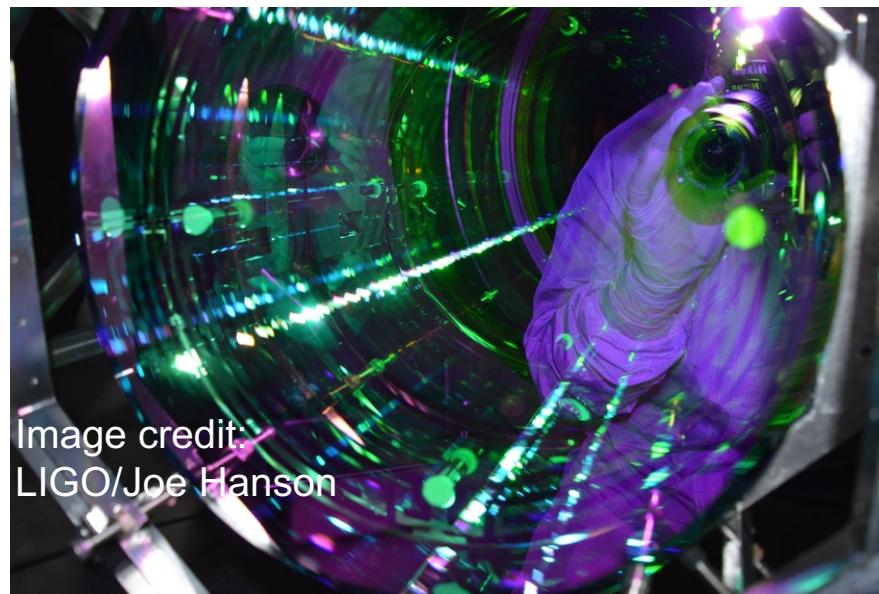
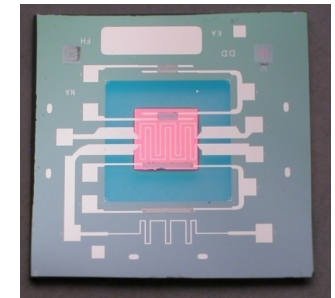


Image credit:  
LIGO/Joe Hanson



Nanocalorimeters:  
Si micromachined chips for thin film heat capacity



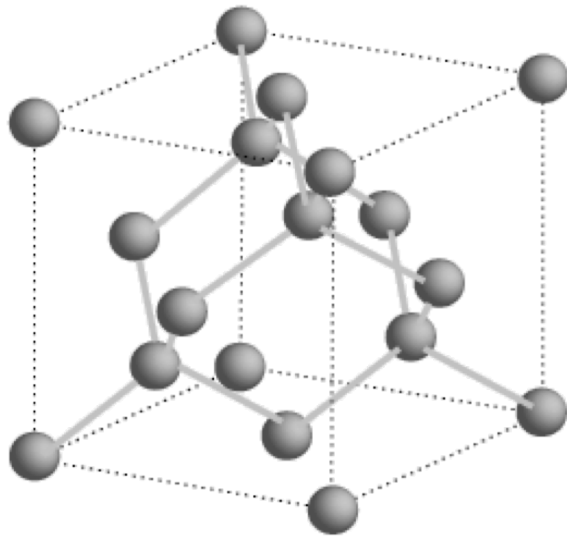
GORDON AND BETTY  
**MOORE**  
FOUNDATION

**LSC** Center for Coatings Research

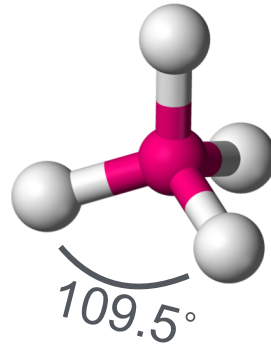


**Current mirror coatings are amorphous (=non-crystalline)  
No long range structural order, but have short range order**

**Traditionally, glasses quenched from liquid  
Also can be made by vapor deposition  
Structure of Silicon**



Xtal Si: diamond structure



Amorphous Si:  
still has  
tetrahedral  
coordination

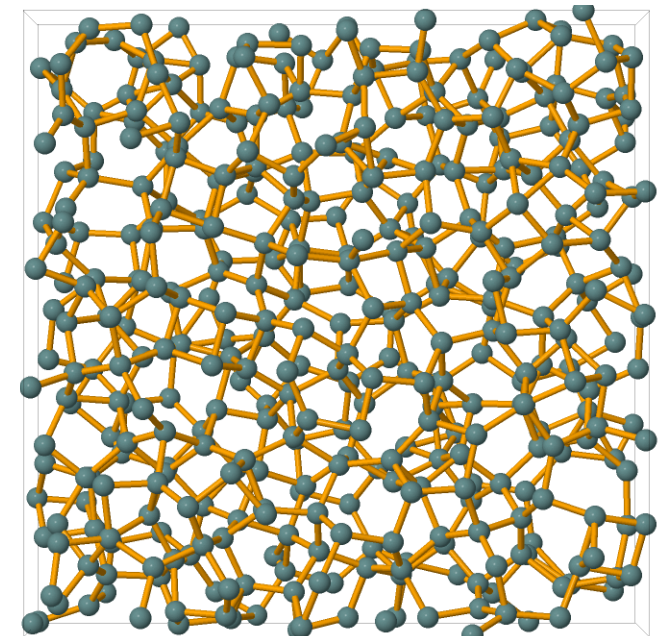
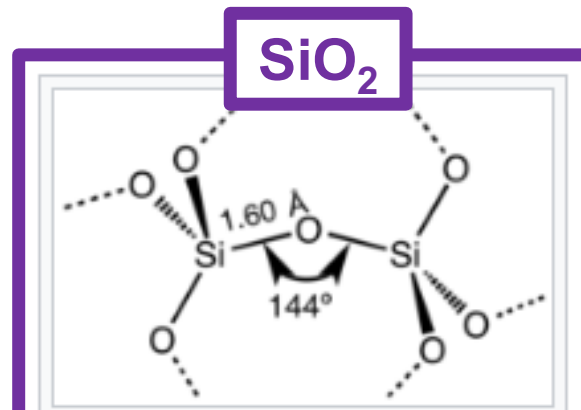
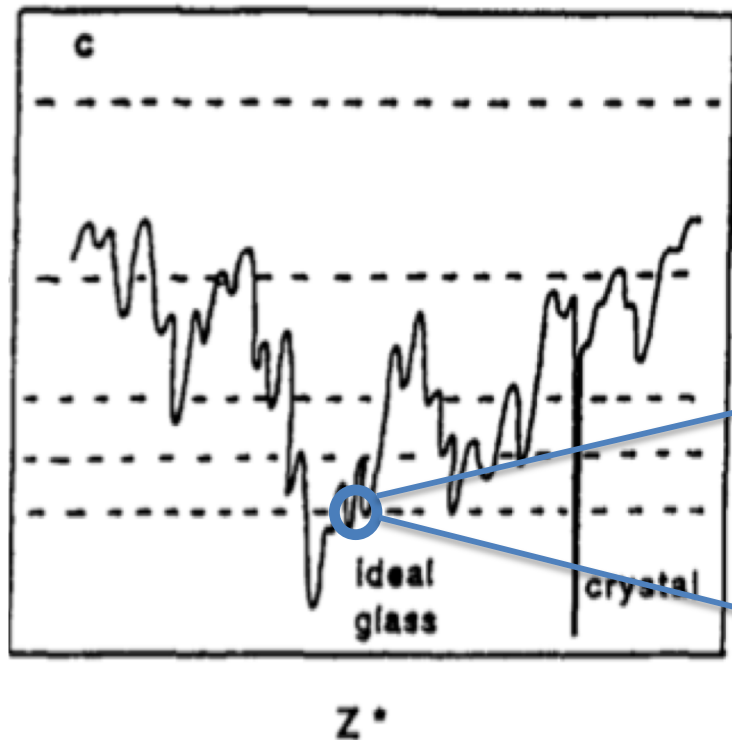


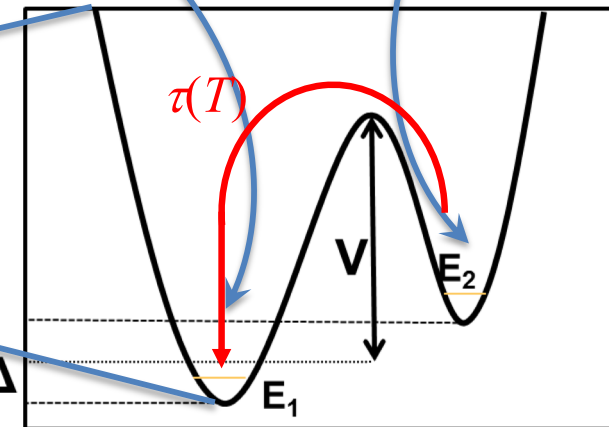
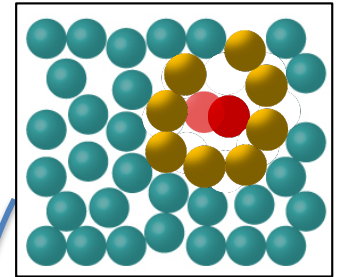
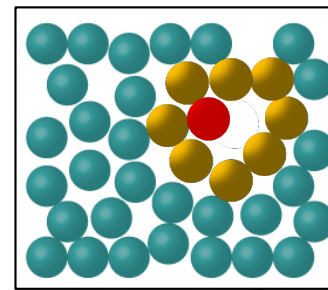
Image credit: Kiran Prasai

**O-Si-O bonds are fixed angle but Si-O-Si angle is quite floppy –  
different energy scales**

# Energy landscape of configurations: “nearby” minima lead to tunneling or thermally-activated motion of groups of atoms



POTENTIAL ENERGY



Anderson, et al., *Philos. Mag.*, **25** 1 (1972)  
Phillips, *J. Low. Temp. Phys.*, **7** 3-4 (1972)

C.A. Angell, *Physica D* **107**, 122 (1997)

## Two-Level Systems from neighboring energy minima in structural landscape:

- At low  $T$ , atomic structure **tunnels** between these  $\gg \mu\text{eV}$  energy splitting  $E_{1,2} \pm \Delta$
- At higher  $T$ , atomic motion is **thermally activated**, requiring  $k_B T \sim$  barrier height  $V$

In both cases, atomic motion leads to dissipation (thermal noise)

For a single  $V$ , dissipation at frequency  $\omega$  will have a peak at temperature  $T$

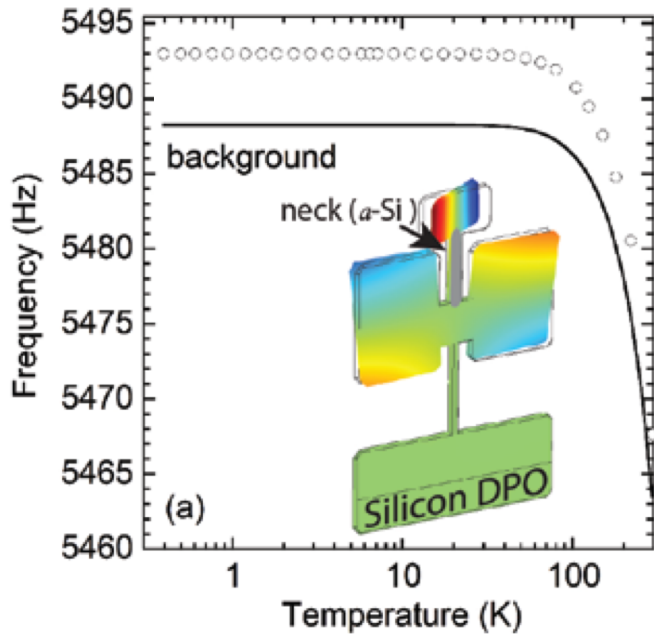
➤ at 1 kHz, 0.5 eV barrier heights has a peak  $\sim$  room temperature

50 meV barrier heights has a peak  $\sim$  30K

A *distribution* of  $\mu\text{eV}$  tunneling-induced energy splitting leads to  $T$ -independent losses



# Internal friction to measure losses



Measure resonant frequency and internal friction (damping)  $Q^{-1}$  of a xtal Si double paddle oscillator (DPO) or cantilever as a function of  $T$  before and after depositing a film, or a silica disk at room temperature using GeNS (gentle nodal suspension)

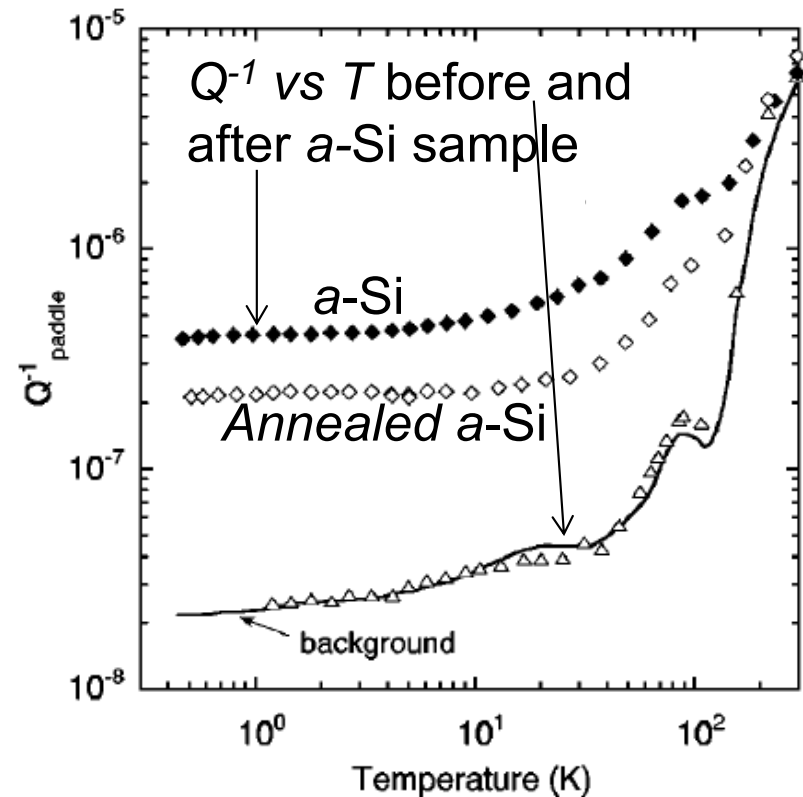
X. Liu and R.O. Pohl, Phys. Rev. B **58**, 9067 (1998)

Internal friction  $Q^{-1}(T)$  has a low  $T$  plateau due to *tunneling two level system (TLS)-phonon interactions*

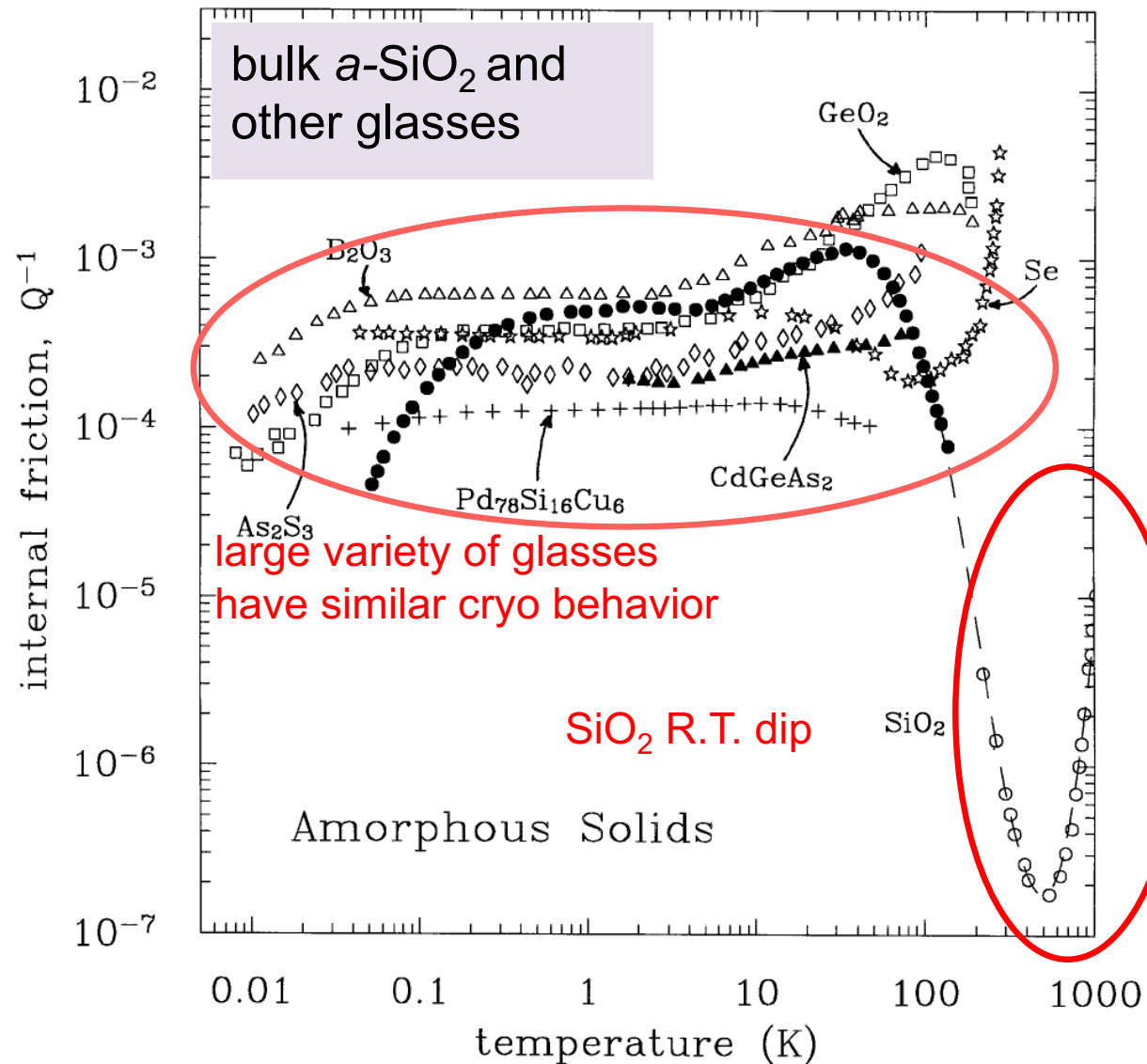
$Q_0^{-1}$  proportional to  $\bar{P}$  (density of TLS) with a *poorly understood TLS – phonon coupling parameter  $\gamma$* , also called the deformation potential

$$Q_0^{-1} = \pi \bar{P} \gamma^2 / 2 \rho v^2$$

$\rho$  is density, and  $v$  is sound velocity

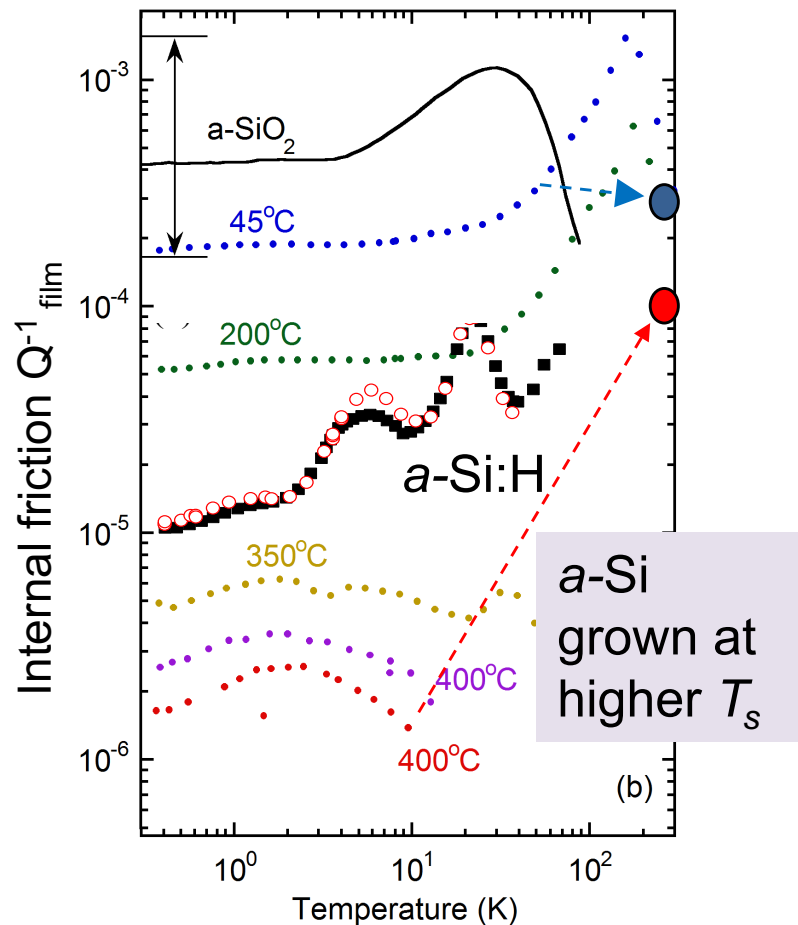


# “Universal” mechanical losses at low T (tunneling) Higher T more variable, including peaks (thermally activated)



K.A. Topp, *Z. Physik B Condensed Matter* 101 235–45 (1996)

# Thin film amorphous Silicon: internal friction $Q^{-1}$ (and excess heat capacity) is strongly reduced (decades) by increased growth $T_s$



$a\text{-Si}$  films grown at lower  $T_s$  have typical (high)  $Q^{-1}$  (like  $a\text{-SiO}_2$ )

$a\text{-Si}$  films grown at 400° C have very low  $Q^{-1}$  at low  $T$  (perhaps good for LIGO Voyager); (at RT better but only 4x)

BUT, large absorption – likely need H to eliminate dangling bonds in  $a\text{-Si}$

$a\text{-Si:H}$  has low  $Q^{-1}$  (although peaks not well understood) BUT has high excess low  $T$  heat capacity indicating local structures that cause loss are present, just not coupling to acoustic waves – why and what does this do to higher temp loss???

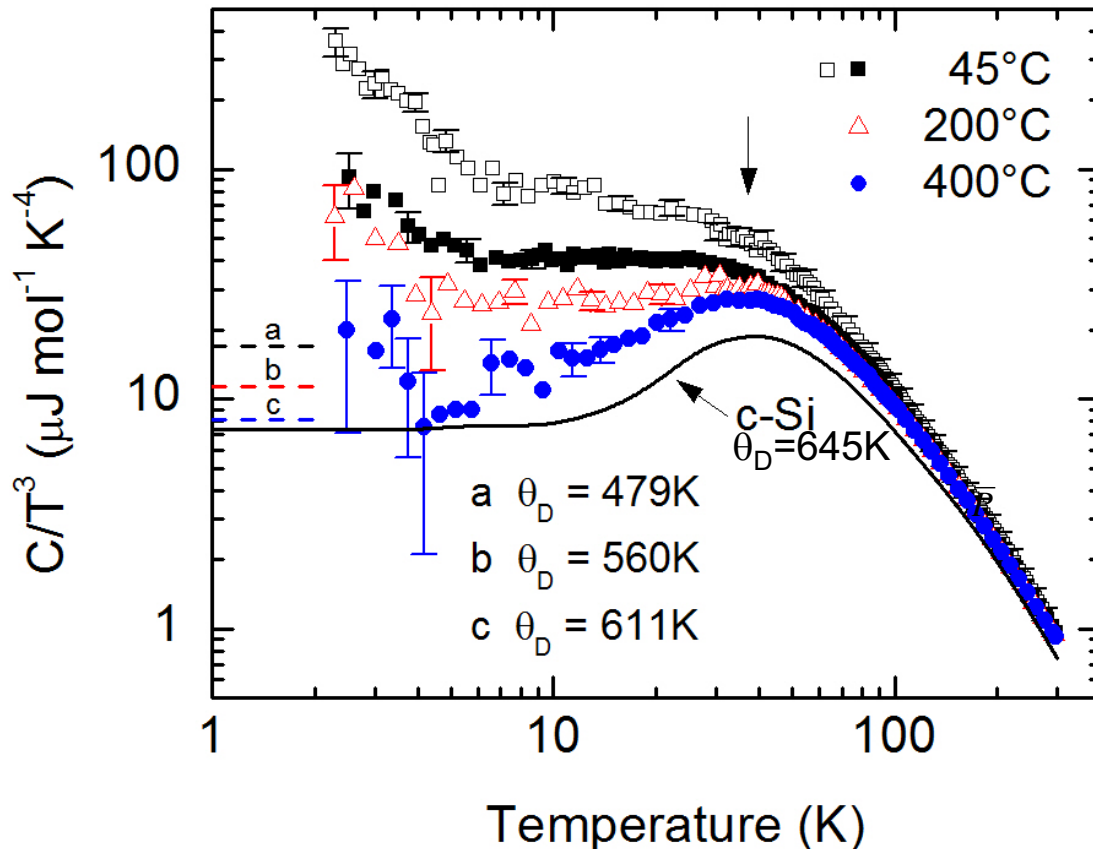
D. R. Queen, X. Liu, J. Karel, T.H. Metcalf, F. Hellman, Phys. Rev. Lett. 110, 135901 (2013).

X. Liu, D. R. Queen, T. H. Metcalf, J. E. Karel, F. Hellman, Phys. Rev. Lett. 113, 025503 (2014).

D.R. Queen, X. Liu, J. Karel, Q. Wang, R.S. Crandall, T.H. Metcalf, F. Hellman, Eur. Phys. Lett. **112**, 26001 (2015).

M. Molina-Ruiz, H. C. Jacks, D.R. Queen, T. Metcalf, X. Liu, Q. Wang, R. S. Crandall, F. Hellman, in preparation

# Variable TLS in a-Si also in specific heat C: excess C at low T



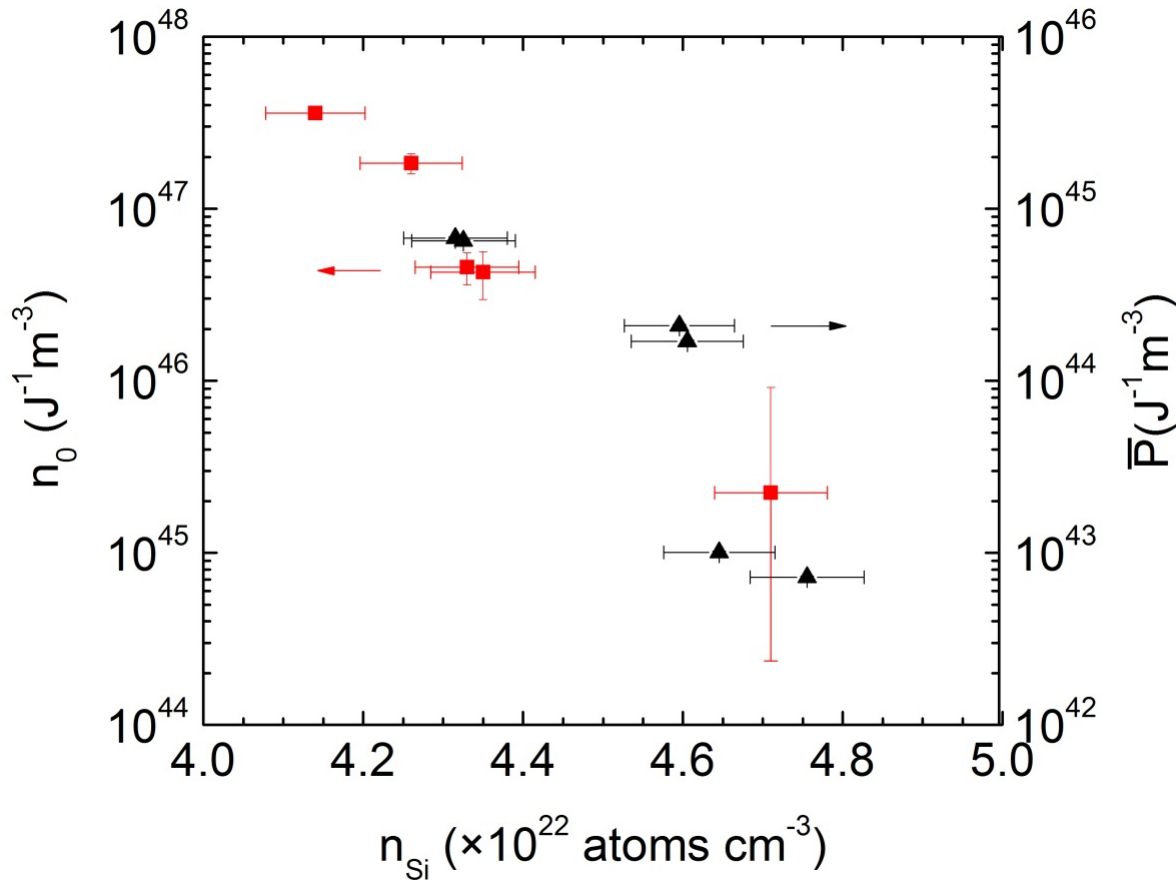
*Linear term in C:  $c_1 \sim n_0$*   
*Excess  $T^3$  term ( $c_{ex}$ )*

Films grown at 400° C have  $C(T)$  only a little above c-Si; small  $n_0$  and small  $c_{ex}$   
 (Also, thermal conductivity shows no plateau)

Films grown at lower  $T_s$  have excess  $Q^{-1}$  and  $C(T)$  above Debye value (from  
 transverse and longitudinal sound velocity measurements)

Fit low T  $C(T)$  to  $c_1 T + C_3 T^3$ ; both  $n_0$  and  $c_{ex}$  depend on  $T_s$  and on film thickness  
 Sound velocity, bond angle distribution, nanovoid size depend on  $T_s$  & **not** thickness

# TLS density from specific heat and internal friction are proportional to each other, *and depend on film density*

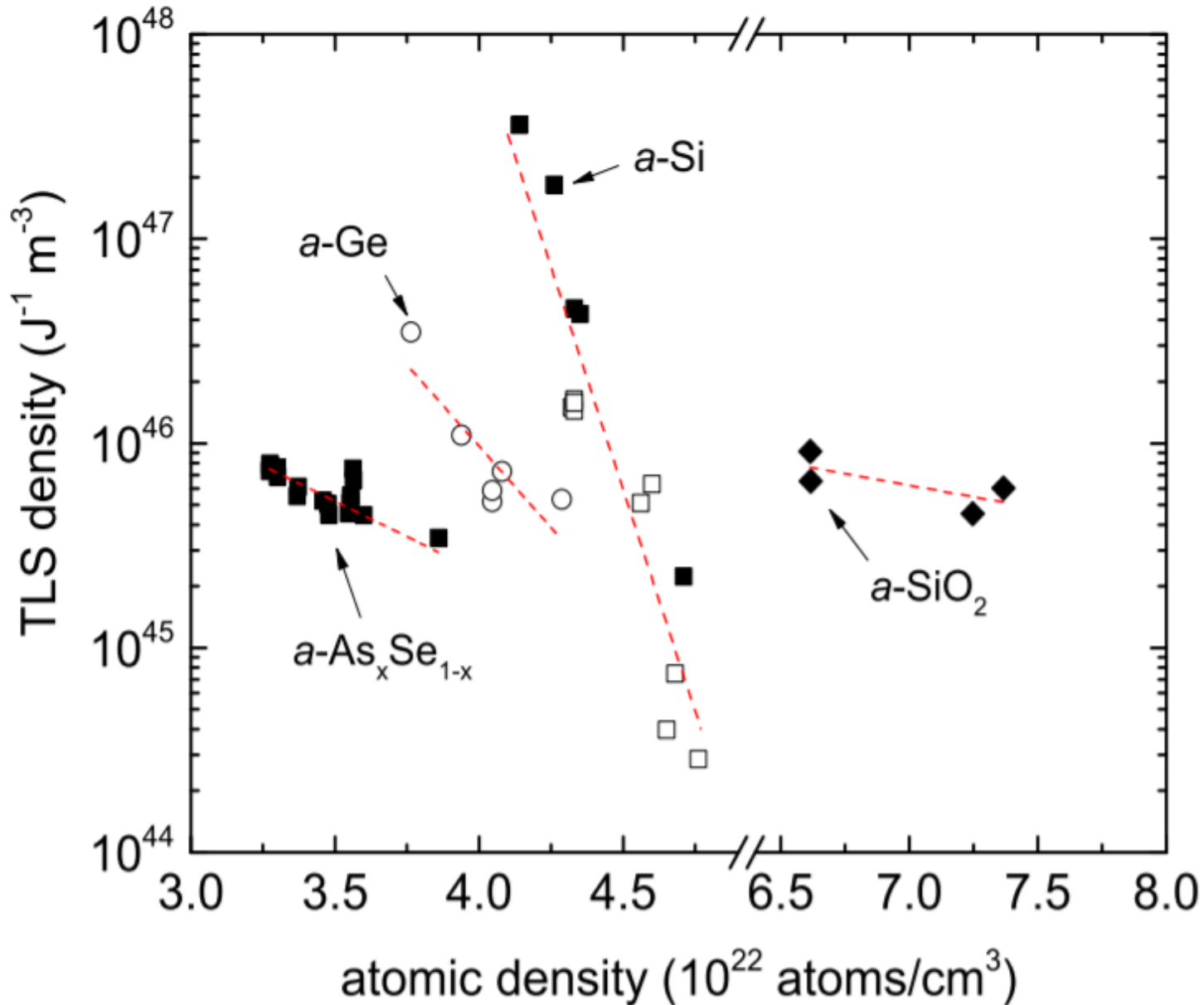


Crystalline Silicon:  
 $n_{Si} = 5 \times 10^{22} \text{ cm}^{-3}$

- $n_o$  and  $\bar{P}$  vanish as  $n_{Si} \rightarrow n_{\text{crystalline Si}}$
- $n_o \sqrt{\bar{P}} \sim 8$  – similar to other glasses.  $n_o$  and  $\bar{P}$  proportional in usual TLS model
- TLS vanish with increasing  $n_{Si}$  – associated with low density regions/nanovoids??
- Correlation is over nearly 3 decades



# TLS (either $n_o$ or $\bar{P}$ ) dependence on atomic density seen in a range of amorphous materials





# Hypotheses re vapor deposited *a*-Si

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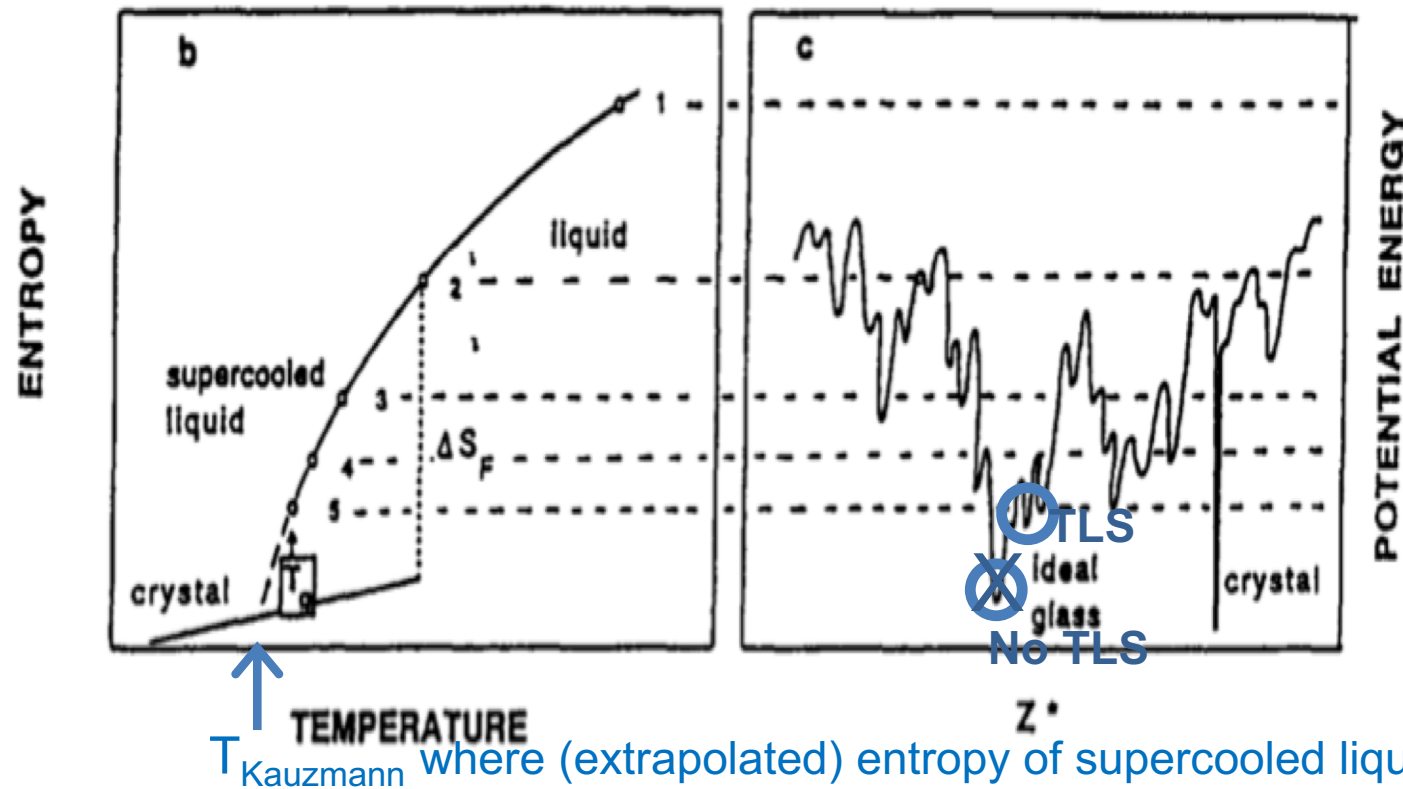
Vapor deposited films of covalent materials such as *a*-Si or *a*-SiO<sub>x</sub> (or Alumina or SiN<sub>x</sub> or Tantalum) have to date not been probed for ideality/ultrastability

The glass transition of *a*-Si has never been measured (because it can't be quenched) but theory suggests 850K (C.R. Miranda and A. Antonelli, J. Chem Phys 120, 11672 (2004)).

**Our growth T to get low TLS is 673K ~ 0.8 T<sub>g</sub>!! (similar to IMC work)**

We have also seen effects of deposition rate and thickness on density and TLS, similar to IMC work; TLS measurements in progress on these other films

# Connection between energy landscape, entropy, and TLS



C.A. Angell, Physica D  
107, 122 (1997)

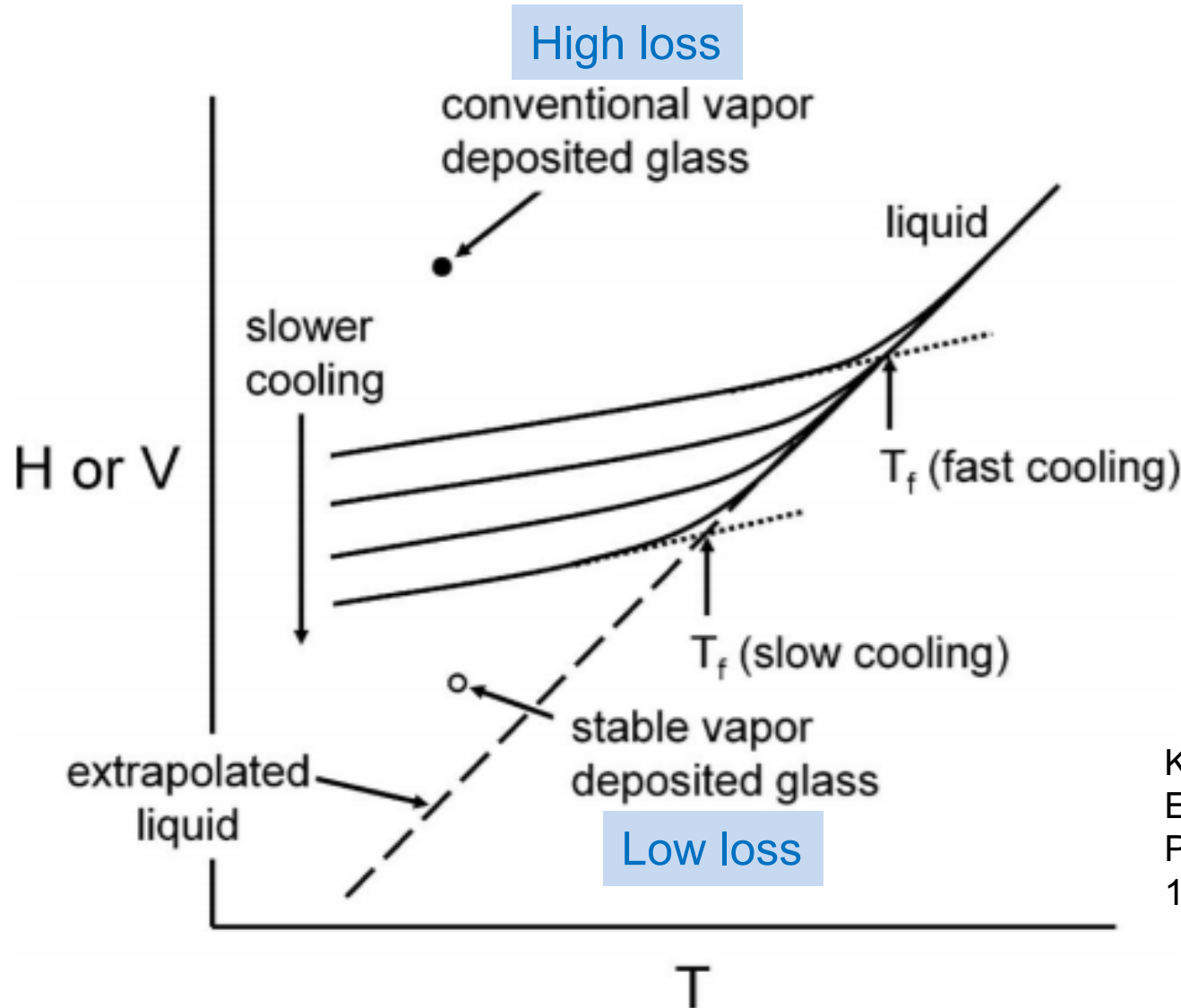
$T_{\text{Kauzmann}}$  where (extrapolated) entropy of supercooled liquid crosses crystal!

The energy landscape (right) as related to the glass transition of a liquid (left). Glasses falling out of the equilibrium supercooled liquid at a given dashed line correspond to configurations in the energy landscape.

Hypothesis: vapor deposition offers a way to directly access low lying (ideal) glass state  
Due to high atomic mobility at film growth surface despite being at low  $T$ .

Hypothesis: Ideal glass has no nearby energy minima, so no TLS, unlike most other states

# Enthalpy or Volume (density) as a function of T starting from liquid Vapor deposited compared to liquid quenching of amorphous material



H and V based on simulations of atoms with simple bonds (Lennard-Jones). Kiran Prasan will talk about *a*-Si.

Kearns, Swallow, Ediger, J. Chem. Phys. 127, 154702 (2007)

**Could this apply to vapor deposited amorphous  $\text{SiO}_2$ ?  
Does this lower losses at all T, some T, for all materials?**



# Comments and Open questions

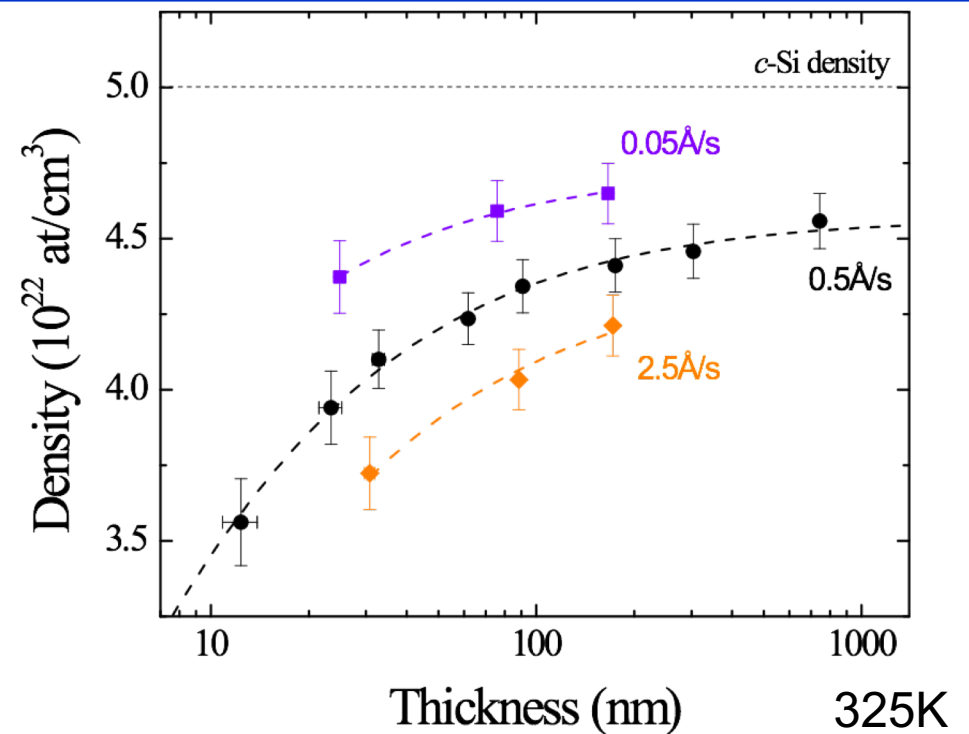
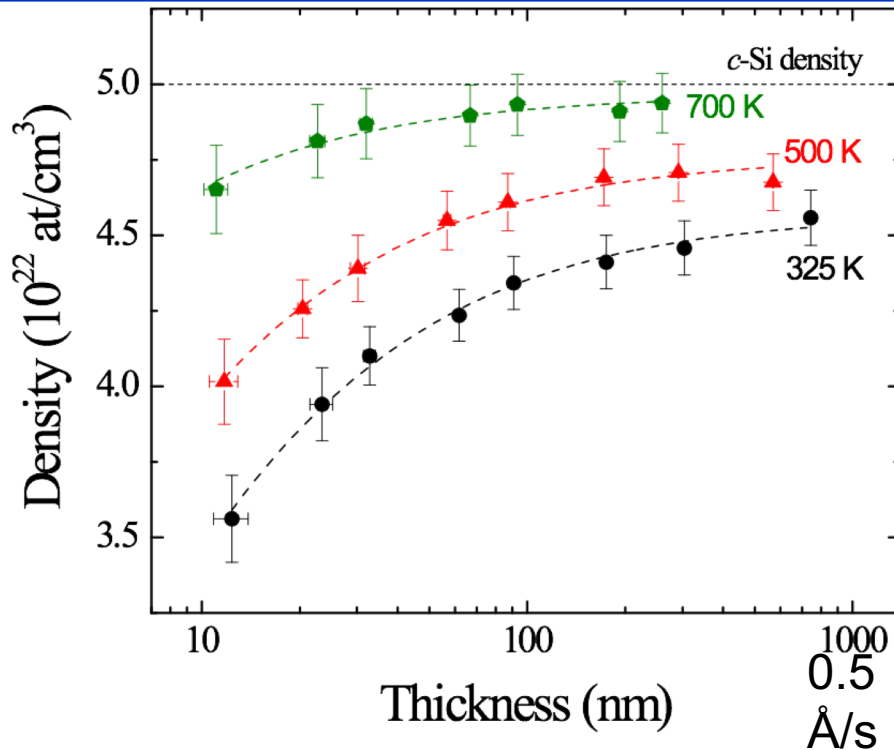
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- Random first order transition (RFOT) theory of glass transition predicts divergence of correlation length  $\xi$  of supercooled liquid as  $T \rightarrow T_K$
- Also predicts TLS density  $\sim \xi^{-3}$ , supporting the hypothesis that TLS density is highly suppressed for deposition  $T$  near  $T_K$
- But, surface diffusion needs to be fast near  $T_K$ , argues  $T_K$  needs to be  $\sim T_g$ , (which means “fragile” glasses like *a*-Si and indomethacin, not *a*-SiO<sub>2</sub>)
- *But, a*-SiO<sub>2</sub> has a high density of low energy floppy modes, so perhaps

## OPEN QUESTIONS

- Is low TLS related to growth near  $T_K$ ? (If (and only if) surface mobility during growth is high). Fragile glasses have  $T_K$  near  $T_g$ , where mobility is high, so low TLS would be correlated with fragility
- Or is low TLS related to nature of bonding: overconstrained (tetrahedral Si) versus underconstrained (Si-O-Si bonds in SiO<sub>2</sub>)
- Test with amorphous SiO<sub>2</sub>, Ta<sub>2</sub>O<sub>5</sub>, Al<sub>2</sub>O<sub>3</sub>, and Se<sub>x</sub>Ge<sub>1-x</sub> alloys!!
- What is the relationship to density in the IBS overdense films, which *could* be higher in the energy landscape with higher TLS.

# Growth parameters substantially modify amorphous Si film density and some measures of structure



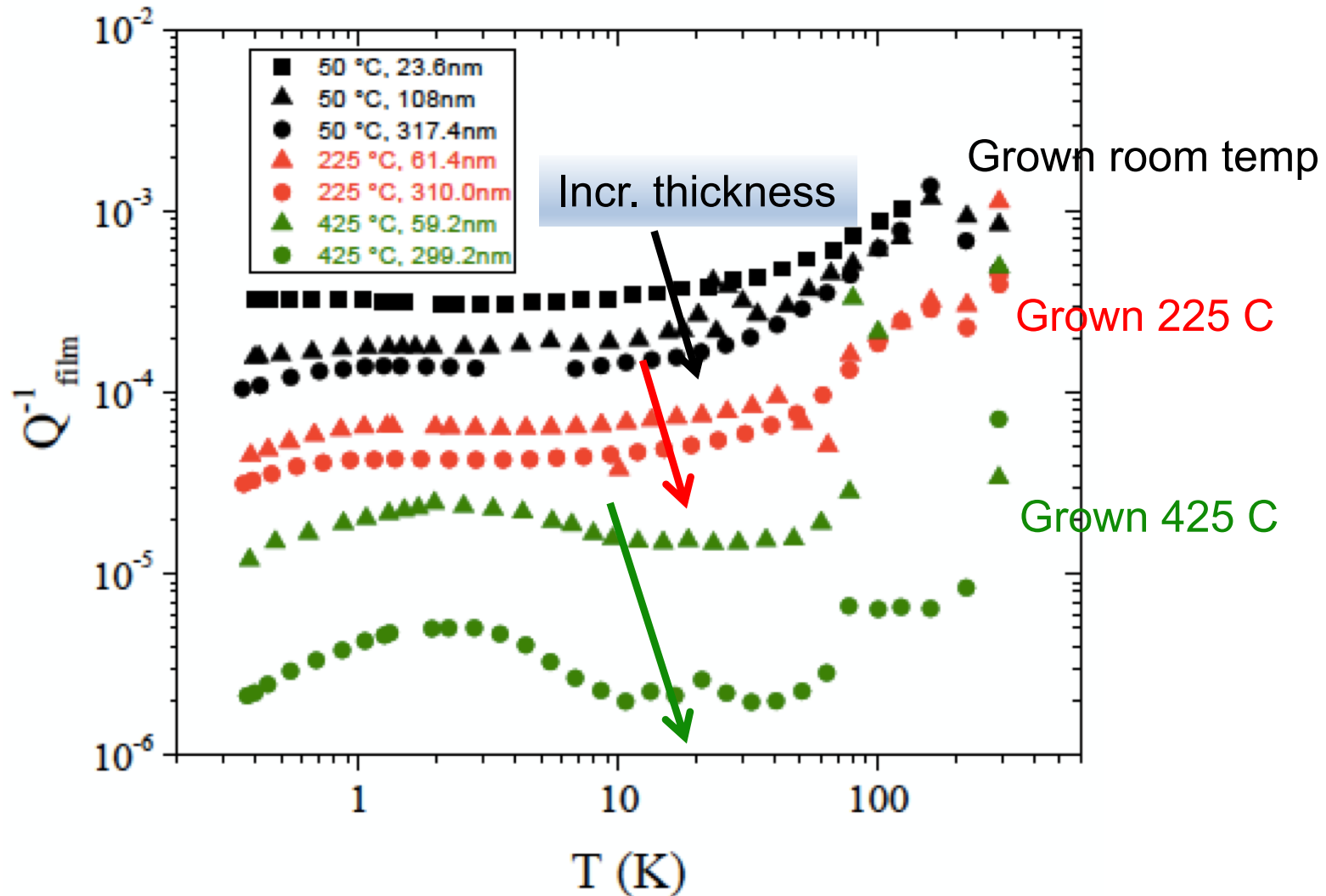
- Thickness, growth temperature  $T_g$ , and growth rate affect film density and roughness; room T growth flattest for all thicknesses; higher  $T_g$  thin is flat, roughens with thickness (1.5 nm RMS at 300 nm)
  - Thinner, low growth T, high growth rate films are less dense
  - On what length scale(s) do density changes occur? Little variation in dangling bond density or macroscale structure
  - *Variations in bond angle disorder, medium range order, nanovoid size and number* (Raman, Fluctuation Electron Microscopy, positron doppler broadening spectroscopy)
- Which of these matter to TLS? Structural work (**Kiran Prasai and Hai-Ping Cheng**)

# Amorphous Silicon losses: thickness also matters

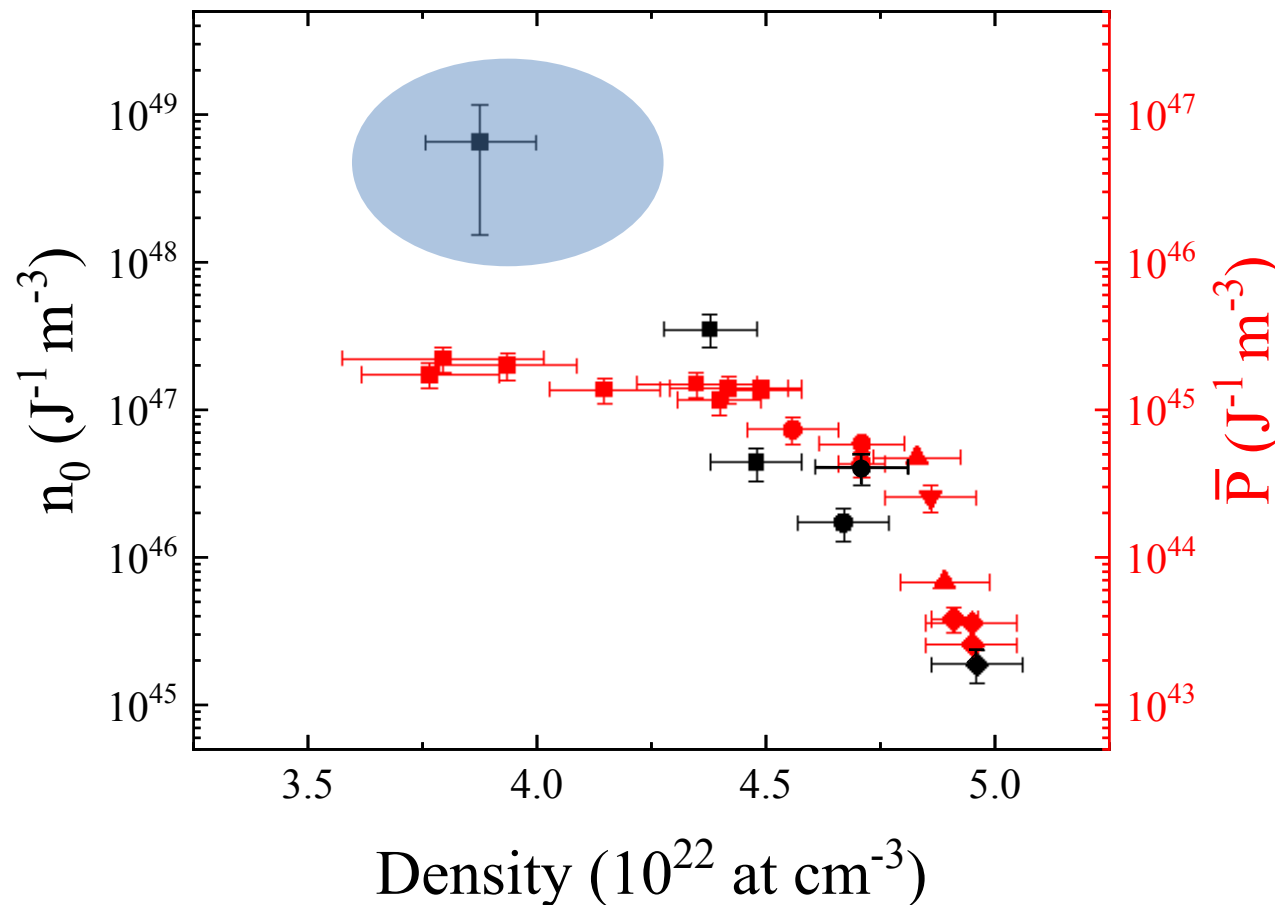
Thin films are more lossy than thick films (not per volume, absolute)

Correlates with atomic density differences (thick films are denser)

Annealing reduces loss, but not much (at low T) compared to growth T effects



# More recent data on internal friction (IF) derived TLS density (specific heat still in progress)



Low density plateau in  $\bar{P}$  shows that IF-derived TLS do *not* continue to increase with lower density samples

Two possible conclusions:

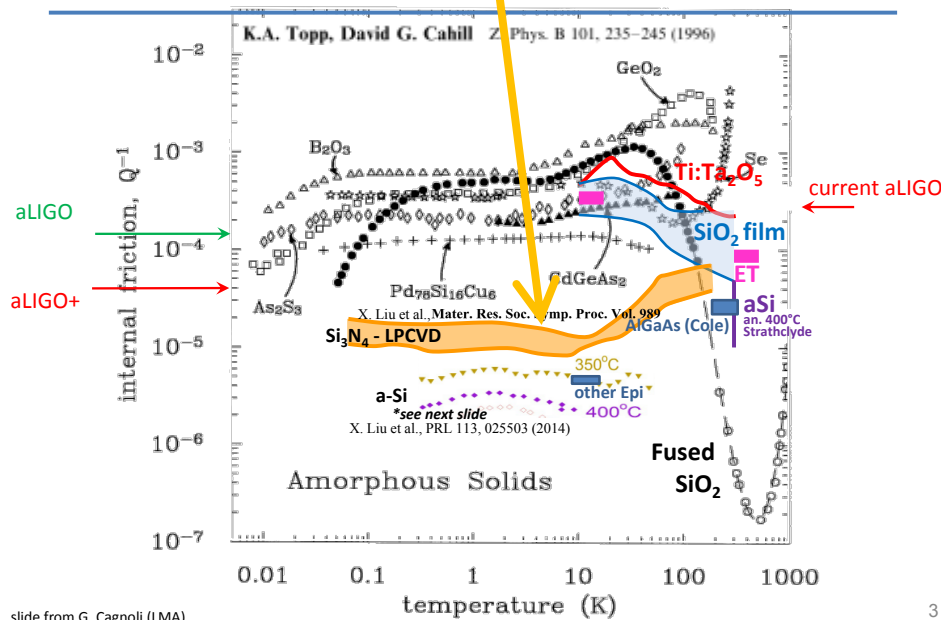
- Larger nanovoids in the lower density (thinner, faster growth rate, lower growth temperature) do not create more TLS (then specific heat  $n_0$  would also plateau)
- TLS decouple from phonons in lower density films (then specific heat  $n_0$  would continue to increase) – one data point suggests this low  $\gamma$  idea

# Other current materials in LIGO coatings:

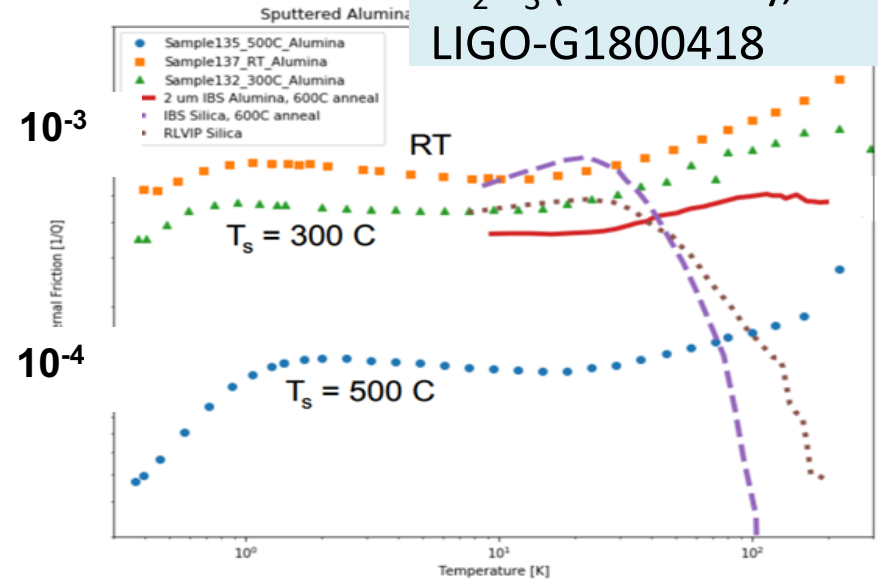
**Alumina** (amorphous  $\text{Al}_2\text{O}_3$ ) has reduced losses (at all T) when grown at elevated  $T_s$

**Amorphous  $\text{Si}_3\text{N}_4$**  – grown by LPCVD, at elevated T. Losses  $< 1 \times 10^{-4}$  at all T

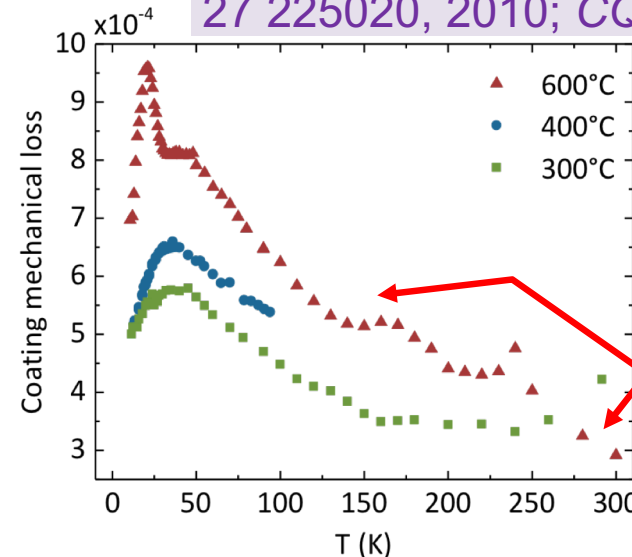
Jessica, Glasgow, LVC G1800585-1



Magnetron sputtered  $\text{Al}_2\text{O}_3$  (Abernathy, NRL) LIGO-G1800418



I W Martin et al, Class. Quant. Grav. 27 225020, 2010; CQG 35 075001



Discussed in next talk  
Kiran Prasai

Annealing improves loss at 300K, increases it at low T

**Tantala ( $\text{Ta}_2\text{O}_5$ )**: annealing more effective, growth T less effective – why?

- Elevated temperature growth of IBS tantala:  $T_s$  up to 400 C) reduces RT loss
- RT loss decreases with  $T_{\text{anneal}}$  until limited by crystallization



# Overview and future LIGO- specific directions

---

**Room T:** Tantalum is improving! Have  $\sim 1 \times 10^{-4}$  loss materials via annealing & doping

- Understand better the structural causes of loss to enable further improvements (Kieran Prasai talk)
- Look for a different high index material that has bond energies like  $a\text{-SiO}_2$  with really low losses. The key seems to be separation of bond energy scales
- Look for a narrow gap (high index) *fragile* glass like  $a\text{-Si}$  with higher band gap

**Low T:** Amorphous silicon is a promising *high* index material (for long  $\lambda$ ) below 50K, with low losses (absorption still needs improvement); losses not yet known at 123K

- What about the *low* index material?  
Silica  $\approx 1 \times 10^{-4}$  at 123K – ok;  $5 \times 10^{-4}$  at 10 K – pretty high (although there is a T in the  $S(f,T)$  so maybe good enough)
  - In many cases annealing improves high temperature loss, but worsens low temperature loss. Strengthening some bonds, weakening others.
  - **Perhaps vapor deposition at elevated temperatures can access an ideal glass, low loss at low T state in  $\text{SiO}_2$  that is inaccessible by liquid quenching or annealing! (in progress)**
- $\text{Al}_2\text{O}_3$  and SiN are alternative good *low* index candidates with low loss
  - May not be necessary to achieve low TLS; **it may be enough to have low coupling constant  $\gamma$ !** (already have seen this effect in  $a\text{-SiN}_x$ , also low density  $a\text{-Si}$  and  $a\text{-Si:H}$ )