

LASER INTERFEROMETER GRAVITATIONAL WAVE OBSERVATORY  
- LIGO -

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LIGO SCIENTIFIC COLLABORATION

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# 1 Introduction

The inauguration of the Advanced LIGO detectors made clear that there are great opportunities to conduct new science by measuring gravitational waves. On September 14, 2015, when only at about 40 percent of the aLIGO design sensitivity, the aLIGO interferometers made the first detection of gravitational waves, GW150914 [16]. During the first and second Observing runs (O1 and O2), LIGO, in collaboration with the Virgo project, detected 4 more binary black hole mergers [17–20] and the merger of two neutron stars [21]. It is certain that improved detector performance will enable more detections and more accurate measurements of the astrophysical parameters encoded in the waveforms of the merging compact objects.

This Instrument Science White Paper describes the work by the LIGO Scientific Collaboration (LSC) to conduct research and development to improve the current generation of LIGO and GEO interferometers and to develop new concepts, prototypes, components, and to perform modeling for future interferometers. The efforts on LIGO Detector Characterization [22] and LIGO Data Analysis [23] are described in companion White Papers, and the coordination of these efforts is described in the LIGO Program Committee document (TBD).

As with Initial and Enhanced LIGO, the success of Advanced LIGO, near-term upgrades (e.g., A+), and future detectors requires a robust instrument science research program. This research program enables us to lay the foundation of upgrades to the detectors and to respond to the many challenges we discover while commissioning and operating some of the most sensitive scientific instruments ever built. This White Paper represents the current best thinking of the LSC Instrument Science working groups:

- The Advanced Interferometer Configurations Working Group (AIC) including Newtonian Noise and Interferometer Simulations.
- The Quantum Noise Working Group (QNWG).
- The Lasers and Auxiliary Optics Working Group (LAWG).
- The Optics Working Group (OWG).
- The Suspensions and Seismic Isolation Working Group (SWG).
- The Control Systems Working Group (CSWG).

The purpose of this white paper is to provide a survey of the current and planned R&D of these working groups and will be updated yearly based on ongoing R&D and both the experimental and observational results obtained.

Figure 1 presents our current best understanding of the fundamental noise sources that will limit aLIGO in the high power, broadband tuning configuration. The design sensitivity is limited by shot noise at high frequencies, by mirror coating Brownian thermal noise in the mid frequencies, and by a combination of thermal, seismic and quantum radiation pressure noises at low frequencies.

Gravitational Wave Interferometer development spans a number of different timescales and research areas. The most immediate effort is to improve the performance and understanding of the existing interferometric detectors and to drive the effort required to reach the aLIGO design specifications. A somewhat longer term effort is to design, build, and characterize

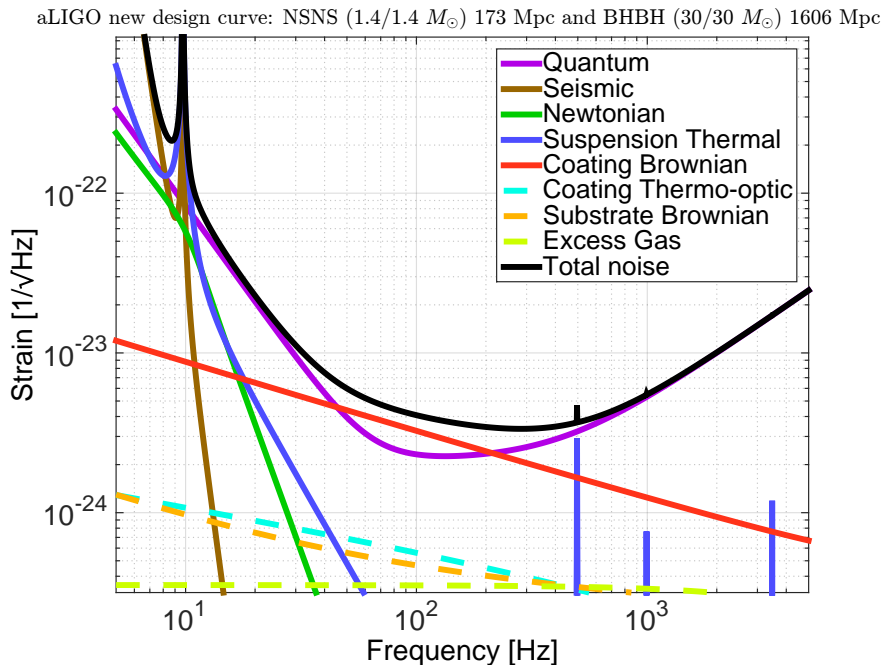


Figure 1: Design Sensitivity Curve for Advanced LIGO. This is the baseline aLIGO Noise Budget (GWINC v2.0). 125 W input power; broadband RSE tuning [1]. There is significant work underway in the Instrument science working groups to achieve this performance within the next two years, and then to move beyond it.

test systems and prototypes for future detector upgrades and entirely new detectors. A final effort, of growing importance as we push the envelope of system performance, is in the physics of materials and condensed matter systems. As this often requires scientists and skills beyond the existing detector groups, the Collaboration must continue to expand these efforts in order to succeed in the larger goal of gravitational wave astrophysics. The research described in this Instrument Science white paper is focused on developing the very best gravitational wave detectors, but as we push the science of precision measurement forward, we are also looking for opportunities to impact other fields of physics.

This White Paper uses a Roadmap to describe the research and development for the instrument over the next 2 decades, 2018-2040, in order to focus effort and aid prioritizing tasks. This Roadmap will be revised and refined yearly.

## 2 Roadmap 2018 - 2040 and Executive Summary

### 2.1 Detector Epochs

The first detection of gravitational waves [16] was a major scientific event and marked the birth of the science of gravitational wave astronomy. The need to probe further out into the universe and thus further backward in time is compelling instrument scientists and engineers to develop and deliver ever more sensitive GW detectors and will do so for generations

to come. Tomorrow’s global network of next-generation detectors will allow us to study novel astrophysics, cosmology and gravitational phenomena as described in the ET Design Study [24] and the LSC technical notes e.g. [25].

The rich physics contained in the gravitational wave signals measured to date suggest great scientific rewards for improved sensitivity across the frequency spectrum. In the current epoch, audio-band upgrades must pursue a general goal to improve the sensitivity across the broadest frequency range including the low frequency cutoff.

The long lead time from conception to operation (Initial LIGO and aLIGO show this to be 15 to 20 years) dictates that planning for new detectors typically begins 2 decades before they become operational. A typical detector cycle includes: Simulation of ideas and concepts; Experimental laboratory work; Conceptual design and prototyping; Proposal writing and submission and engineering; Construction and installation; Commissioning; and Observing.

The research described in this white paper is generally grouped into one of three detector epochs spread over the next 20 years. The first epoch is defined by enhancements to the Advanced LIGO (aLIGO) room temperature detector, in part to bring it to the aLIGO design sensitivity and in part to go beyond the aLIGO design with the A+ upgrade and other subsequent upgrades to the detectors. The second epoch is called LIGO Voyager, and entails a new detector in the existing facilities enabled by cryogenic operation of the optics and new low-mechanical-loss coatings materials. The final epoch requires the construction of new facilities for 3rd generation detectors, and we describe the research needed to enable straw-man designs for detector such as Cosmic Explorer and the Einstein Telescope. These are instruments capable of observing signals originating in the early universe.

The epochs of detector improvement are discussed below, with an emphasis on the critical research required to achieve the desired astrophysical performance. For discussion of the system-level trades used to inform these designs, please see section 3 - Advanced Interferometer Configurations.

## 2.2 Achieving Advanced LIGO Design Sensitivity

The major goal of the LIGO Lab and the LSC Instrument Science teams is to achieve stable operation at the Advanced LIGO design sensitivity by the end of 2020. The LIGO detectors alternate between periods of astrophysical observation and commissioning; the current schedule for the Observing Runs is described in Ref. [26]. At the time of writing this document, the detectors are in a 1 year commissioning period between observing runs O2 and O3. There is another year-long commissioning period scheduled after the completion of the O3 observing run. By the beginning of the fourth Observing run in 2020, the detectors are expected to be operating at the Advanced LIGO design sensitivity, as shown in figure 1.

There are a number of known technical issues which limit the performance of the aLIGO detectors. Moreover, experience has shown that improving the detector sensitivity often reveals unexpected new noise sources. At this time, we know that the following issues will have to be addressed in order to reach aLIGO design sensitivity:

- Fully commission the 70 W high-power laser amplifier.

- Investigate amplifier options to achieve 165W output power after PMC
- Reduce the noise and noise coupling of the laser system.
- Reducing the noise of the current control systems in the 10–25 Hz frequency band.
- Develop improved control techniques to eliminate the control noise in the 10–25 Hz frequency band.
- Improve the angular stability of the detector and develop novel rotation sensors.
- Eliminate the impact of noise from scattered light and excess charge-buildup on the optics.
- Improve the models for interferometer noise coupling.
- Identify and eliminate the (currently) unmodeled noise sources between 20 and 80 Hz.
- Improve the overall duty cycle of the detectors.
- Improve the robust operation of the detector during large environmental disturbances such as high wind and earthquakes.
- Reduce the rate of ‘glitches’ in the detectors.
- Research crackle noise and non-gaussian noise in suspension elements and dissimilar material joints.
- Improve the calibration of the strain channel.
- Control parametric instabilities either through active damping or tuned mass dampers.
- Characterize fully coating thermal noise—understanding the variations in coating mechanical loss measured on different substrates and how this relates to direct thermal noise measurements.
- Improve understanding of the scattering characteristics of the aLIGO coatings, required to assist with fully characterizing scatter loss.

The LSC and the LIGO Lab are working to bring the current Advanced LIGO detectors to the point where they will perform reliably at the Advanced LIGO design sensitivity. Knowledge of the current sensitivity limitations informs the research for next generation instrumentation, and lessons learned from improving the current detectors are already being incorporated into the upgrade plans.

### 2.3 An Upgrade to Advanced LIGO: A+

The LIGO Lab and its international partners have proposed an Advanced LIGO upgrade known as A+. This upgrade is expected to improve the binary neutron star inspiral range by a factor of 1.9 times the aLIGO design specification (to around 325 Mpc) with a binary black hole range greater than 2.5 Gpc, (see Figure 2) at a modest cost. The A+ proposal envisions observations beginning around 2024.

The A+ upgrade proposal includes

- Frequency dependent squeezing using a 300 m filter cavity.
- Design and implementation of a larger diameter beamsplitter.
- Balanced homodyne readout and associated suspended steering mirrors.
- Improved optical coatings including: Mechanical loss reduced by a factor of between 2 and 4.

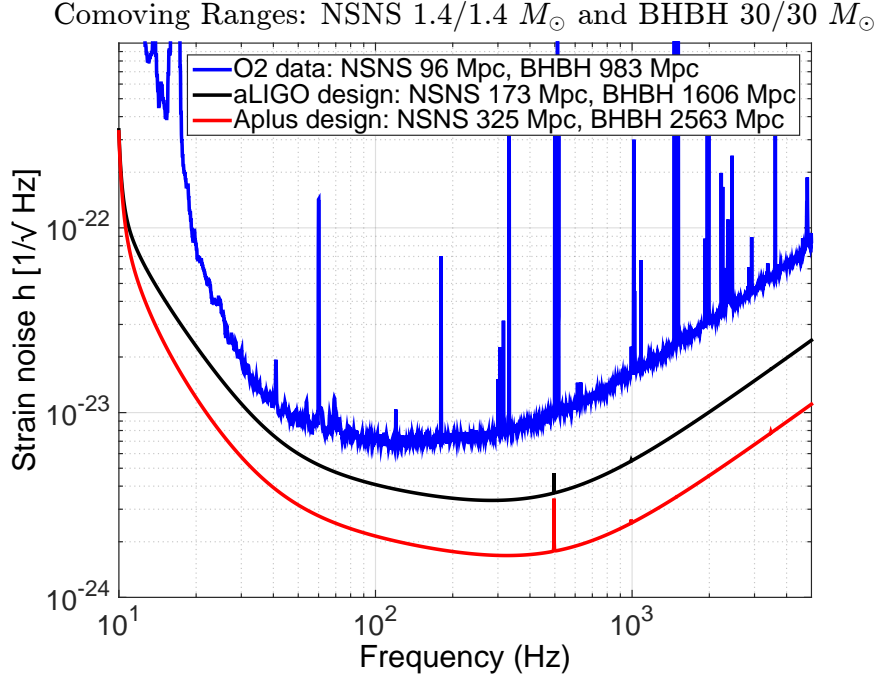


Figure 2: Strain sensitivities from Observation run 2 (O2), Advanced LIGO design and of A+ [27].

- Installation of the new test masses employing fused silica suspension fibers from up-graded pulling and welding systems.

Miller et al [28] have shown that squeezing and coating thermal noise reduction must be combined to achieve maximum sensitivity improvement. Ideally these will both be ready at the same time, but the A+ proposal describes a staged approach if a suitable coating is not yet available.

There is ongoing R&D for A+ which includes:

- Development of improved amorphous coatings to reduce mechanical loss while maintaining optical specified quality at room temperature, aiming for 4 times reduction in mechanical loss.
- Frequency dependent squeezing at 1064 nm.
- Newtonian noise reduction.
- Low noise control of the homodyne readout.
- Reliability testing for higher stress silica fibers.
- Study production of fused silica suspension fibers to ensure frequencies of violin mode of the are sufficiently matched.

The A+ proposal promises an improvement to the detector at a modest cost, but at this point there is no certainty in the funding timeline. Development of backup plans, such as a possible 18 meter filter cavity in the existing buildings, are being examined.

## 2.4 Advanced LIGO Upgrades Beyond A+

Looking past the A+ detector, there are a number of technologies under investigation which could improve detector performance within the current LIGO Observatory facilities while allowing room temperature operation.

- Fused silica suspensions operating at higher fiber stress, able to support heavier test masses (  $\sim 80$ -200 kg).
- Larger fused silica test masses in the range of  $\sim 80$  - 200 kg.
- Improved seismic isolation and suspension control to move the detection band below 10 Hz.
- Charge mitigation techniques including cold gas discharge and possible conductive coatings.
- Newtonian noise reduction.
- Continued development of improved amorphous coatings to further reduce mechanical loss (4 times or greater reduction) at room temperature.
- Development of large crystalline multilayer coatings and the transfer process necessary to apply them to silica substrates.
- Improved substrate surface figure error over the larger area of the larger test masses while managing the residual substrate fixed lens and elastic distortion of the figure error when the mirror is suspended.
- Improved uniformity of coating thickness over the larger area of the test masses.

Successful implementation of these technologies could improve the detector performance beyond that of A+. However, it is likely that we will still be limited by the thermal noise of room temperature optics. To move forward from this point requires a change to the way we currently operate.

## 2.5 LIGO Voyager

LIGO Voyager would be a new detector operating in the existing LIGO observatory facilities. This epoch would involve cryogenic mirrors in the existing LIGO vacuum envelope. A further factor of 3 increase in BNS range (to 1100 Mpc) is envisaged with a low frequency cutoff down to 10 Hz. Simulation and experimentation is underway now so that it could be possible to be operational by end of next decade (*2030 -*), see Figure 5 for the noise budget).

LIGO Voyager will improve sensitivity across the entire LIGO frequency band, 10 Hz to 10 kHz. However, substantial basic research is required to make the system design for Voyager a success because the subsystem improvements are tightly coupled: decisions made about one subsystem likely place requirements on other subsystems. For example, high frequency improvements achieved using increased laser power and low frequency improvements from reduced suspension thermal noise drive the need for low temperature operation. Low temperature will both reduce thermal distortions (optimize thermal conductivity) and lower the thermal noise of the mirrors. It also necessitates a change of the mirror substrate material (to silicon) and may also drive a change to the suspension material. The use of silicon optics drives a change to the laser wavelength to 2000 nm. The need to reduce scatter loss to

maximize the effectiveness of squeezing across the spectrum may also be a driver for longer wavelength operation. In addition to new high-power lasers, the change in wavelength also drives a need for new optical components such as detectors, modulators, and Faraday isolators, all of which must meet LIGO requirements. The development of new optical coatings for low temperature with ultra-low optical and mechanical loss is also required. To achieve the broadband, low-temperature performance shown in figure 5, R&D is needed on (not prioritized):

- Measure low optical absorption and scattering, uniform optical index, and low birefringence across large pieces of bulk silicon.
- Procure, develop, and qualify large test masses.
- Carry out initial cooldown of test masses and suspensions.
- Cool the mirrors in operation with low vibration levels.
- Study additional vibrational noise due to boiling cryogenic fluids.
- Achieve stable control of the mirrors at low temperature.
- High emissivity black coatings for mirror barrels (for radiative cooling).
- Improve controllability of suspensions without violating the load limits of the existing seismic isolation tables.
- Study bond loss for new suspensions: ears, ribbons, etc.
- Characterize of the thermo-mechanical properties of cryogenic materials (thermal expansion, thermal conductivity).
- Research on crackle noise in suspension elements and dissimilar materials joints for cryogenic operation.
- Develop inertial sensors which can operate at cryogenic temperatures.
- Develop passive damping materials which can operate at cryogenic temperatures.
- Develop cantilever springs fabricated from crystalline materials.
- Develop crystalline suspension fibers (ribbons/fibers).
- Longer wavelength laser operating at 200 W.
- Longer wavelength squeezing.
- High power input optics (IO) components (modulators, isolators) at new wavelengths.
- Low-noise, high quantum efficiency photodiodes for longer wavelength.
- High-power low-noise stabilization components at longer wavelength (for spatial, frequency and power stabilization).
- Ion milling level metrology for chosen wavelength.
- Thermal compensation.
- Parametric instability avoidance and mitigation.
- Low phase noise cryogenic interferometer prototype.
- Development of improved optical coatings (amorphous and crystalline) for longer wavelength, optimized for operating at cryogenic temperatures.
- Improved coating thickness uniformity and substrate surface figure error over the larger area of the larger test masses while managing the residual substrate fixed lens.
- Developing silicon substrate super-polishing, surface figuring, bulk scatter and refractive index uniformity.
- Arm length stabilization using a non-harmonically related laser wavelength.



- Contamination control of the low temperature mirror surface.

It is difficult to overemphasize the technical challenge, complexity, and breadth of research and development required to operate gravitational wave detectors at low temperature. It is clear that research must be underway now and that the laser wavelength must be selected soon if technologies are to be ready for large-scale prototyping in roughly 5 years. Some of this work will be leveraged by the progress at Kagra. Even if the start date for Voyager slips past 2030, it is critical that LSC research be supported if these detectors are to be implemented.

The cryogenic operation of these detectors represents the best performance we think can be leveraged from the existing LIGO facilities. To improve the performance beyond this will likely require detectors in the new, larger facilities of next generation observatories.

## 2.6 Next generation of Gravitational wave Detectors and Facilities

Given the long lead-time for new observatories it is timely to make a significant investment in developing and designing the facilities and configurations for third-generation (3G) detectors now. The general expectation is that these instruments must provide a 10 to 30 times better sensitivity than Advanced LIGO across a broad frequency band. Moreover, for increased scientific output and better sky resolution, it is generally agreed that there must be a world-wide network of next generation detectors designed with similar scientific targets and comparable sensitivities, but not necessarily sharing the same detailed design.

The most straightforward way to achieve the target sensitivity is to increase the length of the arms of the new facilities (*e.g.*, 10–40 km). The main limitation to an international network becomes the geographic and political restrictions on the available size of the observatories which are drastically different in Europe, USA, India, Japan, Australia and other potential sites for the NextG detectors. Hence, the paths to the same sensitivity goals will require different approaches and technologies at individual facilities, such as LIGO Cosmic Explorer (CE) in the US and Einstein Telescope (ET) in Europe. It is impossible to predict which set of the currently pursued R&D within the collaboration is best, therefore we present instead a list (not prioritized) of relevant activities and research topics .

- Modeling and simulation of large detectors, underground facilities, and novel readout schemes.
- Studies on the scientific impact vs. cost of various design choices.
- Development of lower cost construction of large vacuum facilities.
- Developing methods of atmospheric Newtonian noise cancellation and/or design of a low-noise infrastructure
- Building quantum back-action (QND) dominated prototype interferometers
- Developing methods of low-loss squeezing injection
- Novel squeezing generation concepts (*e.g.* ponderomotive squeezing)
- New techniques for frequency dependent squeezing without filter cavities.
- Development of 1 kW laser system.

- Enhancing the test masses mechanical response to the gravitational waves using dynamical back-action of light.
- High-frequency sensitivity improvement using negative-dispersion medium in the interferometer
- Developing larger optics ( $\sim 80$  cm ) at room and/or cryogenic temperatures. Technology challenges for manufacturing large optics includes: development of large substrates (silicon/sapphire/silica), figuring of large diameter optics, uniformity of deposition for sputtered coatings and growth/bonding for crystalline coatings.

To the extent possible, the facilities infrastructures will be decoupled from the instruments which will be installed first. The facility should be designed to have a long lifetime, and be capable of housing instruments with sensitivities 10 to 100 times that of aLIGO. Such a facility will, by necessity, be much longer than 4 km. It may be on the surface or it may be underground.

The question of 1–10 Hz sensitivity, in terms of science per dollar, is likely to be a critical factor in the surface vs. underground decision. This is a facility level decision and will impact on the direction of mid-term and long-term R&D. The final designs of these instruments will be decided by the interplay of what can be built (as shown by the instrument science research), what we hope to measure (as illuminated by the ongoing efforts of GWIC to develop the science case for these large detectors), and international agency support to further explore the universe through gravitational waves.

## 2.7 Structure of this White Paper

This vision dictates that

- Work to achieve Design Sensitivity needs to be ready for installation before the beginning of O4.
- Prototyping work for other room temperature upgrades should be underway *soon*.
- LIGO Voyager R&D should be transitioning to the experimental phase in preparation for the design phase beginning around 2021; ideas need to be tested now in order to be mature enough for incorporation into Voyager.
- thinking about the new facility (simulation phase) that will become LIGO Cosmic Explorer needs to begin now.

In the remainder of this paper, the six working groups begin by identifying key issues which are currently limiting the performance of Advanced LIGO and investigations that need to be carried out to overcome these problems. R&D needed for A+, LIGO Voyager and LIGO Cosmic Explorer is then reviewed in detail.

### 3 Advanced Interferometer Configurations

The primary purpose of the AIC group is to coordinate the integration of interferometer systems into coherent wholes, which can then be studied to optimize trade-offs in future detectors. The following sections are organized in terms of the progress of detector technologies presented in the roadmap (section 2), followed by a sections on Newtonian noise mitigation and interferometer sensing and modeling, both of which describe R&D that is relevant to all of the anticipated detectors, such as Voyager and third generation observatories.

#### 3.1 A+

While a few minor upgrades to aLIGO may be possible between major observing runs, such as the installation of squeezed light sources currently taking place in the break between O2 and O3, a dedicated, four-year, upgrade program, called A+, has recently been put forward to the funding agencies, with the target of having A+ operational in about 2023–2024.

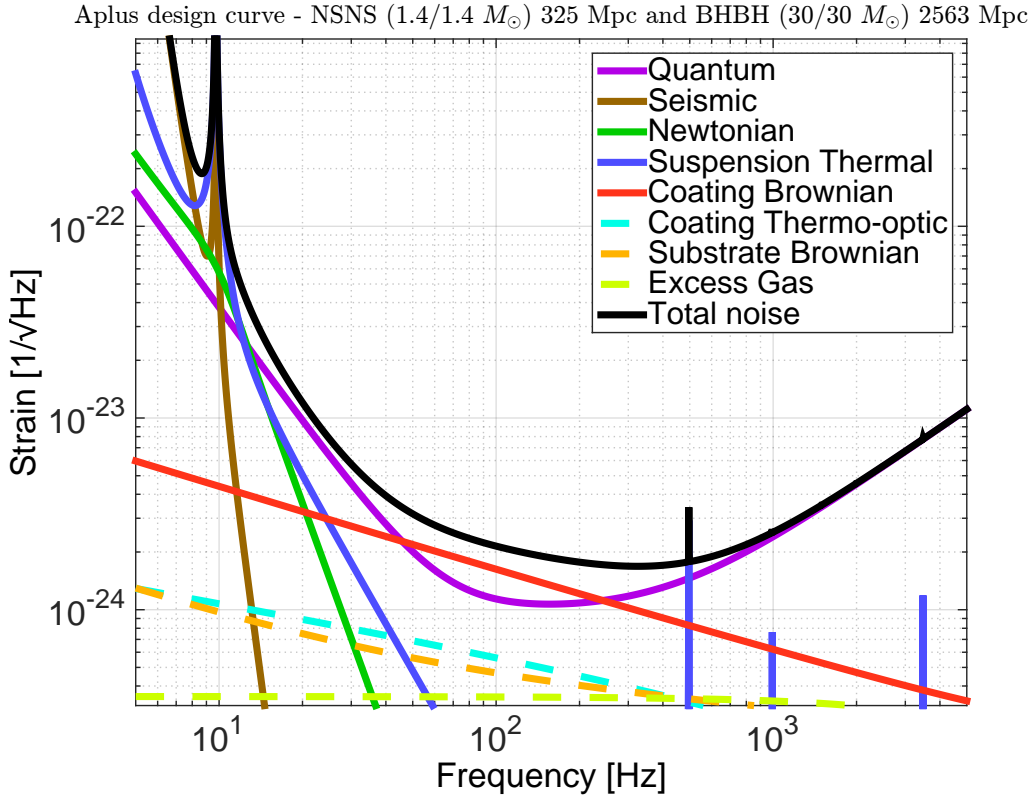


Figure 3: Design sensitivity of the A+ upgrade from [2].

Figure 3 shows the A+ design sensitivity [2] using the parameters as given in Table 3.1. The A+ upgrade targets at an improvement of the two main limiting fundamental noise sources, namely quantum noise and coating Brownian noise.

Once operational, A+ will have a larger reach and therefore will survey a substantially larger

volume of space than the unmodified aLIGO detectors would cover, by a factor of about 4-7 depending on the type of source, in effect collecting many years' science in each observing year. This will come with increased waveform resolution, to discern important physical details of loud event waveforms.

<b>Gwinc Parameter</b>	<b>Value</b>	<b>Comment</b>
ifo.Materials.Coating.Phihigh	$3.6 \times 10^{-4}/4$	tantala mechanical loss
ifo.Materials.Coating.Philown	$5 \times 10^{-5}/4$	silica mechanical loss
ifo.Laser.Power	125W	full power
ifo.Optics.Loss	37.5e-6	75 ppm round-trip
ifo.Optics.BSloss	0.5e-3	
ifo.Optics.PhotoDetectorEfficiency	0.9	Improved readout loss
ifo.Optics.SRM.Transmittance	0.325	SRM transmission
ifo.Optics.SRM.Tunephase	0	SRM tuning
ifo.Optics.Quadrature.dc	$90 \cdot \pi / 180$	Readout phase
ifo.Squeezer.Type	'Freq Dependent'	Squeezing injection
ifo.Squeezer.AmplitudeB	12	SQZ amplitude [dB] (at the OPO)
ifo.Squeezer.SQZAngle	$0 \cdot \pi / 180$	SQZ phase [radians]
ifo.Squeezer.InjectionLoss	0.05	5% loss in SQZ path to IFO
fcParams.L	300m	filter cavity length
fcParams.Lrt	60e-6	round-trip loss in the cavity
fcParams.Te	1e-6	end mirror transmission
<b>Gwinc Output</b>	<b>Value</b>	<b>Comment</b>
Finesse	446	
Power Recycling Factor	43	
Arm power	750 kW	
Power on beam splitter	5.35 kW	
BNS range	325 Mpc	(comoving)
BBH range (30/30)	2.56 Gpc	(comoving, $z = 0.7$ )

Table 1: Gwinc parameters for the A+ design curve [2].

### 3.1.1 Frequency dependent squeezing and balanced homodyne readout

Squeezed states of light [29] have already been used to improve the sensitivity of gravitational-wave interferometers [30,31]. However, any reduction in quantum shot noise at high frequencies is accompanied by a commensurate increase in quantum radiation pressure noise (for frequency independent squeezing).

By reflecting a squeezed beam from a detuned high-finesse optical resonator, known as a filter cavity, one can produce frequency dependent squeezing which can simultaneously reduce shot noise at high frequencies and radiation pressure noise at low frequencies [32, 33], enabling broadband improvements. In order to achieve the required bandwidth of the filter cavity and taking into account that mirror loss, backscattering and phase noise contamination set an upper limit on the allowed optical finesse, for A+ a finesse of 446 and a filter cavity length

of 300 m was chosen.

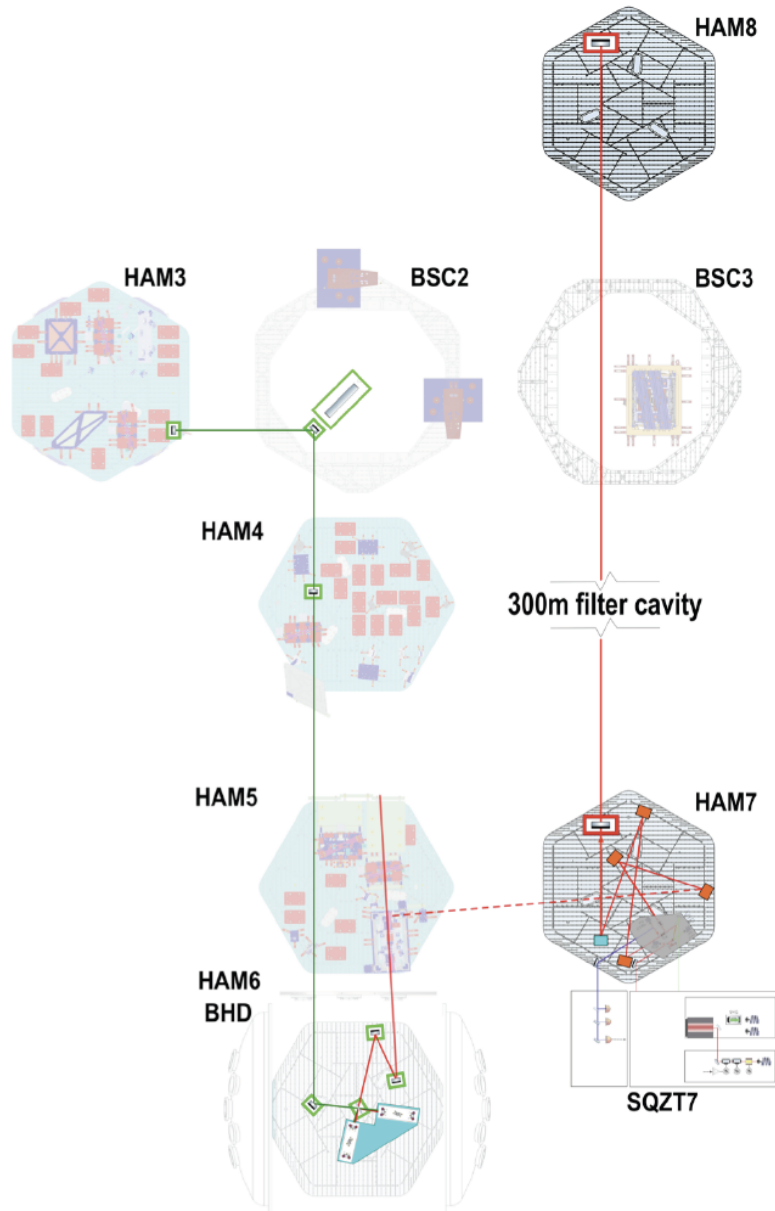


Figure 4: Sketch of optical layout of A+ detailing the filter cavity for the frequency dependent squeezing as well as the beam path relevant for the balanced homodyne readout.

In order to maximise the gain in quantum noise reduction, the A+ upgrades include a number of additional changes aimed to reducing the optical loss in the injection path, the main interferometer and the detection train. First of all the main beam splitter will be replaced by one with a larger free aperture (from currently 370 mm to 450 mm) to reduce clipping losses in the central interferometer from currently about 600 ppm down to about 100 ppm. Additional improvement of optical efficiency will be obtained by application of Faraday isolators with lower loss, improved mode matching between the various optical resonators (e.g. output beam to output mode cleaner etc) and adaptive mode matching elements to be integrated in the output train. Finally, A+ will switch the main readout method from

DC-readout to balanced homodyne detection (BHD). BHD requires a local oscillator beam delivered geometrically separated from the signal beam. As shown in Fig. 4, for A+ the local oscillator beam will be derived from the POP port and then guided via several low-noise suspensions to the detection system, consisting of the separately suspended balanced homodyne beamsplitter and a suspended platform carrying two output mode cleaners and the relevant BHD photo detectors. BHD allows us to eliminate the need for static imbalance in the interferometer arms, and therefore reduces unwanted noise couplings to the GW signal channel. In addition, BHD affords the capability to fine-tune the reference phase in order to optimize squeezed light performance.

### 3.1.2 Coating thermal noise reduction

The optical coatings used in gravitational-wave detectors have extraordinary optical properties: absorption below 1 ppm, scatter losses around 10 ppm and very tightly controlled reflectivities. Conversely, these coatings are mechanically much more lossy than the mirror substrates on which they are deposited. Thus, they constitute the dominant source of thermal noise [34].

Coating research has received considerable attention in the past decade as the use of resonant optical cavities has become widespread in frequency standards, gravitational-wave detectors and other optical precision measurements [35]. Informed by this work, the coatings used in Advanced LIGO are composed of alternating layers of amorphous silica and titania-doped tantalum pentoxide [36,37]. There is significant ongoing effort to develop amorphous coatings with reduced mechanical loss and acceptable optical properties (section 5.1). In addition, there has been a search for new coating materials and technologies for use in future GW detectors.

One potential solution is the use of crystalline coatings. A leading candidate is epitaxial layers of GaAs and AlGaAs, which can be grown on a GaAs wafer and then transferred to a fused-silica substrate. Multilayer AlGaAs Bragg reflectors have been shown to provide at least a factor of three reduction in the amplitude of coating thermal noise [38]. Alternative crystalline materials such as AlGaP/GaP are also under investigation, see [39,40].

While crystalline coatings are promising, they have yet to be demonstrated on a 40 cm-scale mirror. Scaling this technology up presents several technical challenges, both in the manufacturing process and in meeting the extremely stringent surface-figure specification associated with multi-kilometer resonant cavities. As such, **crystalline coatings are not considered practical for use in A+**, though they may be very useful in Voyager or Cosmic Explorer given sufficient and successful technology development.

**Instead the A+ upgrades targets improvements of amorphous coating, ideally providing a mechanical loss of smaller  $10^{-4}$ .** Modeling and tests of amorphous structures suggest that lower-energy microstates of the coatings may be achievable by straightforward modifications of composition and process parameters, in ways that may retain reasonable compatibility with established industrial coating practice. These initiatives are viewed within the theoretical framework of so-called ultrastable glasses (ref), amorphous materials with a very low density of the internal excitation states that contribute to elastic loss. In principle, such stable conditions can be reached by cooling materials slowly from liquid, or

by annealing them after deposition; however, most suitable optical oxides have a competing tendency to crystallize at high temperature, which drastically degrades their optical properties. Alloys can achieve improved intrinsic stability with respect to pure materials; dopants may also hinder crystallization, to allow effective annealing at higher temperature without collateral degradation. Already there are experimental hints that significant improvement can be obtained in tantalum with zirconia doping, a result predicted by computational simulations developed within the LSC (ref). This indicates a highly promising route to the required coating thermal noise performance, within a suitable timescale.

### 3.1.3 Overview of Research and Development for A+

The following list summarizes the R&D expected to be useful for A+, along with more minor improvements which are not directly related to the interferometer design.

- **Coating Thermal Noise Reduction:** Coating thermal noise reduction research is in progress on the aLIGO type coatings and this should be continued. The time scale for A+ is short, and IBS coatings are likely the only viable candidate for producing a mature coating technology. Looking farther into the future, however, Voyager and Cosmic Explorer designs will both benefit from improved coatings (e.g., crystalline coatings or amorphous silicon coatings).
- **Frequency Dependent Squeezing:** Quantum noise reduction R&D is discussed in section 4.2.1. In order to maximize the benefit from squeezing, it is crucial to minimize the losses between the squeezed light source and the gravitational wave photodetector in transmission to the output mode-cleaner (OMC). This requires developing ultra-low loss Faraday isolators, large diodes with high quantum efficiency, and a high throughput OMC. Moreover, at the interface between a squeezed light source and the interferometer, there are problems of alignment control and mode matching. Squeezing drives stringent mode-matching requirements which can only be met by a system capable of on-the-fly optimization. This will require a sensing scheme capable of measuring the mode-matching and feeding back to in-vacuum optics which can change the beam parameters. The mode overlaps of interest are those between the squeezed light source, the filter cavity, the interferometer, and the output mode cleaner (OMC).
- **Balanced Homodyne Detection:** Changing over the main readout system from DC-readout to BHD requires R&D in a variety of areas, including low phase-noise delivery of LO beam to the readout system; integrated arrangement of two output mode cleaners on a pre-isolated platform; low noise in-vacuum readout electronics of the BHD; control strategy for the BHD readout angle etc.
- **Newtonian Noise Subtraction:** All future detectors will benefit from Newtonian noise subtraction using seismometer arrays. While this may not be critical for A+, Newtonian noise coupling predictions vary, as does the seismic environment, making this an important form of risk reduction for both aLIGO and A+.
- **Seismic Isolation Systems:** All three stages of development envisioned here (A+, Voyager and Cosmic Explorer) take advantage of the active Internal Seismic Isolation (ISI) system developed for aLIGO. Although the ISI is designed to support a payload of multiple aLIGO suspensions, detailed mass budgets will be necessary to confirm their fitness. In any case, ISI performance can benefit from improved **vertical inertial**

sensor development, and improved position sensors. These upgrades are **not** intended to decrease in-band seismic noise, but rather to reduce the RMS motion of the ISI platform and minimize noise injection at the suspension resonances (in the 0.2 - 1 Hz band).

- **Auxiliary Optics:** Stray light control will need to be reevaluated for A+ due to its increased sensitivity. The Arm Length Stabilization (ALS) system may need to be redesigned to reduce complexity, possibly by injecting 532 nm lasers from the corner station.
- **Optical Coating Quality:** Study of the three sources of optical loss related to the coatings (coating thickness non-uniformity which leads to figure error, prompt scattering from micro-roughness and scattering from “inclusions” or “defects” in the coatings) will be necessary for A+ to succeed. This is due to the increased sensitivity of the instrument, making scattered light a bigger problem.
- **Charge Mitigation:** Charges on the test-masses are a potentially limiting noise source for aLIGO and future upgrades. Both as risk reduction for aLIGO, and to allow for the improved sensitivity of A+, research related to minimizing charging of (or removing charge from) the test-masses is of interest.
- **Lasers:** The A+ design does not require higher power than aLIGO, but improvements in the PSL design which minimize noise couplings will be very advantageous.

### 3.2 Voyager

The LIGO Scientific Collaboration (LSC) Instrument Science Working Groups held a workshop to study designs for third generation interferometers to be installed in the existing LIGO facilities. Subsequent studies of 3 straw-man designs (known as the Red, Green and Blue designs) showed that all of the designs shared many common requirements, some of which are anticipated in the A+ design described in the previous section. The most promising of the 3 straw-man designs has been dubbed “Voyager”, and adopted as the baseline design for the next major upgrade.

While not intended to exclude other options, a straw-man design for Voyager is used to understand potential benefits and expenses involved in upgrading beyond A+. The Voyager straw-man design mitigates limiting noises of aLIGO by replacing the glass mirrors and suspensions with silicon parts, and operating a detector at the cryogenic temperature of 123 K. The Voyager noise budget and the corresponding sensitivity described in this section are shown in Figure 5. The details of the Voyager design and description of noise sources and design choices is in the working Voyager design document [41].

The main design features of the proposed cryogenic Voyager design are

- **Large cryogenic Silicon mirrors:** In order to obtain broadband improvement of the sensitivity, 120 - 200 kg Silicon mirrors at the operating temperature of 123 K are considered.
- **Silicon cryogenic suspension:** Silicon ribbons (and perhaps blade springs) in the final stage suspension at 123 K are employed.



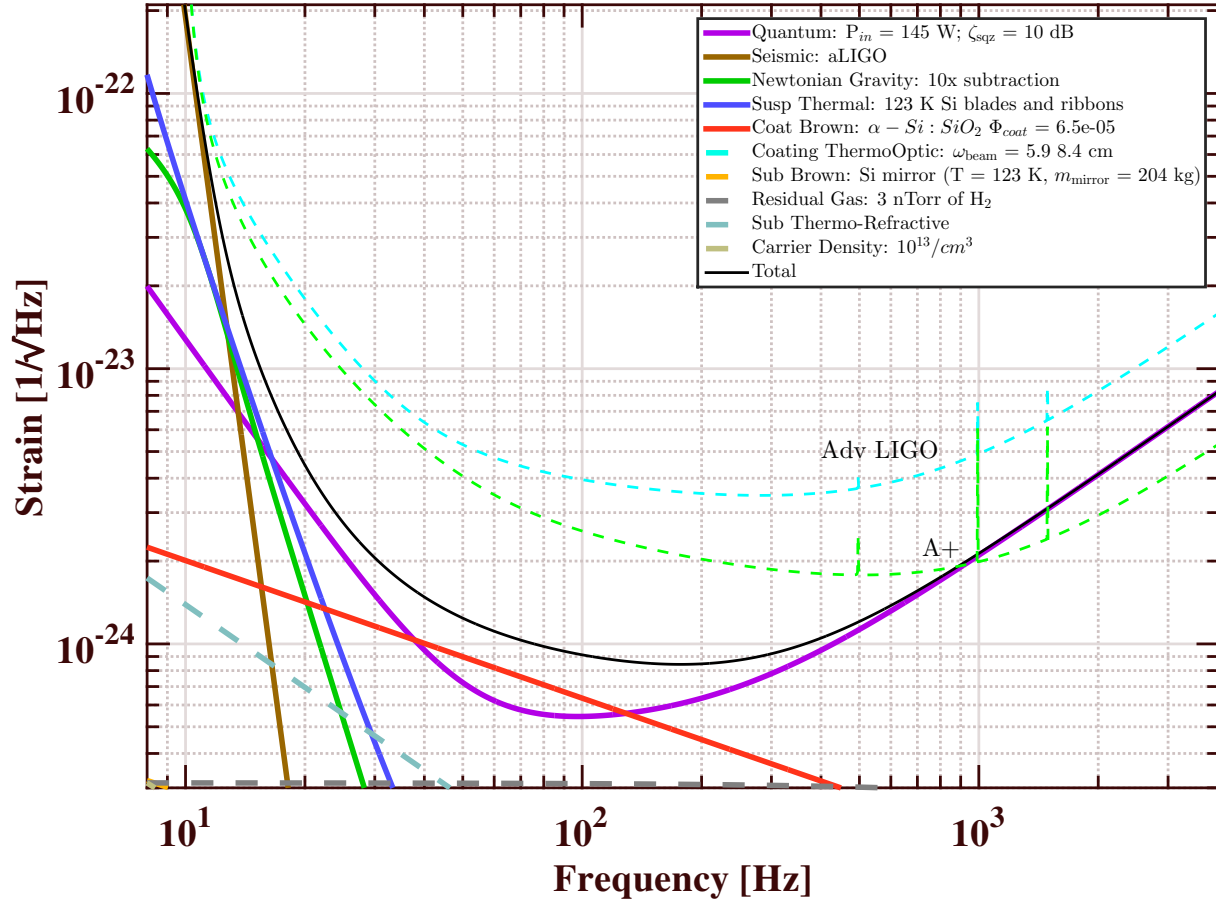


Figure 5: Strain Sensitivity of LIGO Voyager. 200 kg silicon test masses at 123 K, and 3 MW of arm cavity power. See <https://dcc.ligo.org/LIGO-T1400226> for more details.

- **Suspensions:** The Voyager design will require new quadruple suspension systems (SUS) designed to support heavier test masses. In addition to supporting a heavier mass, the redesign efforts should also fix any features found to hinder aLIGO operation (e.g., bounce and roll mode damping, violin mode actuation, gas damping noise, etc.).
- **Newtonian noise:** Newtonian noise subtraction with seismometer arrays is included in the Voyager design, which assumes a factor of 10 suppression.
- **High power 1.5 - 2.0  $\mu\text{m}$  laser:** 1.5 - 2.0  $\mu\text{m}$  wavelength lasers operating at around 200 W are employed. Arm cavity powers will reach  $\sim 3$  MW.
- **Coating Thermal Noise:** The beam spot size is increased by  $\sim 25\%$  relative to the aLIGO size to lower thermal noise while avoiding optical stability issues and the baseline coating is assumed to be amorphous silicon / amorphous silica.
- **Quantum noise:** Squeezed light injection (10 dB) and a 300 m filter cavity for frequency dependent squeeze angle is assumed.

The move to the longer wavelength, motivated by the need for higher thermal conductivity to transport heat away from the optics, along with cryogenic operation offer the wide range

of interesting research areas, listed below. **Note that while the wavelength is specified as 2000 nm for the noise calculations herein, wavelengths of 1500-2100 nm are being considered.** The final choice of wavelength will depend on the material properties of silicon, absorption in the HR coating, the availability of stable high-power lasers, and the quantum efficiency of photodetectors, among other things.

### 3.3 Third Generation GW Observatories

The current facilities, while extraordinary in their capabilities, present significant limitations to gravitational wave astrophysics. In particular, the length of the detectors is well below the optimal value, and the L-shaped vacuum enclosure only allows for the detection of one polarization of gravitational wave signal. A longer, triangular detector would relieve both of these constraints; making an order of magnitude improvement in sensitivity possible, and allowing for both gravitational wave polarizations to be measured with a single detector.

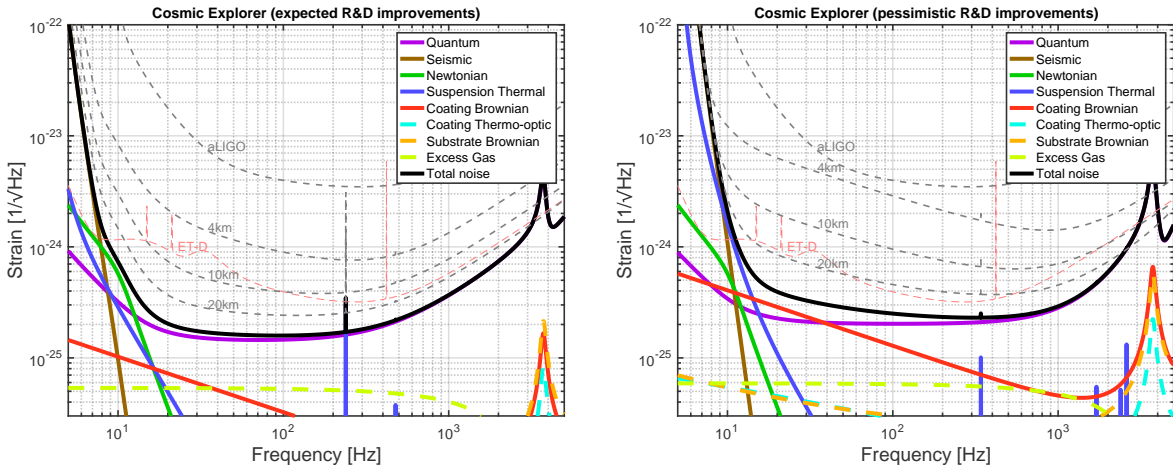


Figure 6: Cosmic Explorer based on Voyager technology (left) has a BNS range of 4.2 Gpc (comoving) and will be able to detect more than 10% of BNS mergers at  $z = 5$ . A more pessimistic CE design (right), based largely on A+ under the very conservative assumption that little improvement is made in coating technology in the next 20 years, serves as a lower bound to the performance of a new facility (BNS range of 3 Gpc and 10% detection at  $z = 2$ ). Binary black hole mergers like GW150914 will be detectable with high SNR out to  $z \sim 20$ . Sensitivity curves for shorter interferometers using similar technology are shown in grey, with coating thermal noise becoming strongly dominant for short detectors. The aLIGO and ET-D sensitivity curves are shown for reference.

Currently, complementary design efforts are being carried out within the LSC for the Einstein Telescope (envisaged to be located in Europe) and Cosmic Explorer (envisaged to be located in the US). The initial design of the Einstein Telescope was developed about a decade ago and is described in detail in the ET design study document [42]. The corner stones of the Einstein Telescope design are: A triangular underground observatory consisting of three detectors of each 10 km arm length to provide full polarisation reconstruction, redundancy and null streams; each detector is comprised of two interferometers, a low-frequency, low power cryogenic interferometer and a high-frequency, high power interferometer.

The Cosmic Explorer design is based on a long, L-shaped observatory located on the earth’s surface [43]. The sensitivity curves presented in figure 6 are computed for a 40 km long detector using A+ technology (CE pessimistic) and Voyager technology (CE). The astrophysical range of such a detector is shown in figure 7. Note that due to the detector’s sensitivity to signals from most of the visible universe, it is necessary to express the range of such an instrument in terms of redshift at the detection horizon.

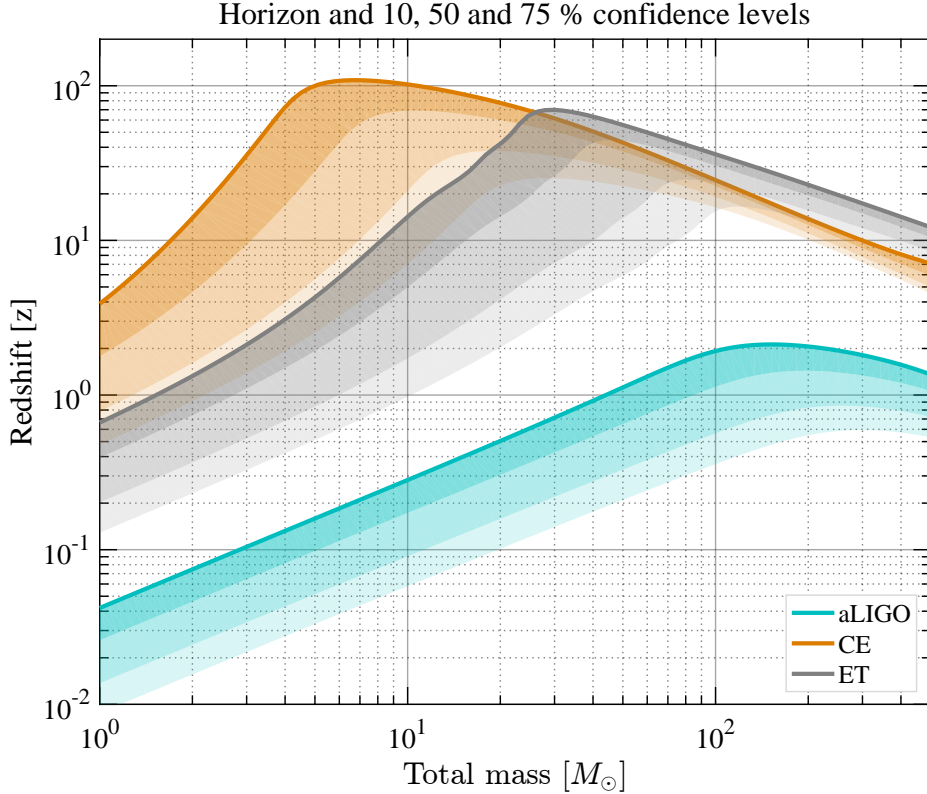


Figure 7: Astrophysical reach of Cosmic Explorer for equal-mass (non-spinning) compact binary inspiral systems. The maximum observable distance is shown as a function of the total intrinsic mass of the system. Barring a large number of primordial black holes, at redshifts larger than  $z \simeq 10$  there will be few sources. Thus a horizon of  $z > 20$  for a given mass should be taken to indicate that essentially *all* compact binary coalescence in the universe will be observable by a network of similar detectors, many with a high signal to noise ratio. Similar curves for the Einstein Telescope are shown for comparison (see figure 6 for relevant sensitivity curves). A Hubble constant of 67.9 km/s/Mpc and a  $\Lambda$ CDM model of expansion was assumed.

The designs used to create the CE sensitivity curves in figure 6 make use of existing technology, or well defined extrapolations of existing technology, as a means of computing a *lower limit* to what can be done in a new facility. This should not be interpreted as a design target: LIGO Cosmic Explorer will make use of the best technology available in time for the final design process. This will certainly incorporate the results of Voyager R&D as well as

other technological developments that occur in the intervening years. The design concept presented here should drive and not limit research into detector technologies and topologies useful in a very long baseline facility.

### 3.3.1 Comparison of current and future detectors

Table 2 shows the comparison of the key parameters of the A+, Voyager and the third generation detectors Cosmic Explorer and the Einstein Telescope.

IFO Cases	aLIGO	A+	Voyager	CE (pess)	CE	ET LF	ET HF
Arm Length [km]	4	4	4	40	40	10	10
Mirror Mass [kg]	40	40	200	320	320	211	200
Mirror Material	Silica	Silica	Silicon	Silica	Silicon	Silicon	Silica
Mirror Temp [K]	295	295	123	295	123	10	290
Sus Fiber	0.6m SiO <sub>2</sub>	0.6m SiO <sub>2</sub>	0.6m Si	1.2m SiO <sub>2</sub>	1.2m Si	2m Si	0.6m SiO <sub>2</sub>
Fiber Type	Fiber	Fiber	Ribbon	Fiber	Ribbon	Fiber	Fiber
Input Power [W]	125	125	140	150	220	3	500
Arm Power [kW]	710	750	3000	1400	2000	18	3000
Wavelength [nm]	1064	1064	2000	1064	1550	1550	1064
NN Suppression	1	1	10	10	10	1	1
Beam Size [cm]	5.5 / 6.2	5.5 / 6.2	5.8 / 8.4	12 / 12	14 / 14	9 / 9	12 / 12
SQZ Factor [dB]	0	6	8	10	10	10	10
F. C. Length [m]	none	300	300	4000	4000	10000	500

Table 2: Baseline parameters for present and future detector configurations

### 3.3.2 R&D for third generation detectors

The R&D required for 3G observatories to suppress fundamental noise sources includes:

- **Large mirrors and relevant metrology :** CE and ET, depending on the exact design flavor, will require silica or silicon mirrors with diameters in the range of 0.5 m to 1.0 m.
- **Silicon cryogenic suspension:** operating at 123 K for CE and 10 K for ET-LF.
- **Suspensions:** need to be able to carry mirrors of several 100 kg for both warm silica (ET-HF) and crogenic silicon (CE and ET-LF). Also need to push suspensions modes to lower frequencies to allow sub 10 Hz sensitivity.
- **Newtonian noise:** Newtonian noise subtraction with seismometer arrays. Decomposition of seismic fields into different wave components. For more detailed description see section 3.6.3 .
- **High power laser:** Laser of various wavelength (1064 nm, 1550 nm and 2000 nm) operating at the several hundred watt level are required to reach arm cavity powers of several MW.

- **Coating Thermal Noise:** The beam spot size will increase relative to the aLIGO size to lower thermal noise. New coating materials are required to provide low mechanical and required optical parameters at cryogenic temperatures (123K for CE and 10K for ET-LF).
- **Quantum noise:** Frequency dependent squeezing with km-scale filter cavities form baseline designs. Alternative QND schemes may provide better sensitivity or reduced costs.

Moreover, there is a huge R&D effort required to reduce a multitude of technical noise source to a level compatible with the envisage 3G sensitivities, in particular in the sub 10 Hz band. This includes amongst others reduction of control and feedback noise, reduction of scattered light, laser noise, photo diode dark noise etc.

Finally, since 3G observatories will by definition be hosted in new facilities, there are various aspects related to site selection which require R&D:

- **Detector Network:** Given certain boundaries what is the best third generation network and how to optimally integrate Advanced + detectors?
- **Observatory Configuration:** Cost benefit analysis of single site with multiple detectors versus several sites with single detectors. This includes research on how to compare in a fair way sensitivity curves from for instance ET (several detectors) versus CE (single detector)
- **Low Frequency Cut-Off:** Analysis of the cost and benefits of extending the sensitivity into the sub-5 Hz range.
- **Integrated Cost Function for various science goals:** How to optimally combine figures of merit of several astrophysical targets (e.g. BNS range, BBH range, stochastic background, high frequency sensitivity for supernova core collapse model separation etc) into a single figure of merit which can guide the design of 3G detectors.

### 3.4 Interferometer Sensing

**High power photo-detection** Unlike previous generations of interferometers, the Advanced LIGO, Voyager, and third generation interferometers will detect substantial amounts of DC photocurrent at signal frequencies as low as 10 Hz. In this region, photodiodes are subject to excess  $1/f$  noise that degrades their performance. Research to characterize, understand, and improve upon photodiodes will be necessary to detect the signals from high power interferometers and at wavelengths from 1 to 2 micron.

In addition to the DC photocurrent, the RF power received by the RF detectors will also increase. Increasing the SNR in the RF detectors will directly reduce the auxiliary controls noise (a limit in nearly all GW interferometers) and so needs to be explored carefully.

**Output Mode-Cleaner** Future interferometers use high finesse mode cleaning cavities to prepare the input laser beam for use in the interferometer and to ensure that the output photodetectors sense only the interferometer’s fundamental spatial mode.

The mode matching between the interferometer and the OMC is particularly critical for future interferometers in which losses will be critical. Output mode matching presents a particular difficulty as it varies with the interferometer thermal state. Current research explores the use of deformable mirrors for adaptive mode matching, together with “Bull’s Eye” wavefront sensing for measurement [44].

**Balanced Homodyne Detection for DARM** To measure high levels of squeezing, allow for a tunable homodyne readout angle, and to avoid a variety of technical noises it would be advantageous to move away from DC readout for the interferometer’s primary output (DARM). Balanced homodyne readout, standard practice in tabletop squeezing experiments, offers many advantages over DC readout (see P1100202, P1300184 and P1500091).

**Wavefront diagnostics** Each of the mirrors in Advanced LIGO will have slightly different absorption characteristics and therefore will react differently when subjected to laser powers projected for Advanced LIGO (see T1100250-v2 [45]) It is useful to develop methods that allow for remote monitoring of the condition of a test mass or beam splitters using optical wavefront sensing methods. On-axis and off-axis Hartmann wavefront sensing have been developed for measuring the absorption-induced wavefront distortion in the test masses and beam splitter. The measured noise limited sensitivity of the Hartmann sensor itself is  $\lambda/15,000$ , and experiments have measured wavefront changes smaller than  $\lambda/3000$  [46]. When applied to off-axis tomographic measurements, the current measured accuracy is  $\lambda/120$ , limited by factors other than the Hartmann sensor itself [47]. Further research is aimed at improving this performance. Avenues for improvement include simple and stable injection schemes for Hartmann probe beams into working interferometers and incoherent probe beam sources with high power and low noise.

In addition to the Hartmann style sensors, we may use phase cameras (essentially multi-pixel RF Wavefront Sensors). These could be used to implement real-time wavefront correction. Primitive phase cameras have been used in iLIGO and iVirgo, but there were problems due to the scanning induced backscatter. Future phase cameras should allow for wavefront sensing of individual sidebands simultaneously without any moving parts.

### 3.5 Interferometer modeling / simulations

Interferometric gravitational-wave detectors are sufficiently complicated optical systems that detailed modeling is required for design and performance studies. The detector behaviour cannot be modelled with commercially available optical simulations because our requirements are very different from conventional systems. For example, the length scales involved in interferometric GW detection span 22 orders of magnitude. A variety of simulation tools has been developed by the GW community, successfully addressing the various problems encountered in interferometric GW detectors.

The planned upgrades to the detectors include either new technologies or envisage pushing detector parameters closer to their limits. In both cases the detector design and commissioning face new challenges, as previously used assumptions and simplifications can no longer be used. Simulation tools and the understanding of interferometer modelling have to be advanced ahead of time in order to be able to provide the essential design and commissioning support.

Simulations which use the FFT for propagating paraxial beams such as SIS [48] or OSCAR [49] can be used to make detailed predictions about the impact of optical phase errors (surface roughness, phase distortions, etc.) and finite aperture sizes. FFT simulations are, however, too slow to simulate interferometer dynamics, and are therefore only used to find steady-state solutions.

For understanding interferometer dynamics over a wide range of conditions which do not allow for linearization (e.g., not only at the operating point), time-domain simulations are the most appropriate tools. This type of simulation is important for lock acquisition studies and existing tools fill the need for most optical configurations [50, 51]. However, like FFT simulations, despite considerable optimization effort, time-domain simulations are slow and able to simulate only the lowest transverse-spatial modes.

For control system development, where we can assume that we are at a stable operating point and linearize around that point, frequency domain simulations such as Optickle [52] or Finesse [53, 54] prove an invaluable tool. Advanced LIGO control system development has depended on the Optickle simulation engine, which has been packaged for LSC design in Looptickle and later Lentickle, and for ASC design as Pickle. In addition to control system development, frequency domain simulation tools are also used to compute limitations to interferometer sensitivity due to fundamental noise sources (e.g., quantum noise) and technical noises (e.g., laser noise and auxiliary length degree of freedom control noise). Future tools should provide a simple and effective way to combine optical, mechanical and electronic systems to estimate control noise in the presence of multiple cross couplings.

The commissioning of Advanced LIGO, as well as the design of further upgrades have shown that the combination of frequency domain simulations with paraxial beam propagation has become crucial to understand interferometer limitations. Geometric instabilities and thermal beam distortions, both due to higher laser power, can dominate the interferometer behaviour. Modelling tasks related to alignment sensing and control, parametric instabilities, or simply the effect of mode mismatch or beam clipping on the control systems and ultimately on the gravitational wave signal have taken center stage. Simulation tools such as Finesse and MIST [55] use higher-order mode expansion to describe the laser beam properties and can combine the speed of frequency domain tools with the power of paraxial beam propagation. Such modal-modelling tools have become key tools in the commissioning of Advanced LIGO to date and will be used to transfer and apply aLIGO knowledge into the upgrade design process. Further development is ongoing to improve their performance and to add more advanced features.

The configuration level simulation, discussed in section 4.4, operates at a higher level in that the response of a given optical configuration is symbolically computed and parameterized [56]. These symbolic computations can be compared with the more detailed numerical results given by frequency domain codes. While inappropriate for detailed simulation of the optical

plant, this approach is very effective for an initial exploration of a parameter space for a variety of optical configurations. Further development of the GWINC configuration level simulation to accommodate the variety of optical configurations is currently under way for future detectors and will aid in downselection and optimization.

To date these simulations have focused on the core of the interferometer, the optics, and have largely neglected the surrounding mechanical systems. A missing piece in detector simulation is a comprehensive mechanical simulation tool for the vibration isolation and suspension design which includes the capability to handle a variety of mechanical systems and the ability to compute of thermal noise for any given configuration.

In addition, development of the various simulation tools, future detector design and study will benefit greatly from a dedicated effort to organize and document these tools. Experience has shown that smaller, dedicated tools have been more effective than those tools trying to combine everything into one system. However, most tools come only with rudimentary or outdated documentation, and code reviews or formal testing of simulation results are not common practice. In particular the main and established software tools should provide:

- An active maintainer who is responsible for the current code base and who can be identified and reached by any users of the software.
- A single webpage hosting the master version of the tool or code.
- Documentation about the implemented models and their limitations, including descriptions of mathematical algorithms, or parameter sets used.
- Training material, such as a set of examples and tutorials for new users, especially graduate students.

The effectiveness of simulation tools for the development of future detectors will be enhanced with some more high level organisation to coordinate the above requirements.

Simulation tools are strongly defined through the context in which they are used. Many of the priorities for the research and development directly translate into a priority task for modelling work or a required effort in developing new capabilities for simulation tools. Below we list a subset modelling challenges.

**Commissioning Advanced LIGO to reach design sensitivity** will continue for several years, with modelling tools being essential for diagnostics and developing mitigation strategies. Advanced LIGO has seen a rapid initial commissioning phase leading to the first detection. As expected, progress has now slowed because the origins of excess noise become harder to identify or eradicate closer to the final design sensitivity.

Advanced LIGO and future detectors will operate at higher light power so that thermal distortions of the optics and the change of the mirror suspension response due to radiation pressure will play a dominant role. In particular the transfer functions of any displacement signal into the detection ports will show the optical spring effects and higher-order mode resonances. Thus to model noise coupling of auxiliary degrees of freedom or auxiliary optics



system into the gravitational wave channel we require models that include thermal effects and radiation pressure at the same time. For the rapid development of the control systems and noise mitigation strategies during the detector commissioning it is desirable that such simulation software is a fast, flexible and easy to use tool. Specific priority tasks include:

- Alignment sensing and control, modelling of mode sensing schemes, possibly using alternative sensors, in particular to improve noise around soft mode. Adding additional sensors to simulation tools, matching the models to experimental results.
- Noise couplings at low frequencies: modelling of auxiliary optics and thermal distortions to mitigate coupling into longitudinal signals at low frequencies, modelling of beam jitter and beam size fluctuations. Including a more detailed mechanical coupling and feedback in optical modelling tools.
- Noise Budget Software: a simple tool (preferably matlab or python) to estimate the current noise budget of the LIGO interferometers based on current data and parameters. Software to quickly and easily make measurements of noise budget traces (parent spectrum and non-linear coupling functions).
- High-power operation: modelling of parametric instabilities and thermal lensing to improve performance of mitigation systems, such as active wavefront control systems. Provide and test convenience functions for beam tracing and mode couplings.
- Knowledge transfer and training: providing documentation of modelling result, easy to use parameter files for the LIGO detectors and a forum in a simulation group will be important for recording the knowledge gained from commissioning and operating the first advanced interferometer and to train young scientists, towards the design of upgrades and new detectors.

**The A+ upgrade** introduces significant changes to the output optics, adding a balanced homodyne detector and frequency dependent squeezing. The design for these systems is under way. During the A+ upgrade a down-selection of coating options is required with the initial coating design being driven by models. The necessary modelling tools are available but require the development indicated above. Current priorities for this phase include:

- Frequency dependent squeezing, modelling of non-linear optics in the squeezer and defects in the filter cavity and injection path to optimise system for maximum effective squeezing. Modelling of control signals for filter cavity and of injected squeezed light, especially alignment sensing and control.
- Balanced homodyne detector, in particular coupling of signals and noise due to imperfections in core optics and auxiliary output optics. Modelling of control for auxiliary optics, the OMCs and the local oscillator phase.
- Coating thermal noise, combine various models into unified and accessible coating simulator, provide guidance for coating down selection.

**Voyager, and possible further upgrades of current facilities** . Each improvement in detector sensitivity will uncover new unwanted noise coupling mechanisms that typically require post-hoc mitigation strategies. This implies that for the designs of future detectors, we must model and study potential designs at greater depths than previous detectors. Voyager will introduce significant changes, such as cryogenic suspensions and a different laser wavelength. The main challenge will be the operation at very high circulating power (several MW). For the design of medium term upgrades such as Voyager a more consistent modelling of control noises is required to reduce the risk of low-frequency excess noise already at the design stage.

- High-power operation at low optical loss. Small asymmetries in absorption can increase the effective optical loss for squeezed states. Modelling of squeezed light in higher-order modes, improved thermal compensation systems, improved arm and mode matching techniques.
- Parametric instabilities control, modelling of advanced mode dampers, development of a control scheme that allows tracking and selective damping of a large number of modes, requires more detailed modelling to predict individual modes.
- Scattered light control, modelling of backscatter of detection optics and internal interferometer scatter. Include injection of noise with specific coherence into interferometer modelling tools.
- Control design: better models of control schemes can be achieved by developing more effective tools for the analysis of in-loop cross coupling of a mixed mechanical, optical and electronic system. It is expected that modern control schemes as mention in ?? will be used in a subset of systems at this time.

**Third generation designs** will introduce new facilities and therefore can potentially make use of very different optical configurations (very long arms, speed meters, very long auxiliary cavities). Quantum noise reduction presents a challenge that will most likely have the strongest impact on the interferometer design. Newtonian noise reduction will be required. At the moment the main priority for interferometer modelling is to capture the knowledge from the current design and commissioning work to inform 3G designs. Other priority modelling task for this stage include:

- Advanced quantum noise schemes, development of a robust 'fundamental' quantum limit. Modelling of quantum correlations through complex MIMO (multiple in, multiple out) systems.
- Modelling of non-linear optical elements, such as crystals (squeezing), active optomechanical elements (unstable filters) etc. Development and implementation of realistic linearised couplings for these elements into optical models.
- Modelling of long-baseline and high-cavity finesse operation, including geometric consideration, control schemes and thermal effects.

- Study of optical configurations which rely strongly on polarisation schemes, requires the addition of light polarisation to interferometer models.
- Newtonian noise reduction, advanced modelling of local sensing (6D) and global control strategies, a more advanced implementation of mechanical systems and seismic and Newtonian noise coupling in interferometer models. Simulation of Newtonian noise based on ground noise measurements.
- Investigation of alternative beam shapes for thermal noise reduction, modelling of auxiliary optical systems to study feasibility

### 3.6 Newtonian Noise

Fluctuations of the gravity field due to terrestrial sources cause random displacement of test masses. This is generally known as Newtonian noise (NN), and it is predicted to limit the sensitivity of Advanced LIGO detectors between about 10 Hz and 30 Hz once their target sensitivity is reached [57]. According to current models, it is the contribution from seismic surface fields that will dominate at the LIGO sites [58]. Specifically, one can neglect atmospheric NN, NN from seismic body waves, and NN from vibrating structures. The main activities today focus on the development of a NN cancellation system as possible upgrade of the Advanced LIGO detectors based on monitoring of the seismic field by a seismic surface array [59].

Since NN increases steeply towards lower frequencies, and since its mitigation is a major challenge, NN must be considered early in the design of third-generation detectors, which will potentially extend the detection band to frequencies below 10 Hz. The outcome of these early analyses might strongly influence the detector design, i.e., they will inform decisions whether the detector should be built underground or above surface, or how aggressively the community should push the development of technologies to gain low-frequency sensitivity. Current efforts include the modeling of NN in underground environments, and characterization of underground seismic fields at the former Homestake mine [60]. Another major concern is that atmospheric NN might become a limiting noise source, which could potentially set an ultimate sensitivity limit especially to surface detectors [61].

#### 3.6.1 A+

**Status quo** A first comprehensive estimate of various NN contributions at the LIGO sites showed that seismic (ground vibration) NN strongly dominates (by about a factor 10 and more) over other NN contributions such as sound NN, and vibration of infrastructure [58]. These studies justify the focus of current activities on seismic NN characterization and cancellation. Only if NN is to be canceled by about a factor 10 (or more), more careful studies of sub-dominant NN contributions will be required.

Two array measurements were carried out at the LIGO Hanford site. The first array consisted of 44 Wilcoxon 731-207 sensors that were deployed in 2012 at the EY end station and took data for about a year [59]. The second array consisting of 30 L-4C sensors together with a tiltmeter from the Washington group was deployed in August 2016 at the corner station and also took data for about a year. Analysis of the array data showed that seismic fields are highly homogeneous, stationary in the NN band, and the dominant component is produced by local sources. This suggests that it will be relatively easy to cancel NN at the Hanford site, and that it is straight-forward to compute near-optimal array configurations in advance.

Further evidence of the efficiency of a future cancellation system came from noise-cancellation tests using a compact tiltmeter as target and seismometers as reference channels (to be published). It had been shown previously that ground tilt under a test mass is fully coherent with acceleration NN under certain ideal conditions, and that it therefore serves as an ideal proxy of the actual NN and that a tiltmeter can be used to probe NN cancellation well before it is seen in GW detectors [62]. These tests were carried out with the result that the tilt signal fully canceled in tiltmeter data. Performance limitation came from the tiltmeter

instrumental noise.

**Summary of required R&D activities** The development of a NN cancellation system as possible upgrade of Advanced LIGO requires activities in the following categories:

- Data characterization in the NN band
- Deployment of seismic arrays
- Static Wiener filters
- Optimization of array configurations
- Implementation in existing data-acquisition systems

**Data characterization in the NN band** It is important to understand other noise contributions in the NN band to be able to predict and later analyze the performance of the NN cancellation system. Such studies should include long-term correlations between seismic sensors and the GW channel, search for coincident transients in the NN band between seismic channels and GW channel, and also the investigation of other than gravitational couplings between seismic data and the GW channel. Some of the required tools can be derived from existing software developed in various LSC groups such as STAMP-PEM, KleineWelle, h-veto, etc.

**Deployment of seismic arrays** Since one of the results of the array measurements is that the dominant seismic sources in the NN band are all part of the site infrastructure, one might argue that properties of the seismic fields at other stations can be inferred from the existing array data. While this might be true for the remaining Hanford end station, differences in geology between the Livingston and Hanford sites are substantial. Therefore, the recommendation is that at least one array measurement should be carried out at the Livingston site, preferably at the corner station, using a similar number of sensors as used at Hanford, ideally including a tiltmeter, and with data being recorded for at least two weeks.

**Static Wiener filters** The simplest filter type that can be expected to achieve some cancellation of NN is the static Wiener filter [63]. It takes the data from a seismic array as input and maps them linearly to a single output, which represents the best linear estimate of the NN associated with the seismic field.

One needs to keep in mind that the optimization of the filter in this context does not include the optimal positioning of seismic sensors, on which the filter performance will strongly depend [57]. This problem is addressed in a separate task.

Preliminary results obtained from the Hanford array measurements suggest that better cancellation performance can be achieved if the Wiener filter is recalculated once per day, but the effect is minor (to be published). Such findings are important for the implementation of a noise filter since more options exist to implement a static filter.

The standard representation of a Wiener filter is a finite-impulse response (FIR) filter, but alternatives should be investigated. One might consider other static filters, Kalman filters, or adaptive filters. Even neural networks might be used to represent a filter, which has the potential advantage that a non-linear filter can in principle deal with more complicated structure in the data such as repetitive short transients.

**Optimization of array configurations** The optimization of array configurations turns out to be the most challenging task for the development of a cancellation system. This will be true for A+, Voyager, and third-generation detectors. An array can be completely ineffective to cancel NN if the sensor positions relative to the test masses are chosen badly. Optimization can be based on models of the seismic field, which are necessarily simple. In this case, algorithms have been developed to calculate optimal array configurations.

The problem is when the seismic field has certain features that are difficult to include in a physical model. Above all, seismic fields with significant inhomogeneity can typically be modeled only with computationally expensive finite-element simulations, which is not compatible with finding optimal configurations of arrays with many sensors being by itself a computationally expensive operation.

The easiest way to base the optimization on measurements is to pick the most effective seismometers from an array that has more sensors than required for cancellation and discard the rest. This has been successfully tested with the Hanford arrays. The open problem is to make use of array data to predict where the optimal sensor locations are, which do not have to coincide with the position of any of the sensors used for this analysis. A solution to this problem might lead to improvements of a cancellation system for A+, and Voyager, but it is absolutely essential for realizing a cancellation system for underground detectors, or surface detectors with very ambitious targets for NN suppression.

It should be emphasized that the two array measurements at Hanford suggest that optimization can in fact be based on relatively simple, homogeneous models of the seismic field, which greatly simplifies the problem.

**Implementation in existing data-acquisition systems** This task concerns the strategy for developing a cancellation system. A cancellation system can be implemented by (a) using stored data, (b) using static filters that subtract in real time before data are stored, e.g., to improve sensitivity of low-latency GW searches, and (c) using filters that are adaptive or regularly updated. It seems feasible to start with either (a) or (b), and develop the system towards (c) as upgrade. Each option requires the development of software that should work autonomously even for option (a). An additional issue about option (c) is that it might require significant computational resources, which are not available in the existing data-acquisition system.

### 3.6.2 Voyager

**Status quo** Little work has been done so far that is specifically relevant to Voyager. One of the main tasks will be to understand the composition of the seismic field in detail in terms of

surface and body waves. This can be done with existing Hanford array data. Work that has already started concerns the filter design. Neural networks were trained with reinforcement-learning algorithms with the goal to substitute the static Wiener filters. Such nonlinear filters might have fundamental advantages.

**Summary of required R&D activities** The development of a NN cancellation system for Voyager would result from continuous enhancements of the A+ cancellation system. In order to achieve a factor 10 suppression of NN, the following additional tasks need to be included:

- Modeling of NN from seismic body waves and its cancellation
- Potentially, deployment of borehole seismometers
- Optional: modification of site infrastructure
- Optional: advanced filters (nonlinear, adaptive, ...)

**Modeling of NN from seismic body waves and its cancellation** It is conceivable that the contribution of body waves to seismic NN can be significant if noise suppression factors of 5 or higher are to be achieved. In order to be prepared for this case, models have to be refined that combine NN from surface and body waves. The models must also provide methods to investigate possible array configurations to cancel NN from body waves if necessary. This task should also include an estimation of the effect of seismic scattering on NN and its cancellation, and the impact of local, near-surface geology.

**Potentially, deployment of borehole seismometers** A possible result of the previous task is that a three-dimensional seismic array is required to cancel seismic NN in Voyager. This means that borehole seismometers need to be deployed. Since it is probably too costly to deploy borehole arrays for site characterization, the optimal three-dimensional configuration must be estimated from models and surface-array data.

**Optional: modification of site infrastructure** A result of the previous array measurements at Hanford is that the dominant seismic sources in the NN band all form part of the site infrastructure. This is above all the air-handler fans at the end stations, and a variety of sources at the corner station. Any decrease of the seismic perturbation caused by these systems would result in a similar decrease of NN. It is therefore reasonable to consider modifications of the site infrastructure as part of the NN mitigation effort.

One approach would be to relocate the dominant seismic sources to increase their distance to the test masses, or to redesign their supports so that less energy is coupled into the seismic field. It was also suggested to dig recesses around the test masses that would remove most of the mass bearing the density perturbations of the seismic field. This could greatly reduce NN from seismic surface waves. These schemes are more costly than NN cancellation, but they would guarantee a certain NN reduction, and they are complementary to NN cancellation.

**Optional: advanced filters (nonlinear, adaptive, ...)** Many potential alternatives to static Wiener filters exist. Wiener filters can be continuously recalculated, one can implement Kalman filters or generic, (quasi) linear adaptive filters, or even nonlinear filters for example based on neural networks. Enhanced filter designs might provide optimal noise cancellation even if two-point correlations of the seismic field change slowly, and some of the design options might even be able to deal with sudden changes of the field associated with seismic transients. It is therefore of interest to investigate alternative filter designs.

### 3.6.3 Third Generation

**Status quo** Two third-generation projects are currently being pursued. First, finite-element simulations are carried out to investigate the impact of stratification and surface topography on NN and its cancellation. Second, an underground array was deployed and decommissioned at the the Sanford Underground Research Facility (the former Homestake mine). Analysis of the Homestake array data is expected to provide important insight into the composition of underground seismic fields. The Homestake array was taking data for about a year, and was mostly composed of STS-2 seismometers with 15 underground and 9 surface stations.

**Summary of required R&D activities** Research on NN for third-generation detectors has various new aspects compared to the A+ and Voyager cases. First, it can influence site selection. Avoiding a site with high seismic noise is the best way to mitigate NN. This also concerns the question whether a third-generation detector should be constructed underground. Furthermore, atmospheric NN needs to be considered. In general, greater attention needs to be paid to all tasks since an overall mitigation goal of more than a factor 100 NN suppression relative to the predicted level at the LIGO sites might be requested. In addition to all the tasks listed in the A+ and Voyager sections, NN research for the third generation includes:

- Modeling of atmospheric NN and its cancellation
- Modeling of NN from vibrating structures
- Identification of sites with low NN and/or beneficial to NN cancellation
- Study of underground seismic fields
- Algorithms to analyze data from three-dimensional seismic arrays
- Potentially, development of technology to cancel atmospheric NN, e.g., LIDAR based
- Optional: designing low-noise site infrastructure

**Modeling of atmospheric NN and its cancellation** It is possible that atmospheric NN will become a limiting noise source of third-generation detectors depending on how much the observation band will be extended towards lower frequencies. In order to develop a cancellation system of atmospheric NN, it is important to improve our current models.



Especially, hydrodynamic simulations are required that can model turbulent air flow and atmospheric temperature fields. This task must also include a simulation of atmospheric NN cancellation based on yet-to-be-developed technologies.

**Modeling of NN from vibrating structures** Newtonian noise from vibrating structures such as vacuum tanks or parts of the suspension system can typically be neglected, but this has only been shown to be true for Advanced LIGO (and this conclusion probably extends to Voyager). More careful analyses have to be done for third-generation detectors, where marginal limitation by NN from structural vibrations is conceivable. If identified as a relevant NN contribution, then cancellation of this component will be relatively straight-forward since a few sensors should suffice to monitor the vibration for the purpose of NN cancellation. While rough estimates are easy to obtain, a more detailed analysis requires finite-element simulations of the vibrating structures.

**Identification of sites with low NN and/or beneficial to NN cancellation** Site characterization forms an important part of the development process towards the third generation. Candidate sites have to be evaluated based on many criteria, including an estimation of NN and the impact of local geology and topography on NN cancellation. With respect to NN, the site should have lowest possible seismic noise in the NN band, and be supported by a homogeneous ground with flat topography. The last two conditions make sure that seismic scattering can be neglected, which might otherwise cause NN cancellation to be unfeasible. Characterization of underground sites can be particularly time consuming and costly, and requires special, robust equipment for data acquisition.

**Study of underground seismic fields** There have been very few experiments with underground seismic arrays, especially with installations of high-quality, broadband seismometers. It is therefore important to gain experience with underground arrays, and to learn how to characterize an underground site and how to use the data to calculate accurate NN models.

The array design needs to be based on prior estimates of seismic speeds and depends on the frequency band where data are to be analyzed. As a rule of thumb, the array needs to have a spacing that does not exceed half a wavelength of the shortest waves of interest, and it should ideally have a diameter larger than the length of the longest waves of interest.

**Algorithms to analyze data from three-dimensional seismic arrays** The analysis of data from a three-dimensional array is a complex task. The most important goal is to understand the average composition of the seismic field in terms of shear waves, compressional waves and surface waves so that the performance of a NN cancellation system can be predicted. The analysis can be further complicated by seismic scattering, and local seismic sources. Other goals include the estimation of seismic speeds, the measurement of propagation directions of waves in the field, identification of seismic sources, and cycles due to changes in seasonal and anthropogenic activity.

**Potentially, development of technology to cancel atmospheric NN, e.g., LIDAR based** No scheme has been proposed so far shown to be effective to cancel NN from the atmosphere, which means that atmospheric NN should currently be considered as ultimate, low-frequency sensitivity limit. It is known that a surface array of microphones is ineffective to cancel NN from infrasound. A new idea is to monitor the atmosphere by a LIDAR system. This technique is promising since it would give information about the two main contributions to atmospheric NN: infrasound NN and NN from advected atmospheric temperature fields. First estimates showed that current LIDAR systems do not have the required sensitivity at least to monitor infrasound, and so it needs to be investigated how much the sensitivity of these systems can be improved. It is possible though that current systems are already able to monitor temperature perturbations with sufficient accuracy. It is also necessary to simulate NN cancellation with LIDAR to learn how practical constraints such as spatial resolution and range affect cancellation performance.

**Optional: designing low-noise site infrastructure** Constructing a new site offers the possibility to improve its design compared to the existing sites with the goal to minimize NN. The goal is to keep seismic noise introduced by the site infrastructure at a low level. Seismic noise can be introduced in the form of vibrating machines such as air-handler fans, or by coupling of wind turbulence with the building walls. Seismic sources should generally be kept at greatest possible distance from the test masses, which not only reduces seismic noise at the test masses, but also facilitates cancellation of the associated NN. Wind-generated seismic noise is not an issue for the advanced detectors, but it should be investigated again for a third-generation detector if built at the surface.

## 4 Quantum Noise Reduction

### 4.1 Introduction

The 2<sup>nd</sup> generation detectors, which are now being pushed towards design sensitivity (Advanced LIGO) or being assembled (Advanced Virgo, and KAGRA), are expected to be limited by quantum noise over nearly the entire GW band (10 - 10000 Hz). By quantum noise, we refer to the quantum uncertainty of the electro-magnetic field at Fourier (side-band) frequencies in the GW detection band that beats with the monochromatic laser field (carrier field) to produce shot noise (photon counting noise, quantum measurement noise) and radiation pressure noise (quantum back-action noise).

To upgrade those detectors it is essential to reduce the quantum noise. In this section techniques are discussed that go beyond scaling-up the carrier light power for reducing the (signal-normalized) shot noise and the test masses for reducing the (signal-normalized) radiation pressure noise. Both reduce the effects of quantum noise in 'classical' ways.

There are several 'nonclassical' approaches that have been proposed within the community [64–66]. They can be put into four categories:

- (i) Injection of externally produced squeezed vacuum states of light into the signal output port [67–70]. When the optimal frequency-dependent rotation of the squeeze angle is achieved either using optical filters, or, by conditional squeezing and the interferometer as an optical filter [71], this technique yields a broadband suppression of quantum noise.
- (ii) Internal manipulation of the GWD cavities' gain bandwidth product by using negative dispersion media and internal squeezed state production including ponderomotive squeezing [72–76].
- (iii) Reshaping the quantum noise spectral density by dynamical back-action from test-mass motion to intra-cavity light power and vice versa (coherent feedback) producing e.g. the optical bar effect or the optical spring effect (associated with detuned signal recycling) [77, 78].
- (iv) Reducing quantum back action by performing a QND measurement, e.g. by using a speed-meter interferometer [79–82], or measuring the frequency-dependent mixture of phase and amplitude quadratures that is free from back-action (variational measurement) [74], or using auxiliary quantum systems with effective negative mass to cancel back-action noise [83]

All four categories utilize correlations within the light's quantum uncertainty (with respect to the electric fields at different phases, the so-called field quadratures) to gain the nonclassical sensitivity improvement. Generally, optical states having quantum correlations are more sensitive to optical loss than coherent states. This requires future detectors producing much less optical loss for light fields that enters the signal output port, travels along the arms and is back-reflected and detected by the photo diode. All categories also require introducing additional optics, increasing the complexity of the detectors, or even require a completely new interferometer topology.

The categories are not mutually exclusive and can be combined in different ways. In particular the injection of squeezed states can be combined efficiently with all other concepts.

Injection of squeezed vacuum states of light (for the reduction of shot noise) has been implemented in GEO 600 [31, 84] and was tested in one of the LIGO detectors [30]. In both cases only the shot noise was targeted, since radiation pressure noise was not an issue. This will change in future GW detectors. They will require filter cavities to avoid radiation pressure noise that is increased due to shot noise squeezing. Squeezed state injection, completed by filter cavities for the generation of frequency dependent squeezing and the mitigation of radiation pressure noise, is a potential technique for A+ and Voyager and beyond. All other approaches are potentially implementable on LIGO Cosmic Explorer but are too immature to be considered for A+ or Voyager. Intensive research and development is needed.

#### 4.1.1 A+

For A+ we propose the injection of frequency dependent squeezing at 1064 nm and a measured shot-noise squeezing of up to 6 dB. A list of loss contribution that allows for 6 dB of squeezing is given. This list is input for other groups, in particular to "optics" and "AIC". From Table 9 it is likely that 6 dB is achievable. Reaching the design performance of the aLIGO input Faradays and the aLIGO OMC, careful mode-matching and using now available photodiodes with quantum efficiency (QE)  $> 99\%$  [85] should see 6 dB shot noise reduction using currently available squeezers.

From Table 9 we see that 6 dB noise reduction with a total efficiency of 77% (losses less than 23%) requires phase noise to be controlled to better than 17 mrad. This is regularly achieved in small scale laboratory demonstrations. Operation of the squeezed light source in vacuum, recently demonstrated in [86], might be helpful for achieving this goal for A+.

R&D experimental tests include control system development for frequency dependent squeezing, optimizing squeezed light sources at 1064 nm and the characterization of low-loss Faraday isolators and mode-matching efficiencies.

#### 4.1.2 Voyager

For Voyager we propose the injection of frequency dependent squeezing at some frequency around  $2\ \mu\text{m}$  and a measured shot-noise squeezing of 8 to 10 dB. A list of loss contribution that allows for such squeezing levels is given Table 9. We see that the total efficiency required for 10 dB is 91.3% requiring phase noise to be controlled to better than 4.4 mrad. Achieving the levels of stability will likely require operating the squeezer inside the vacuum envelope. Smaller and more compact designs of squeezed light sources will also help to reduce phase noise.

Required R&D experimental tests: development of an audioband squeezed source of greater than 12 dB from 10 Hz to 10 kHz at  $2\ \mu\text{m}$ ; control system development for frequency dependent squeezing at  $2\ \mu\text{m}$ , sourcing and characterize low-loss Faraday Isolators at  $2\ \mu\text{m}$ , mode-matching efficiency, Brillouin scattering, and phase noise. Basically the state of the art for 1064 nm needs to be achieved at  $2\ \mu\text{m}$ . Up to 12.3 dB of squeezing at 1550 nm has been demonstrated at the AEI, but only down to about 10 kHz. The extension to 10 Hz, however, should be straight forward using the techniques invented for 1064 nm. High efficiency generation and detection of squeezed states should be possible up to a wavelength of

about  $2.1\ \mu\text{m}$ . The upper limit is set by the current availability of high-quantum-efficiency photo diodes.

### 4.1.3 3rd generation GW detectors

The longer term goal of 15 dB observed quantum noise reduction is extremely ambitious. From Table 9 we see that the total efficiency required is 97.4% requiring phase noise to be controlled to better than 2.5 mrad. Loss tolerances on the isolators, OMC, modematching etc far exceed anything currently achieved. 15 dB of squeezing was recently observed at the AEI [85] at 1064 nm at MHz frequencies. Similar results should be possible at 1550 nm and down to a few hertz, if photo diodes of the same high quantum efficiency are available. A photo diode quantum efficiency of greater 99.5% needs to be made available with surface areas suitable for GW detectors. Recently, wavelengths around  $2\ \mu\text{m}$  were discussed. At these wavelengths squeezed light sources as well as photo-electric detection needs to be researched.

Given the timeline for 3rd generation detectors, LIGO Cosmic Explorer and Einstein Telescope, the full list of quantum noise reduction concepts should be explored in addition to frequency dependent squeezing and need to be extensively researched. The next section summarizes the currently proposed concepts.

Recently, a variety of new configurations featuring QND measurement of speed have been invented (see Fig. 8). Some of these concepts offer not only improved quantum noise and relaxed requirements for filter cavities, but also the possibility of reduced coating thermal noise [10] and the introduction of negative inertia [6, 9]. Some can be implemented with minimal changes to the main Michelson interferometer infrastructure by means of using 2 orthogonal polarizations of light at the same time and by adding a few polarization optical components to the readout port [8, 11, 12]. Individual design studies are required to evaluate the technical challenges, infrastructure requirements and various benefits of these configurations.

## 4.2 State-of-the-art squeezed-light sources

Squeezed vacuum states of light [70] are produced by cavity-enhanced type I optical-parametric amplification (OPA), also called cavity-enhanced type I parametric down-conversion (PDC), sometimes also type I optical-parametric oscillation (OPO) below threshold. The most efficient source is a (second-order nonlinear) periodically poled KTP (PPKTP) crystal placed in a cavity that is pumped with second-harmonic light (below oscillation threshold). At threshold such a device produces, theoretically, an infinite squeeze factor, if the optical loss is zero and if the pump power is infinite. Great progress has been made over the last 10 years in the generation of squeezed vacuum states of light at 1064 nm and 1550 nm [69]. For LIGO Cosmic Explorer squeezed-light sources at wavelengths around  $2\ \mu\text{m}$  need to be developed and tested. PPKTP is expected to be again a suitable low-loss material.

1. Both ANU and AEI have measured squeezing of around 10 dB down to 10 Hz at 1064 nm. The AEI has observed greater than 12 dB at MHz frequencies at both 1064 nm and 1550 nm [87, 88]. Recently, the AEI together with the ILP Hamburg observed 15 dB of squeezing at 1064 nm at MHz frequencies [85].

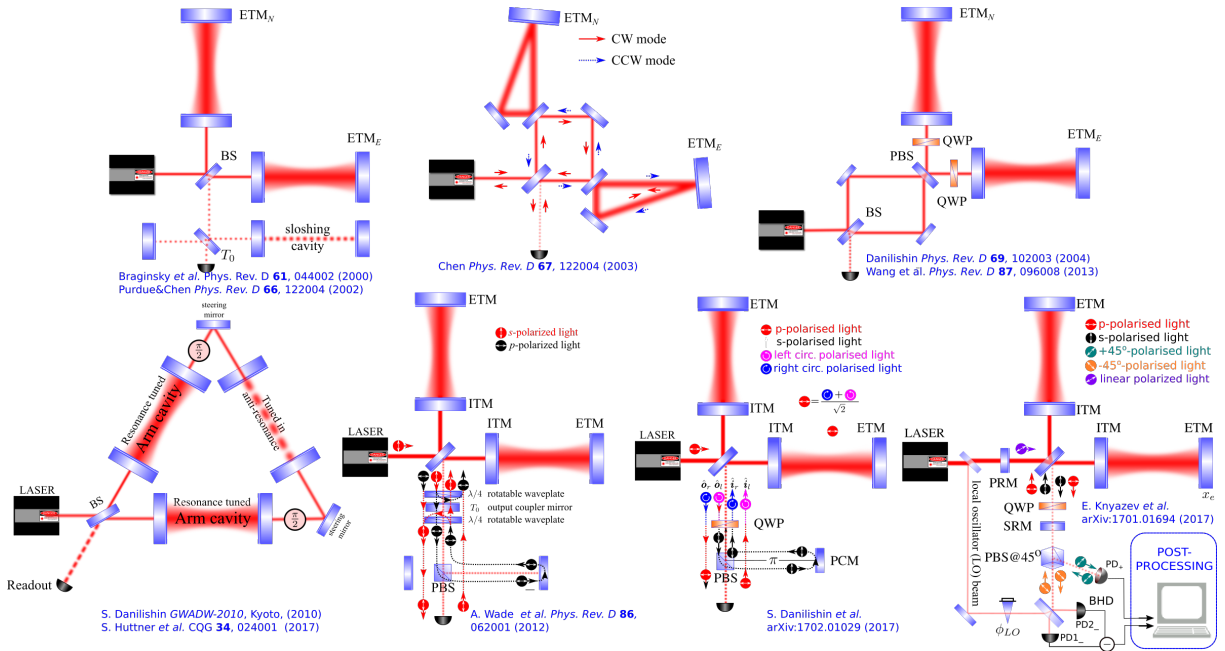


Figure 8: Different incarnations of speed meters [3–12].

2. GEO 600 has been operated now for several years [31] at a time regularly observing more than 3.5 dB sensitivity improvement with a highest value above 4 dB [89, 90]
3. An ANU designed squeezer was installed on the LIGO H1 detector operating in the S6 configuration. Sensitivity enhancement was observed above 200 Hz with 2.1 dB improvement above a few hundred hertz. Importantly, no excess noise was seen below 200 Hz and no additional glitching was found.
4. Photodiodes with quantum efficiency in excess of 0.99 at 1064 nm are now available [85].

Any absorption/scattering between the squeezing resonator and detection reduces the level of squeezing. Starting with a suitable source of squeezed light, Table 9 presents a comprehensive list of potential loss sources along with levels required to observe 6 dB, 10 dB and 15 dB of quantum noise reduction. The final column of the table presents the numbers achieved during the H1 squeezing test. The numbers are presented in terms of efficiencies (1 - loss). A filter cavity (see section 4.4.1) is needed to achieve broadband quantum noise suppression or to at least reduce the squeezing ellipse to a coherent state at frequencies where shot noise does not dominate.

As the level of squeezing is increased the requirement on stability of the squeeze angle is increased. As the angle rotates away from the detected quadrature angle the level of measured squeezing is reduced. In the presence of phase fluctuations the average squeeze factor is reduced. The effect gets worse as the level of anti-squeezing increases. Due to optical loss, the anti-squeeze factor is always larger than the squeeze factor. The tolerance on phase noise will decrease as larger squeeze factors are applied. This is depicted in Fig. 10 where measured squeezing loci are plotted for different values of total loss and total phase noise.

	Loss source	H1 experiment	Near term goal (6dB)	Longer term goal (10 dB)	Dreaming(15dB)	
1	OPO escape efficiency	96%	98%	99%	99.8%	
2	Injection path optics	80%	99.7%	99.7%	99.99%	
3	viewport		99.8%	99.8%	99.99%	
4	3 faraday passes		94%, 94%, unknown	97% each (aLIGO input Faradays)	99% each	99.7 % each
7	RF pick off beamsplitter (beam for ISCT4)		98.8%	99%	99.5%	99.8%
5	Reflection off of Signal recycling cavity@100 Hz	arm cavity and michelson =98%	97.5%(Tsr=35%)	99.2% (Tsr=50%)	99.5%	
6	Circulator for filter cavity	NA	98%	99.5%	99.8%	
8	Squeezer mode matching to OMC	71% (inferred from total)	96%	98%	99.7%	
10	OMC transmission	82%	97%	99.5%	99.7%	
11	QE of PDs		99%	99.7%	99.99%	
	Total efficiency (escape * detection)	40-45%	77.6%	91.3%	97.4%	
	Total phase noise allowable		17mrad	7 mrad	2.5 mrad	
	Measured squeezing (dB)		6	10	15.25	

For numbers in the H1 column- red numbers were measured during the experiment, numbers in black were made before installing components or estimates of losses that were too small compared to the total to measure accurately with the uncertainty in our measurements.

Figure 9: Table 11 Maximum acceptable losses to achieve 6 dB, 10 dB and 15 dB of observed quantum enhancement. From Dwyer [13]

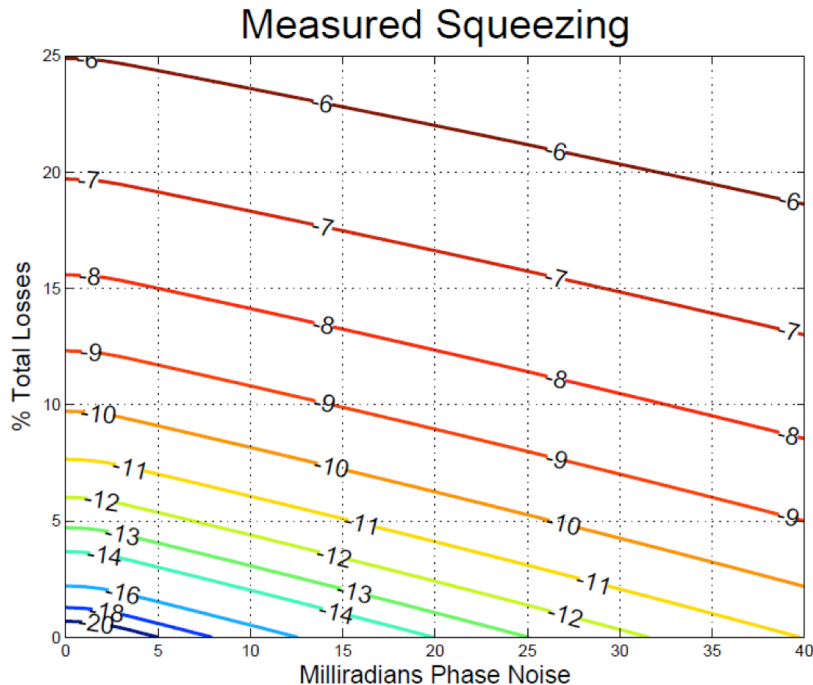


Figure 10: Measured squeezing for different values of total losses (the product of escape efficiency and detection efficiency) and total phase noise. For each value of phase noise, the nonlinear gain in the squeezing resonator is optimized to maximize the measured squeezing, but the nonlinear gain is capped at 90, (pump power at 80% of threshold). Operating the squeezing resonator closer to threshold doesn't improve the measured squeezing very much, but could make stable operation of the squeezer difficult. From Dwyer [13]

#### 4.2.1 Low-loss input and output optics

Utilizing externally generated squeezed vacuum states of light requires additional input and output optics. The most prominent ones are the Faraday rotator in combination with one polarising beam splitter for input/output coupling the squeezed mode, the detuned narrow linewidth filter cavity for engineering the frequency dependent rotation of the squeezing ellipse that allows for simultaneous squeezing of shot noise and radiation pressure noise, and the output modecleaner right before photo-electric detection. All these components should have a transmission of greater 99%, where the exact value depends on the squeeze factor aimed for and very low scattering. The Faraday rotator needs to provide precisely 45 degree rotation with high transversal homogeneity and very good anti-reflection coatings (cf. section 7.2.2). The extinction ratio of the polarising beamsplitter is another issue as well as the impedance-matching of the output modecleaner.

A general issue of GW detectors at signal frequencies below 100 Hz is given by parasitic interferences due to carrier light that is back-scattered from moving surfaces [91,92]. Generally, parasitic interferences can be mitigated through optics having higher surface qualities, by dumping and absorption of scattered light, and by reducing the motion of back-scatter surfaces. Recently, a new readout scheme was proposed and experimentally investigated that allows for a subtraction of back-scatter noise, if a model can be fitted to the parasitic interference [93]. The injection of squeezed vacuum states requires new optics that can produce



additional scattering. In order to avoid parasitic interferences from the squeezing resonator, the optical pass between the squeezing resonator and the signal recycling/extraction mirror must be kept constant. Additional Faraday isolators in between may be required.

Recently, it was shown in proof-of-principle experiments that the back-scatter mitigation readout in [93] can also be extended to the nonclassical regime using entangled light [94, 95]. Two sources of squeezed vacuum states of light overlapped on a (balanced) beam splitter produce two entangled light fields [96–98].

The filter cavity concept is outlined in 4.4.1, also see the AIC Chapter. For a given filter bandwidth  $\gamma_{\text{filter}}$  (to be determined by the needs of input/output filtering), when realized by a cavity of length  $L$ , the total loss  $\mathcal{E}$  is determined by

$$\mathcal{E} = \frac{4\epsilon}{T} = \frac{\epsilon c}{\gamma_{\text{filter}} L} \quad (1)$$

where  $T$  is the input-mirror power transmissivity [related to bandwidth by  $\gamma_{\text{filter}} = Tc/(4L)$ ] and  $\epsilon$  is the loss per round-trip. It is therefore the ratio  $\epsilon/L$  that determines the goodness of the filter. Since the per-round-trip loss  $\epsilon$  depends on the beam spot size, which in turn depends on  $L$ , an optimization is needed to find out the optimal length and design of filter cavities [99].

The effect of losses is further amplified if back-action evasion is required, in which case the signal strength in the quadrature being detected is significantly less than conventional situations. A rule of thumb for this limitation is available from Kimble et al. [74], where we have

$$\sqrt{S_h/S_h^{\text{SQL}}} \geq (e^{-2q}\mathcal{E})^{-1/4} \quad (2)$$

where  $\mathcal{E}$  is the power loss, and  $e^{-2q}$  is the power squeezing factor. Assuming  $\mathcal{E}$  to be 0.01, and 10 dB squeezing, we have a SQL-beating limit of 0.18.

Figure 11 shows the effect of round-trip loss on the degradation of squeezing for a filter cavity with 22 Hz detuning and bandwidth. The colors correspond to the quantum-noise spectral density along the minor axis of the noise ellipse relative to coherent vacuum noise. It can be seen that round-trip loss over 50 ppm degrades the squeezing to a level comparable to coherent vacuum around frequencies near the filter resonance. Here and in the following, round-trip loss are considered excluding the transmission of the input mirror. The input transmission is comparatively large (about 550 ppm for the filter calculation of Fig. 11) and does not directly contribute to squeezing degradation. Round-trip loss is dominated by optical scatter loss from imperfect mirrors in combination with finite aperture size. The goal is to minimize scatter loss.

Scatter loss in cavities is determined by properties of the individual mirrors such as mirror size, mirror curvature, mirror surface aberrations and defects, and by the optical path length of the cavity. Scatter loss in cavities is not fully understood yet. A concerted effort of high-finesse cavity experiments, scattering measurements, numerical simulations and theoretical work is required to form a consistent picture between loss models and observations. Practically speaking, ultra-low losses (around 1 ppm) have been achieved on the mirrors of fixed cavities [100, 101]. However, the lowest loss measured on the large, test-mass-sized beams are more usually in the 50-100 ppm range. FFT simulations have shown that the loss for large

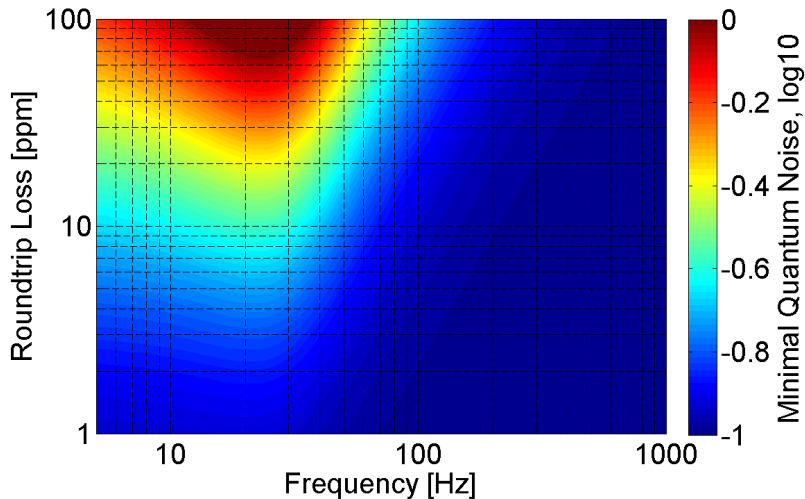


Figure 11: Degradation of 10 dB external squeezing reflected from a 300 m filter cavity as function of round-trip loss (in addition to input transmission) and frequency.

beams is dominated by the large scale figure error of the substrate, while the losses for small beams are dominated by point defects in the coatings. Since the low frequency performance of the QND schemes so strongly depends on the loss for intermediate sized ( $\sim$ mm) beams, it is vitally important to develop ultra-low loss mirrors for this beam size. The modern polishing technology is already good enough.

Additional more specific problems that need to be investigated include small-angle scatter loss in cavities, coherent loss versus incoherent loss, and round-trip loss as a function of cavity aspect ratio. Coherent scatter loss designates the effect when the spatial pattern of higher-order modes resonating inside the cavity matches the pattern of light scattered from a mirror. In this case scattered light can build up resonantly provided that the overall loss in these higher-order modes is not too high. In general, higher-order modes experience more loss since they describe a wider spatial intensity distribution that directs additional light power beyond the mirror aperture compared to the target mode of the cavity, which is usually the Gaussian shaped  $TEM_{00}$  mode. For this reason it is also necessary to understand how much round-trip loss is influenced by the cavity aperture, or better by the cavity aspect ratio, which is the aperture divided by the cavity length. Coherent scatter loss can potentially lead to a significant increase of the total scatter loss, and it is no more describable by simple perturbative scatter models.

Losses in mirrors are also addressed in the Optics section 5.

### 4.3 High-quantum-efficiency photo-electric detection

#### 4.3.1 PIN photo diodes

Photo-electric detection in the output of a GW detector is done with unity gain semiconductor PIN photo diodes, where PIN stands for positive-intrinsic-negative. Such a detector

ideally elevates exactly one photo electron to the conduction band for every incident photon. An imperfection is treated as optical loss. Already GW detectors operated with light in coherent states require photo diodes with high quantum efficiency since the square root of its signal normalized shot-noise spectral density scales with  $1/\sqrt{n}$ , where  $n$  is the number of *detected* photons per second. For a GW detector that uses any nonclassical techniques the effect of imperfect quantum efficiency is stronger [70].

Recently, the quantum efficiency of a custom-made InGaAs photo diode was calibrated to  $(99.5 \pm 0.5)\%$  [85], where two identical such photo diodes were used in a home-made balanced homodyne detector arrangement. The two photo diodes were manufactured from the same wafer material as the photodiode that has been used in GEO 600 since 2011 [31, 84]. The photo diodes allowed for the observation of 15 dB squeezing from a squeezing resonator with about 1% escape efficiency, negligible phase noise and access noise, and propagation loss of another 1% including imperfect modematching between the squeezed vacuum mode and the optical local oscillator of the balanced detection scheme.

At 1550 nm custom-made photo diodes with quantum efficiencies greater 99% seem might already exist. An indication for that is the observed squeeze factor of 12.3 dB at this wavelength [88]. The response of standard InGaAs photodiodes show a cutoff at 1600 nm, but extended InGaAs is available that allows photo-electric detection to wavelengths as long as 2.5  $\mu\text{m}$ . For LIGO Cosmic Explorer it needs to be investigated up to what wavelength quantum efficiencies of close to 99% can be realized. Further discussion on the development of high efficiency photodiodes can be found in section 7.3.2.

### 4.3.2 Balanced homodyne detection (BHD)

All table-top experiments that observed squeeze factors of greater than 8 dB used two identical PIN photo diodes as a balanced receiver. In this arrangement the dim squeezed signal beam is overlapped with a modematched local oscillator of the same carrier wavelength on a balanced beam splitter and the different voltage from the photo diodes is recorded. Ideally the signal beam does not contain any carrier component. In any case, the power ratio of local oscillator (LO) and signal beam should be as high as possible.

The main advantage of a balanced homodyne detector compared with a single photo diode is that an arbitrary (frequency independent) field quadrature can be detected, which is required in many nonclassical approaches, such as detuned signal recycling, variational input/output and speed meters. Ideally the local oscillator mode corresponds to the carrier mode inside the interferometer to avoid any differential phase noise. Tapping such a local oscillator, e.g. from the AR-coated surface of the beam splitter or the POP port, and delivering it with the required low phase noise to the detection system is the main challenge and it will involve additional hardware in terms of low noise suspensions for beam delivery and an addition output mode cleaner [102].

The BHD scheme, however, provides additional advantages. The local oscillator, assumed to be in the ideal transverse mode, acts as an output mode cleaner. Also does the BHD scheme allow for an interferometer operation exactly at dark port.

Recently, noise couplings in the BHD scheme was experimentally investigated in the context

of GW detectors [103]. Additional research is required to solve the problem of tapping the LO inside the interferometer.

## 4.4 Interferometer techniques supporting nonclassical quantum noise reduction

### 4.4.1 Frequency-dependent squeeze angle (input filtering)

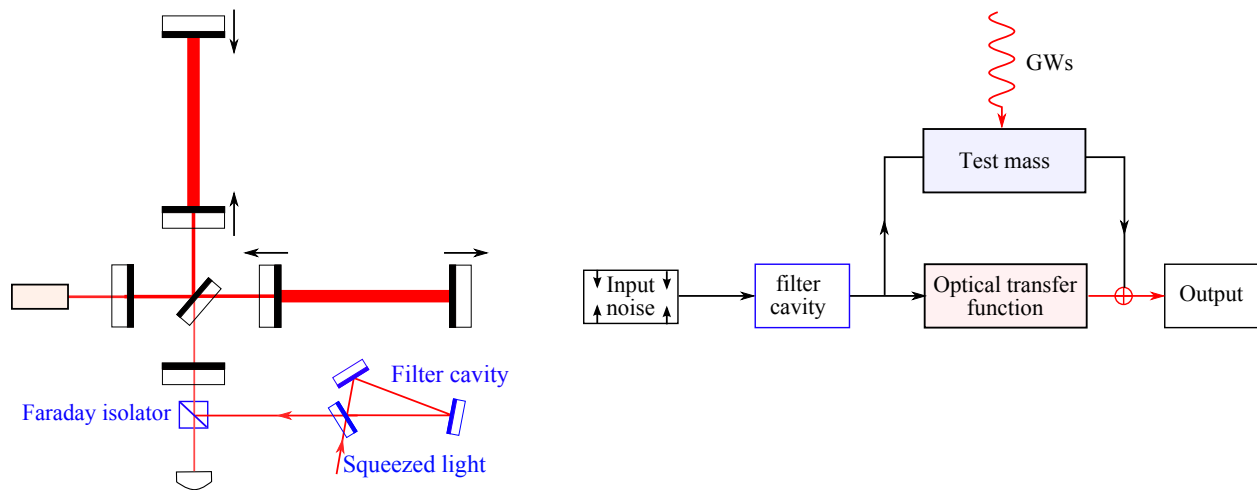


Figure 12: Schematics showing the frequency-dependent squeezing scheme (left) and its associated flow chart (right).

The first scheme is ‘frequency-dependent squeezing’ (as discussed in the previous sections). This term refers to a broadband spectrum of squeezed vacuum states for which the squeeze angle is frequency dependent and optimized for the cancelation of back-action noise and for achieving the optimum signal to noise ratio. The frequency dependence of the squeeze angle can be achieved by reflecting off the squeeze spectrum from a detuned narrow band cavity. There might be other ways how the same frequency dependence can effectively be achieved. Research on alternatives might be very valuable for future gravitational-wave detectors.

As shown schematically in Fig. 12, it utilizes an optical (filter) cavity to rotate the amplitude and phase quadratures, or equivalently the squeeze angle, in a frequency-dependent way. If the parameters of the filter cavity is appropriately specified, one can rotate the squeeze angle such that the quantum noise spectrum is reduced by an overall factor that is equal to squeeze factor.

For illustration, in Fig. 13, we show the resulting noise spectrum in *the ideal case without optical loss*. As we can see, the squeeze angle rotates in such a way that at low frequencies the fluctuation in the amplitude quadrature is squeezed—thus reducing the radiation-pressure noise, while at high frequencies the phase quadrature is squeezed—thus reducing the shot noise. In order to achieve the desired rotation of squeeze angle, the filter cavity needs to have a bandwidth that is comparable to the detection bandwidth—this indicates a high-finesse cavity is necessary if the cavity length is short. The specification for the filter cavity can almost be analytically calculated by using the method outlined in [81].

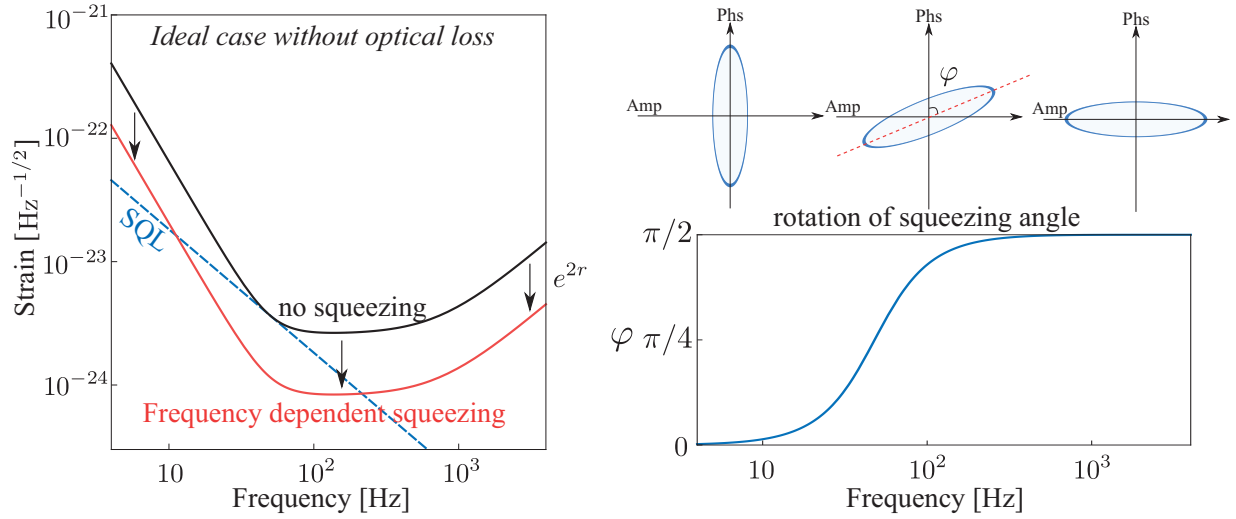


Figure 13: Noise spectrum for frequency-dependent squeezing (left) and rotation of the squeeze angle (right).

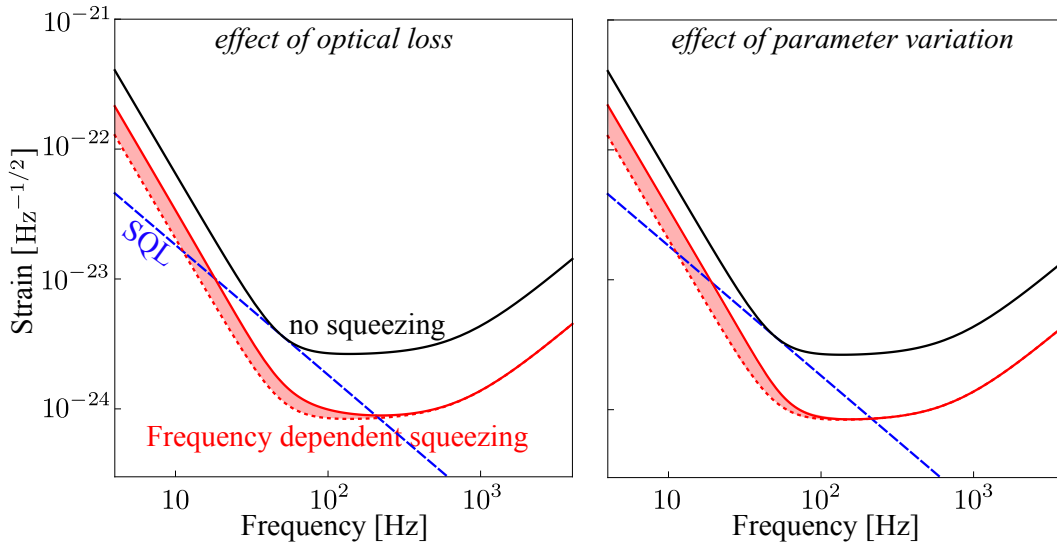


Figure 14: The effect of optical loss (left) and the parameter variation of the filter cavity (right) for input filtering. The shaded regions illustrate the degradation in sensitivity. Here we have assumed the total optical loss of 20% (round-trip loss multiplied by the number of bounces inside the cavity) and parameter variations of 10%.

In reality, the optical loss will affect the performance of input filtering, as shown in the left panel of Fig. 14. Additionally, parameters of the filter cavity, in particular, the transmissivity of the mirrors and the detuning, cannot be exactly set to the optimal value, and their variation will influence the sensitivity in a similar manner to the optical loss, as illustrated in the right panel of Fig. 14.

#### 4.4.2 Frequency dependent readout phase (output filtering)

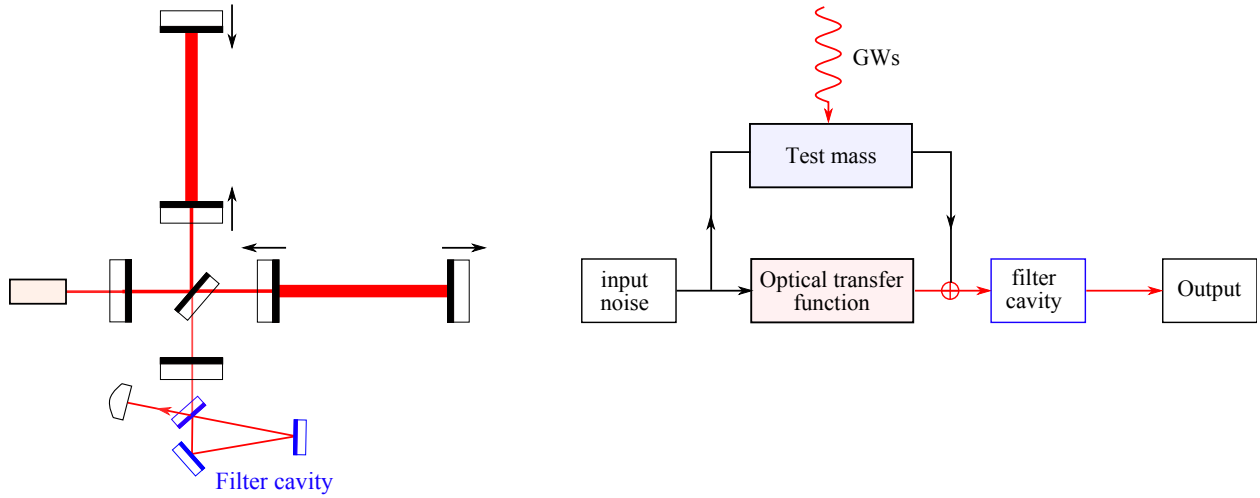


Figure 15: Schematics showing the frequency dependent (or variational) readout scheme (left) and its associated flow chart (right).

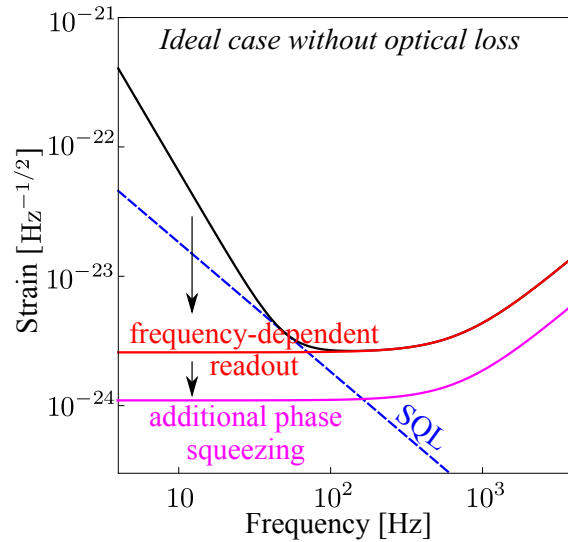


Figure 16: The noise spectrum for the frequency-dependent readout scheme.

A close related counterpart to the input filtering is the variational readout, and as shown schematically in Fig. 15, it uses an optical cavity to filter the detector output which allows one to measure different optical quadratures at different frequencies. The filter cavity has the same functionality as in the case of the frequency-dependent squeezing—the only difference is that it rotates the optical quadratures of the output instead of input. In the ideal case, this scheme can coherently cancel the radiation-pressure noise at low-frequencies, and give rise to a shot-noise only sensitivity [74]. In Fig. 16, we show the resulting noise spectrum in the ideal lossless case. In reality, due to the presence of optical loss and parameter variation of the filter cavity, such a cancellation cannot be perfect, as illustrated in Fig. 17.

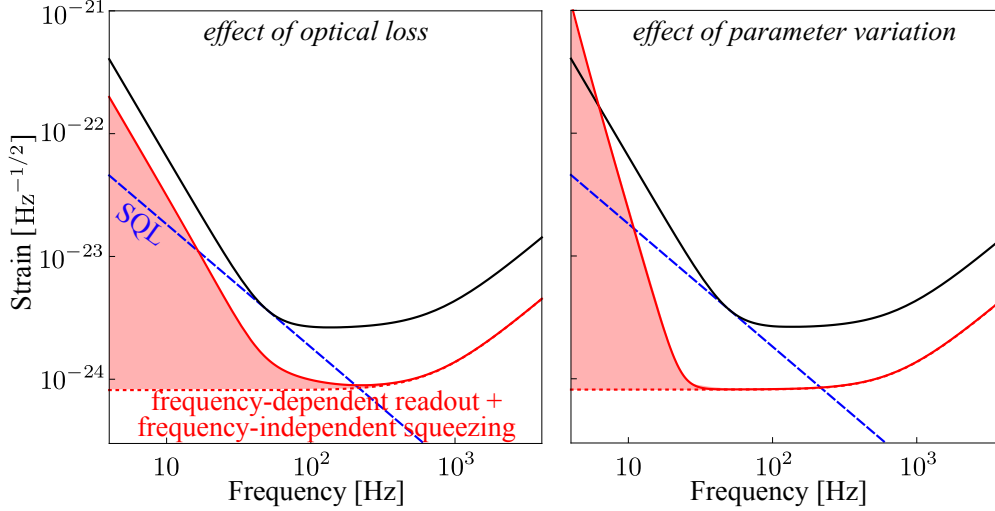


Figure 17: The effect of optical loss (left) and parameter variation of the filter cavity (right) for frequency-dependent readout (output filtering). Similar to Fig. 14, we have used a total optical loss of 20%. In contrast, the parameter variation is chosen to be only  $10^{-4}$  in order to produce reasonable sensitivity, as it is much more sensitive than input filtering.

As this technique leads to reduced signal (but increased signal/quantum noise) at low frequencies it is far more sensitive to optical losses compared with input filtering.

#### 4.4.3 Conditional Squeezing

*Conditional squeezing* is a new concept that was proposed to achieve frequency-dependent squeezing without an additional filter cavity [71]. The concept is based on Einstein-Podolsky-Rosen entanglement that exists within a single squeezed cavity mode as demonstrated in [104]. The entanglement is utilized by separating the upper and lower sidebands and using two balanced homodyne detectors with different local oscillator frequencies. The new concept uses two local oscillators with a frequency difference of about a multiple of the free spectral range of the signal extraction cavity. The observed frequency dependent squeezing is conditional because it is revealed only after combining the noise from both readouts.

Implementation of this new concept in a large scale GWD implies separate readout and injection channels for both entangled beams. A thorough analysis of the effect of losses throughout the interferometer, arm length asymmetries, and imperfect separation of the signal and idler beams were considered for the GEO600 interferometer in [105]. This study has shown that conditional squeezing results in comparable performance to frequency dependent squeezing with filter cavities, yet at a price of 3 dB extra squeezing and with rather stringent constraints on injection and readout loss. Experimental table-top tests of the concept are underway in ANU and in Hamburg.

#### 4.4.4 Twin signal-recycling

Twin signal recycling (TSR) refers to detuned signal recycling in which upper *and* lower signal sidebands are resonantly enhanced simultaneously providing a spectral density which

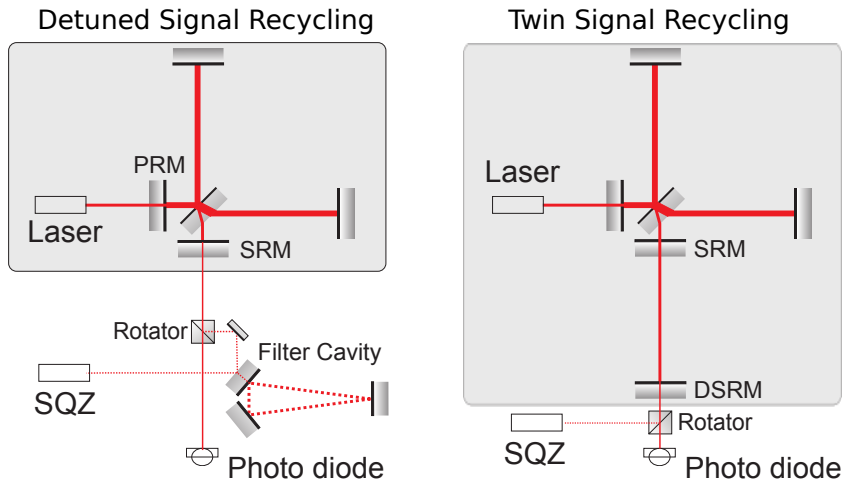


Figure 18: (Color online) Top left (shaded): Topology of the current gravitational wave detector GEO 600. The mirror in the laser input port (PRM) realizes so-called power-recycling. The signal-recycling mirror (SRM) in the output port establishes a carrier light detuned single-sideband signal recycling cavity. Bottom left: Extension for a broadband shot-noise reduction utilizing squeezed states. Right: Topology, proposed here. Two optically coupled cavities are formed with the help of an additional mirror DSRM. Their resonance doublet enables detuned dual-signal-recycling resulting in lower shot noise. Squeezed states can be used without additional filter cavity.

is reduced by up to a factor of two [106]. TSR requires a an additional signal recycling mirror, most likely in a long-baseline arrangement. Squeezed states, however, can be used without setting up another long baseline filter cavity [107].

#### 4.4.5 Speed meter

Normally, a GW detector measures the test mass position at different times to infer the signal. However, position at different times does not commute with the Hamilton operator of a free mass. According to quantum measurement theory [108], such a measurement process inevitably introduces quantum back action and perturbs the test mass motion. (In the context of GW detectors, the back action is the radiation-pressure noise.) In order to evade back action, one needs to measure the conserved dynamical quantity of the test mass—the momentum. The latter is (approximately) proportional to speed, which is why a *speed meter* is ideal for measuring gravitational waves with greatly reduced radiation-pressure noise [79].

Several practical speed-meter configurations were proposed over the last 15 years (see Fig. 8) which fall into 3 distinct categories by the shape (frequency dependence) of their response to the GW signal. These are the *Sagnac-type speed meters* [5, 7, 12, 82], *sloshing speed meters* [8, 10, 109], and a recently proposed scheme that uses *EPR-type* entanglement between different polarizations [11].

In Fig. 19, we show the two simplified schemes of a sloshing and Sagnac-type speed meters. The sloshing speed meter, proposed in Ref. [109], uses a sloshing cavity. We can gain a



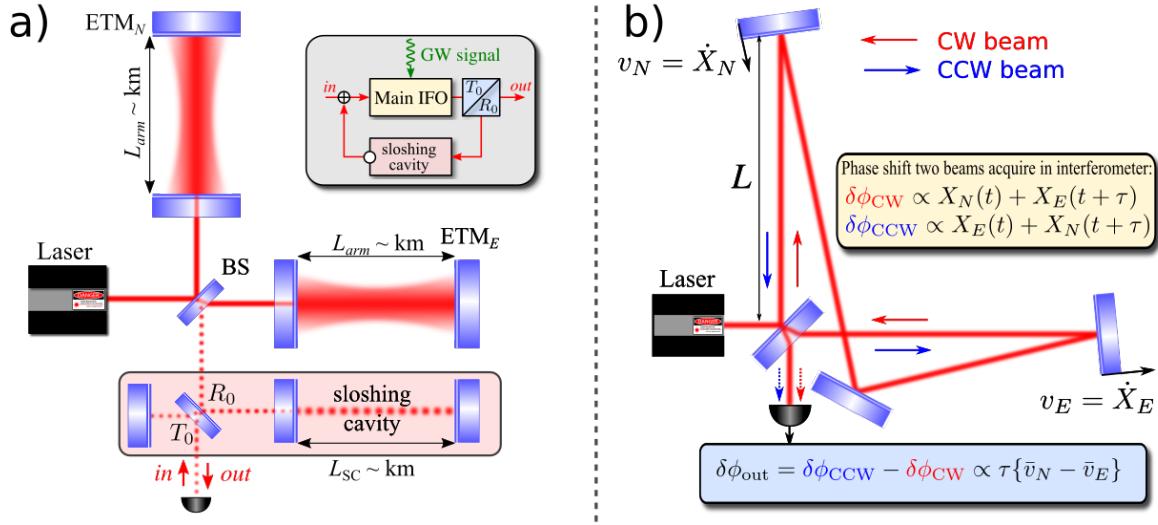


Figure 19: Two variants of implementation of speed meter interferometers, (a) the sloshing speed meter, and (b) the Sagnac speed meter. Inset in the grey rectangle in (a) represents the block diagram of the sloshing speed meter principle of operation. Here ETM stands for end test mass, BS is a beam splitter, and  $T_0 = 1 - R_0$  is the (power) transmissivity of the output coupling mirror.

qualitative understanding of how such a scheme allows us to measure the speed of the test mass. Basically, the information of test mass position at an previous moment is stored in the sloshing cavity, and it coherently superposes (but with a minus sign due to the phase shift in the tuned cavity) with the output of the interferometer which contains the current test mass position. The sloshing happens at a frequency that is comparable to the detection frequency, and the superposed output is, therefore, equal to the derivative of the test-mass position, i.e., the speed.

The zero-area Sagnac speed meter derives the speed signal from the subtraction of the two counter-propagating light beams at the main beam-splitter. The light beams visit both arms and interact with the test masses sequentially, but in an opposite order. As a result, the phase of the beams carries the information about displacement of the test masses, but with a delay  $\tau$  that takes light to travel between the arms. When the two beams recombine at the beam splitter, the phase at the readout port of the interferometer turns to be proportional to the relative mean velocity of the differential arm motion,  $\delta\phi_{out} \propto \tau\{\bar{v}_N - \bar{v}_E\}$ .

The typical noise spectrum of speed meters is shown in Fig. 20. The low-frequency spectrum has the same slope as the standard quantum limit, which is a unique feature of speed meter. When the optical power is high enough, we can surpass the standard quantum limit.

As a rule, no filter cavity is needed for the speed meter configuration, as the radiation pressure noise at low frequencies is cancelled. However, such a cancellation is achieved by choosing the homodyne detection angle that is deviated from the phase quadrature, therefore decreasing sensitivity at high frequencies. The deviation is proportional to the optical power. With frequency-dependent squeezing, we can reduce the effective optical power seen by the

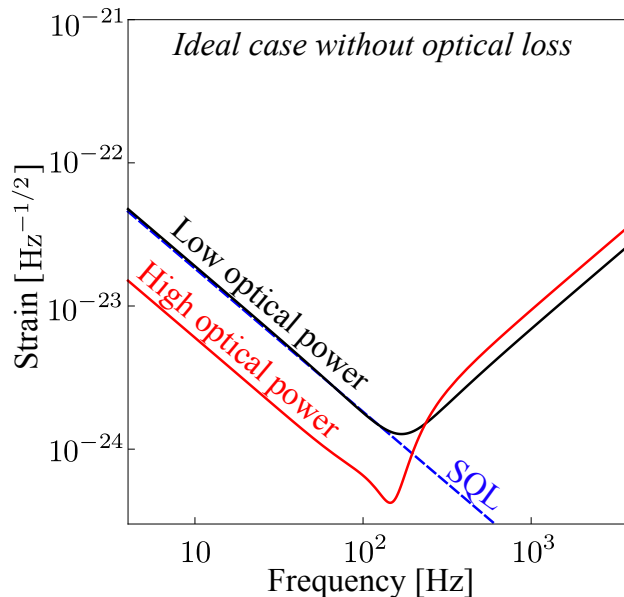


Figure 20: Plot showing the noise spectrum for the speed-meter configuration for two different optical powers.

test mass, which allows us to measure the quadrature closer to the phase one, and thus enhance the high-frequency sensitivity. Similarly, the frequency dependent readout allows us to cancel the low-frequency radiation pressure noise without sacrificing the high-frequency sensitivity by rotating the readout quadrature to the phase one at high frequencies.

#### 4.4.6 Local Readout / Optical Bar

The local-readout scheme is shown schematically in Fig. 21. It is a special case of a dual-carrier scheme—the second carrier light is only resonant in the power-recycling cavity and is anti-resonant in arm cavity (barely enters the arm cavity). The local readout scheme was first proposed in Ref. [110] and was motivated by trying to enhance the low-frequency sensitivity of a detuned signal-recycling interferometer, which is not as good as the tuned signal-recycling due to the optical-spring effect.

The idea of the local readout is based on that of the *optical bar*, proposed by Braginsky and co-authors [77]. At frequencies below the optomechanical resonance, the input test masses (ITMs) and end test masses (ETMs) are connected in a rigid way via dynamical backaction. For this reason, at these low frequencies, the distance to the ETMs due to a gravitational wave result in a local motion of the ITMs. This first of all reduces the sensitivity at these frequencies. The idea behind the local readout is to measure the motion of the ITMs locally and to use this information in the data processing.

### 4.5 Development of a QND apparatus which is quantum radiation pressure dominated

The implementation of these advanced ideas in the next generation detectors, in particular in the Einstein Telescope and LIGO Cosmic Explorer, cannot be done without first demon-

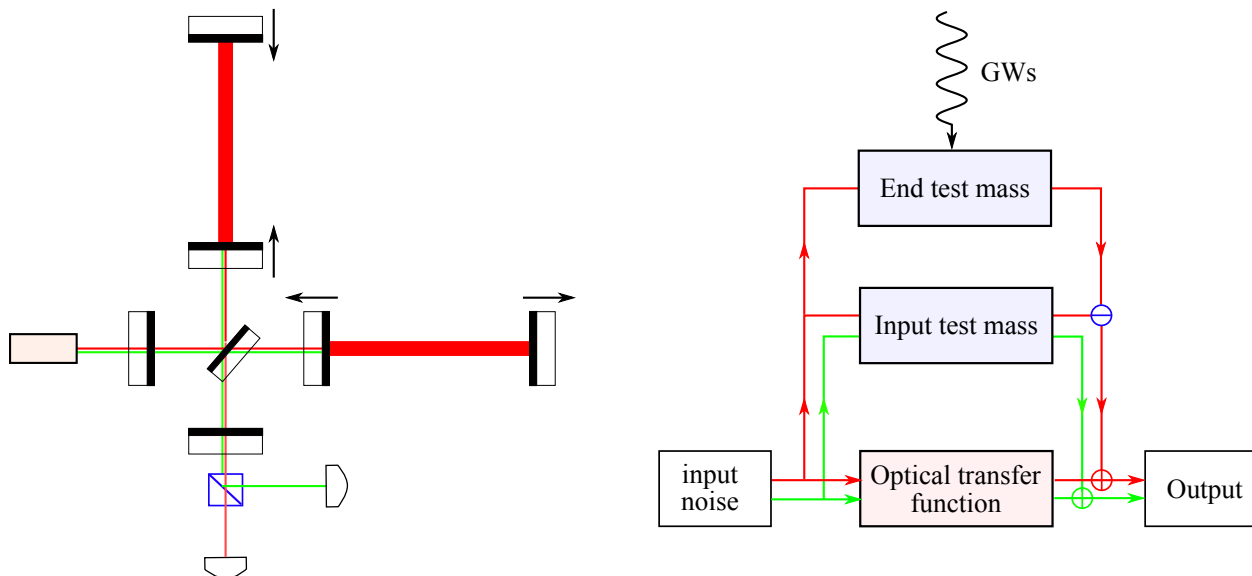


Figure 21: Diagram of the local-readout topology (left) and the resulting feedback loops (right).

strating them in a lab on the prototype facilities. The first experimental demonstration of quantum radiation pressure noise in a broadband regime at room temperature in a cavity with a movable 55-ng mirror was done at LSU this year [111], but to date no experiment has observed quantum radiation pressure noise on gm-scale let alone kg-scale mirrors, let alone reached the ‘naive’ standard quantum limit. Efforts to demonstrate QND interferometry are crucial to understanding what is achievable, to learn about problems which could mask such phenomena and how to beat such limits.

There are major activities planned or underway at the AEI 10 m prototype; MIT; the University of Tokyo and LSU. The Glasgow 10 m, the Gingin Facility and the ANU have embarked on testing optical spring dynamics. More effort is needed toward observing QRP noise. The first serious effort to demonstrate a quantum speed meter is under construction at the University of Glasgow. As such experiments run up against excess noise sources and thermal noise they will inform activities across other working groups.

Obviously, without QRP noise and SQL limited apparatus, no direct tests of these ideas can be performed. However, measuring transfer functions, demonstrating low-loss manipulation of squeezed states and variational readouts can be performed with shot noise limited systems. Plans are underway for such an experiment at MIT and AEI. More effort is needed.

#### 4.6 Numerical Optimization and Comparison

To find out the possible candidate that is feasible for upgrading the sensitivity of aLIGO, we made a systematic comparison of those configurations mentioned before, and the results are reported in the LIGO document [112]. Here we will just show the main result and one can refer to the document for more detail.

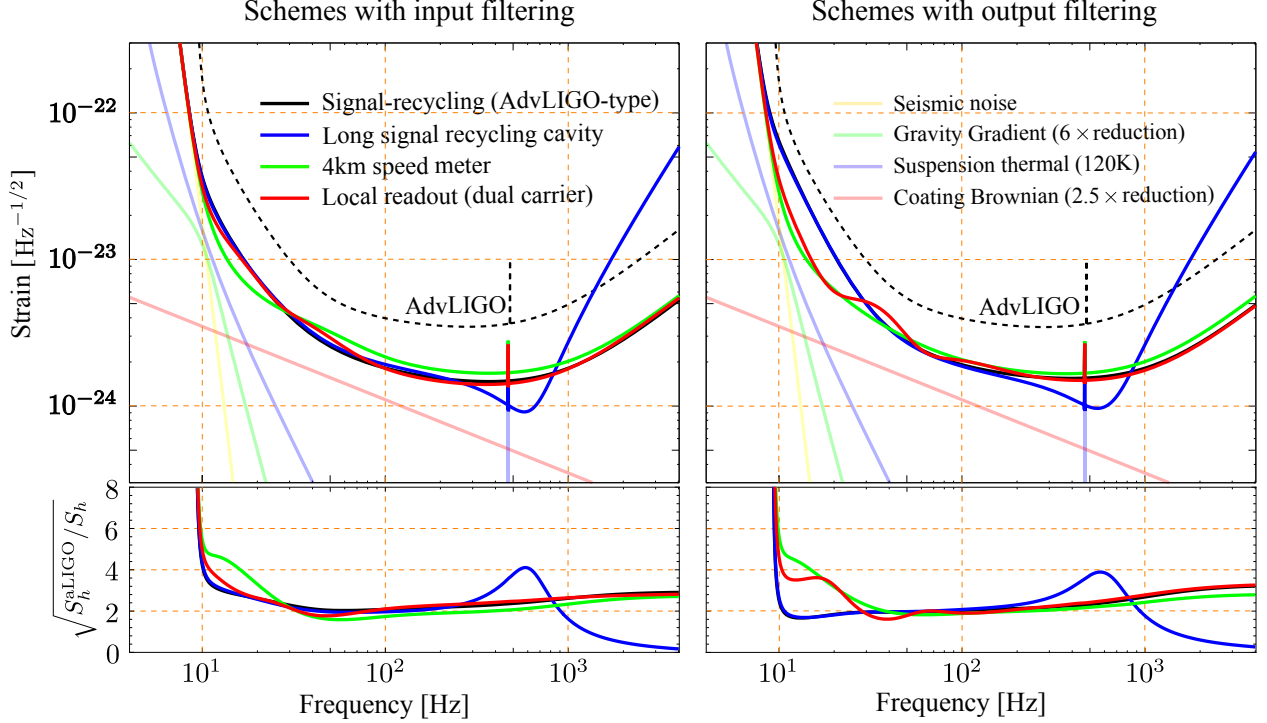


Figure 22: The optimized total noise spectrum for different schemes assuming a moderate improvement of the thermal noise compared with aLIGO baseline design. The lower panels show the linear strain sensitivity improvement over aLIGO.

The final optimization result critically depends on the cost function. In the literature, optimizations have been carried out by using a cost function that is source-oriented—trying to maximize the signal-to-noise ratio for particular astrophysical sources. Here we apply a rather different cost function that tries to maximize the *broadband* improvement over aLIGO. This follows the same philosophy of designing aLIGO which aims at a factor of 10 broadband improvement over initial LIGO.

For optimization, we also take into account the various classical noise sources (to be distinguished from the quantum noise). In particular, we consider Brownian thermal noise in the mirror suspensions [113, 114], seismic vibrations propagating to the mirror [115], terrestrial gravitational fluctuations [58, 116], and Brownian thermal fluctuations of the mirror (and mirror coating) surface [36, 117].

The results of the numerical optimization are shown in Figs. 22 and 23, where we plot the total noise spectra (the quantum noise + the classical noises) for different configurations with frequency dependent squeezing (input filtering) and frequency dependent readout (output filtering), respectively. In producing Fig. 22, we assume a moderate reduction in the thermal noise and the same mass and optical power as those for aLIGO. In producing Fig. 23, we assume a more optimistic reduction in the thermal noise, the mirror mass to be 150 kg and the maximum arm cavity power to be 3 MW.

As we can see, by adding just one filter cavity to the signal-recycled interferometer (aLIGO topology), we can already obtain a broadband improvement over aLIGO. Limited by thermal

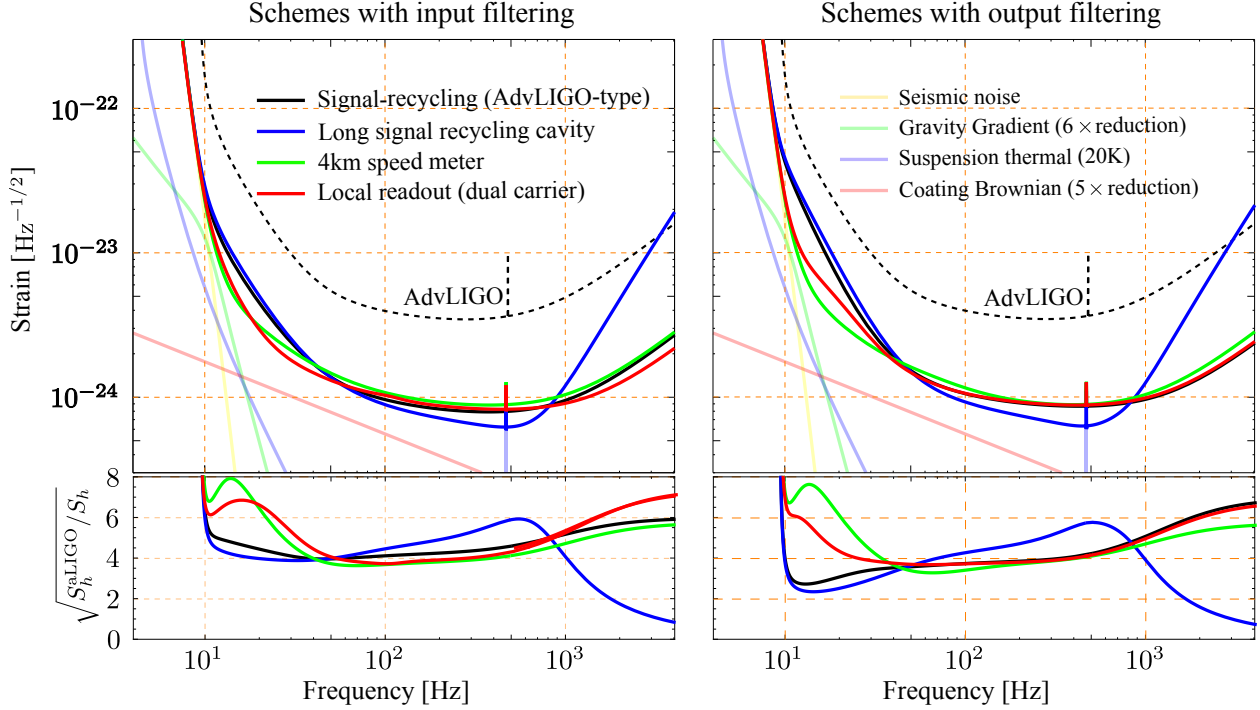


Figure 23: Optimization results for different schemes assuming more substantial thermal noise improvements, increasing the mirror mass from 40 to 150 kg, and increasing the arm cavity power from 800 to 3000 kW.

noise at low frequencies, the difference among these configurations is not very prominent. This leads us to the conclusion that adding one input filter cavity to aLIGO seems to be the most feasible approach for upgrading in the near term, due to its simplicity compared with other schemes. If the low-frequency thermal noise can be reduced in the future, the speed meter and the multiple-carrier scheme can provide significant low-frequency enhancement of the sensitivity. This extra enhancement will, for some low enough thermal noise, be enough to compensate for the extra complexity.

## 4.7 Development of other signal-to-quantum-noise enhancement techniques

### 4.7.1 External source of ponderomotive squeezing

Squeezed light can be produced by mirror motion under radiation pressure — this is called ponderomotive or opto-mechanical squeezing. Ponderomotive squeezing was demonstrated recently [118]. A squeezed factor of 1.7 dB at about 1.5 MHz was achieved, in a setup with a SiN membrane. A ponderomotive squeezing experiment has been going on at MIT [119]. See Fig. 24 for a sample configuration.

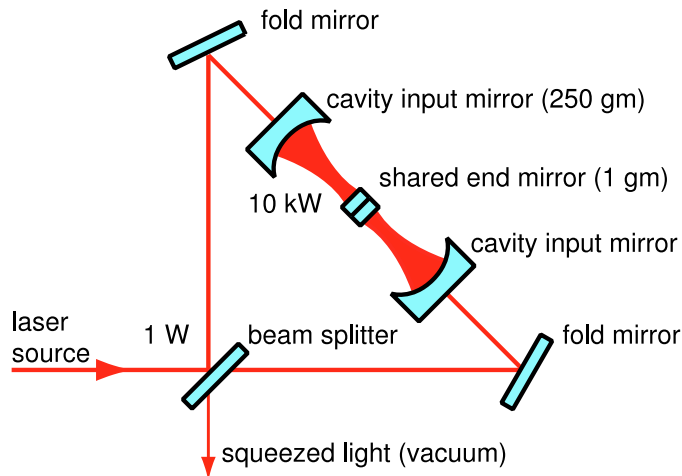


Figure 24: Ponderomotive Squeezer, taken from Corbitt et al. [14]

#### 4.7.2 Intra-cavity squeezing and white-light cavity

Independent of the injection of externally produced squeezed light, intra-(signal-recycling) cavity squeezing might result in an additional signal to quantum noise improvement [76]. Intra-cavity squeezing modifies the peak sensitivity - bandwidth product in a non-trivial way, from which the sensitivity of signal-recycled gravitational-wave detectors could benefit. Such intra-cavity squeezing can also be used to control and modify the optical spring effect by adjusting the squeezing angle and strength [120]. Theoretical as well as experimental work is required to research and develop appropriate schemes based on crystal squeezers.

It was theoretically shown that negative dispersion can result in a white light cavity without significantly adding quantum noise [121–123]. Traditionally, the type of mediums which provides negative dispersion are atomic media, but they introduce an additional constrain on the wavelength a gravitational-wave detector can be operated at and a thorough investigation of all noise sources of such systems is required. It should be investigated whether solid state optical systems at suitable wavelengths can also provide the same property. Recently, unstable opto-mechanical filters were proposed [124], however, such filters are susceptible to the thermal noise of the mechanical oscillator, which puts stringent requirements on the quality factor of the mechanical oscillator and the environmental temperature that the filter operates at. This thermal noise issue might be mitigated with the idea of optical dilution, which uses the optical spring effect and is currently under study. Additionally, the details of the control scheme for suppressing the instability still needs to be carefully researched.

#### 4.7.3 Opto-mechanically Induced Transparency

Opto-mechanically induced transparency (OMIT) [125] corresponds to electromagnetically induced transparency (EIT) in atomic vapors with more freedom in choosing the laser wavelength. OMIT can slow down light, possibly improving high finesse optical filters [126]. With enough slowing, one can almost eliminate the need for long filter cavities. This research needs to be pushed into a direction that does not introduce an unrealistic constrain on the gravitational-wave detector wavelength [127].

#### 4.7.4 Negative-mass/frequency spin systems

Multi-atomic spin ensembles proposed in [128] and demonstrated experimentally in [129, 130] possess a set of unique features which make them attractive for use in a GW detector. Under certain conditions, the collective spin of the atoms behaves like a harmonic oscillator with the frequency that can be made both positive and negative. Moreover, the interaction Hamiltonian of this spin system with light has the same shape as the opto-mechanical interaction of a mirror with probing light.

If the effective frequency of a spin oscillator is negative, the effective radiation-pressure noise introduced into the probing light has the sign opposite to one in the case of an ordinary opto-mechanical system. Therefore, the negative-frequency atomic spin ensemble can be used to cancel the quantum back action noise.

Atomic spin ensembles, however, are very selective in terms of light wavelength, as it has to match the certain transition frequencies between the Zeeman energy levels, and to date no atomic ensemble that works with 1064 nm, 1550 nm or 2  $\mu\text{m}$  wavelength is identified. This problem can be solved by using entangled dual-frequency squeezed beams like in the case of conditional squeezing. A scheme of GW detector based on this idea was proposed in [83]. It offers a similar broadband suppression of quantum noise as the conditional squeezing scheme but with an advantage of a simpler readout optics. The preliminary estimation of loss influence on its performance was done in [83], yet the design of a practical scheme based on this principle requires further theoretical and experimental R&D and prototyping.

#### 4.7.5 Other

There is the general question of whether atom interferometry offers a competitive topology for a third generation detector. Whilst this research area is not a fundamental activity of the LSC we should keep a watching brief on this technology and provide scientific support and advice when and where we can.

## 5 Optics

The Optics Working Group (OWG) of the LSC pursues research related to the development and implementation of optical components for ground-based gravitational wave detectors. This includes work on optical components being installed in Advanced LIGO, to better understand their behavior during commissioning and operation; possible upgrades to subsystems of Advanced LIGO including core optics, input optics, and auxiliary optics; and longer term research into ways around significant limitations in current detectors to be implemented in future generations of interferometers.

The OWGs work can broadly be divided into the following categories:

**Mirror coating research** The high-reflection (HR) coatings on the test masses must satisfy a number of performance criteria including low absorption, low scatter, high uniformity, high reflectivity at two wavelengths, low mechanical loss, and low thermo-optic noise. Of these, mechanical loss and optical absorption provide the greatest sensitivity limits, and thus the most significant opportunity to improve performance.

The majority of coating research is focused on understanding and reducing the thermal noise arising from the HR coatings because it is a leading noise source in Advanced LIGO's central frequency range and a major design hurdle for all future detectors. However the anti-reflective (AR) coatings also present some challenges such as optical absorption.

Coating research encompasses the following areas:

- **Materials Investigations:** Measurements of the physical characteristics (Mechanical loss, Thermo-optic loss, Young's modulus, optical absorption etc) for coating material candidates for A+ LIGO and LIGO Voyager. This research includes investigations of how these characteristics change through processing, composition, temperature, etc.
- **Structural Measurements and Modeling:** Research to determine the coating structure by measuring or modeling the distribution of atoms and bond configurations, which can then be used to identify correlations to mechanical loss and other material properties. These investigations aim to understand the mechanisms of coating mechanical loss, and identify ways of reducing its effect through a materials-by-design process.
- **Direct Thermal Noise Measurements:** Use of precision interferometry to directly measure thermal noise from coatings made from new materials, techniques, designs, etc. to verify predictions coming out of material investigations. This is an important step to both test predictions and test coating performance before use in actual gravitational wave detectors.

**Mirror substrate research** Mirror substrate research encompasses the following broad areas:



- Substrate material properties:** The mirror substrates for gravitational wave detectors must meet many requirements including low mechanical loss, low optical absorption, low scatter losses and appropriate figure. While previous experience with fused silica substrates is directly relevant for A+ LIGO, the proposed change to silicon substrates for LIGO Voyager requires significantly more research and development. The proposed change from silica to silicon substrates for Voyager is largely driven by the increase in mechanical loss of fused silica at cryogenic temperatures. In addition, silicon mirrors will experience less thermal lensing for the same laser power, potentially allowing the use of higher powers in the arm cavities with silicon mirrors. Silicon is not transparent at 1064 nm, so a change in interferometer laser wavelength, possibly to 1550 nm or higher, would also be required. Low optical absorption is important in a cryogenically-cooled mirror to minimize the heat load and allow the required operating temperature to be maintained. In addition, other optical effects in silicon need to be considered, including two-photon absorption and noise associated with both intrinsic and generated free carriers.
- Parametric instabilities:** The build-up of parametric instabilities in the arm cavities related to the high laser power levels are a potential problem from high optical power in aLIGO and beyond, with parametric instability already having been observed in the Livingston aLIGO detector. These undesirable effects result from exchange of energy between light stored in cavities and acoustic modes of the mirror which define the cavities, and can result in high excitation of particular acoustic modes of the mirrors, leading to control problems and in extreme cases loss of lock of the interferometer. Research focuses on understanding parametric instabilities and developing methods of reducing the effects.
- Charging:** There are various mechanisms by which the mirrors in an interferometer can become charged, and interaction between charges on the mirror and charges on nearby surfaces can generate force noise on the mirror. In addition, charging may lead to interferometer control problems through interactions with the electrostatic drives. Charging research focuses on measurements of charge noise and identification and testing of methods for charge mitigation.

## 5.1 Advanced LIGO +

### 5.1.1 Mirror Coating Research for A+ LIGO

Coating thermal noise is related to several properties of the coating including the mechanical loss, the Young's modulus and the thickness, and to the radius of the interferometer laser beam and the temperature of the mirror. For A+ LIGO, a number of different techniques are expected to be capable of providing some reduction in coating thermal noise, including using a larger beam radius, optimizing coating designs (including composition, layer thicknesses and possible exploitation of differences in the loss angle associated with shear and bulk motion) and possible reductions in mechanical loss via optimizing doping, annealing or via the use of promising alternative coating materials. It seems likely that a combination of a number of these techniques may be used to produce the best possible coating for the A+ LIGO detectors.

### 5.1.2 Coating Research: aLIGO Coatings

It is important to ensure that the thermal noise and optical performance of the aLIGO mirror coatings is fully understood, in order to enable an accurate baseline for the development of coatings for A+ LIGO. Understanding the variations in coating mechanical loss measured on different substrates, and how this relates to direct thermal noise measurements on test mirrors and to the magnitude of thermal noise expected in the aLIGO mirrors is a research priority. Similarly, a full understanding of the scattering characteristics of the aLIGO coatings is required to assist with fully characterizing scatter loss, and potential related noise sources, in the detectors.

### 5.1.3 Coating Research: Increased Laser Beam Radius

The use of a larger laser beam radius in the interferometer (along with the use of appropriately larger mirrors) is one of the most obvious methods for reducing coating thermal noise. However, the use of larger mirrors will have a significant impact on the mirror suspensions and seismic isolation systems, so detailed trade-off studies in collaboration with the suspensions working group are required. In addition, detailed discussions with coating vendors are required to determine how much development is required to enable coatings of the required uniformity to be deposited on larger optics.

### 5.1.4 Coating Research: Mechanical Loss

For room temperature mirrors at 1064 nm, the coatings with highest optical quality are ion beam sputtered (IBS) amorphous coatings, where the high index material is typically a metal oxide (ie.  $\text{Ta}_2\text{O}_5$  or  $\text{TiO}_2$ ) and the low index material is usually silica.

The thermal noise due to the mirror coatings can be calculated from the coating mechanical loss using the Levin method [131]. Two methods are currently employed to measure the mechanical loss. The first method measures the loss via ringdown of small coated cantilever samples over the temperature range of 0 – 300 K. These measurements map out the Debye loss peaks and determine the activation energy of the loss mechanism. This knowledge can then be used as a constraint in models of the coating structure (see Section 5.1.8). By accurately modeling the coating structure, one may be able to identify the microscopic causes of internal friction and design coatings with minimal loss. Thus, while temperature dependent coating loss measurements are also of interest for cryogenic coatings for LIGO Voyager, these measurements are also relevant to understanding loss mechanisms in room temperature coatings for A+ LIGO.

The second method measures the mechanical loss of coated thin silica disks at room temperature over a wide frequency range (typically 3 – 20 kHz). The frequency dependence allows separation of the bulk, surface and thermoelastic losses. Characterizing the mechanical loss at a wide range of frequencies and temperatures is valuable both as a search for new materials and to better understand the causes of thermal noise in amorphous thin film oxides.

The high index materials that have been explored include: tantala, titania, niobia, hafnia, and zirconia. Silicon nitride,  $\text{Si}_3\text{N}_4$ , also may be promising. Low index materials include

silica and alumina. As applied, the high index materials have losses ranging from  $3 - 7 \times 10^{-4}$ .

The loss in the silica coating layers, while several times less than the high-index layer, is not negligible. For coating geometries, the loss in silica is dominated by surface and stress-induced losses. The latter loss can be minimized through a slow annealing process. However, this process typically destroys a multilayer coating by either crystallization of the high index material or adhesion failure due to differential thermal expansion. Thus the loss in the silica layers will depend on the geometry and thermal history of the multilayer coating. The loss in alumina is dominated by thermoelastic loss and is not currently suitable for room temperature coatings.

The most commonly used model for loss in these glassy materials [132] assumes an asymmetric, double-well bond potential formed by two nearly degenerate bond states. In silica, for example, the O bond potential has an angular dependence that is described by nested, asymmetric potentials [133], where the activation energy depends on the atomic distribution. In this model, one may minimize the loss by either removing the potential's double-well (i.e., doping to change the bond potential) or by removing the asymmetry (i.e., annealing to allow the material to assume its lowest energy state).

**Doping** Doping can reduce the mechanical loss by favorably changing the bond potential and by stabilizing the matrix against crystallization. For example, doping the tantala layers with titania has been shown to reduce mechanical loss by about 40%. This reduction allowed titania-doped tantala/silica to be selected for the Advanced LIGO coatings. Silica-doped titania has shown promise for reduced thermal noise, and was considered a fallback coating for Advanced LIGO. A ternary alloy of titania/tantala/silica as the high index material may allow for benefits from each material. Silica, while effective at stabilizing high-index materials and thus allowing for higher annealing temperatures, also reduces the index of refraction, thus requiring thicker coating layers. Effective medium theory may be of use in modeling alloys and doping, and comparing Q measurements to predictions of effective medium theory are an important next step in verifying this approach.

The use of dopants in tantala, titania, hafnia and zirconia are being explored as a means of understanding and reducing the mechanical dissipation. While it is crucial that any new material also satisfy the stringent optical requirements of LIGO interferometers, there are advantages of initially just pursuing lower mechanical loss.

**Shear and bulk loss angles** Amorphous materials, including fused silica, tantala, titania doped tantala, and other oxides under consideration as coating materials have two independent elastic moduli, and thus two independent loss angles. Recent work at Caltech by Yanbei Chen's group suggests that it should be possible to design coatings with reduced Brownian thermal noise by having the thermal noise generated by different loss angles partially cancel out. Whether this is practical or significant depends on the values of the loss angles for the coating materials, which can only be determined experimentally. Measurement of torsional and bulk mechanical losses of coating materials will provide necessary input for these designs.

**Annealing** Ion beam sputtering forms coatings with high compressive stress, and comparatively high mechanical loss and absorption. Annealing these coatings will reduce the stress and correspondingly reduce the absorption and mechanical loss. Typically the annealing temperature is limited by either the crystallization temperature for the high index material or shear failure from the differential expansion of the coating layers. If these failure mechanisms can be avoided, annealing could potentially provide significant improvements in mechanical loss and absorption.

The mechanical loss in silica coatings have be reduce by about  $8\times$  through high temperature annealing. In addition a zirconia/silica coating has been demonstrated to survive a  $1000^\circ\text{C}$  annealing without shear failure. Annealing has also been shown to reduce the absorption in zirconia coatings by a factor of several. Work is currently underway to use doping to stabilize the coatings against crystallization, which leads to unacceptable scatter. Silica doping suppresses crystallite growth during annealing in titania [134], hafnia and zirconia [135]. Coating designs to withstand high temperature annealing could see significant improvements in their absorption and their mechanical loss. Planar layered composites consisting of alternating nm-scale Titania (or hafnia) and silica (or Alumina) layers can be annealed to much higher temperatures compared to (un-doped) materials [136, 137]. Recent results have suggested that heat-treatment of zirconia-doped tantala coatings can significantly reduce the mechanical loss, and heat-treatment at temperatures up to  $800^\circ\text{C}$  is possible without crystallization occurring. Further measurements and repeat coating runs are required to fully characterize this material. It is interesting to note that zirconia was also identified as an interesting dopant by atomic modeling work carried out within the LSC. This work showed that titania doping increased the flexibility of ring-structures found in the coating, and suggested that zirconia doping would increase the flexibility further, possibly resulting in lower mechanical loss [G1300379]. The strategy of using a doped material consisting of the correct proportions of two high-index materials to stabilize against crystallization appears promising and should be further investigated with other material combinations.

**Elevated temperature deposition** Deposition of amorphous silicon coatings onto heated substrates can reduce the mechanical loss by orders of magnitude [137], possibly due to the material forming a more ordered amorphous state referred to as an ‘ideal glass’. Evidence suggests that a larger loss reduction is possible with elevated temperature deposition than with post-deposition annealing, possibly due to the coating atoms have more freedom to move during the deposition process. Since post-deposition annealing is already known to reduce the mechanical loss of oxide coatings such as silica and tantala, deposition at elevated temperatures may also be a promising technique for significantly reducing the loss of these materials. Alternative methods of providing more energy to the coating atoms as they are deposited on the substrate may also be of interest, for example the use of ion-assisted deposition where a secondary ion-source is used to bombard the substrate with an ion beam during coating deposition.

**Interface Effects** Early experiments on loss in Initial LIGO’s tantala/silica mirror coatings [138] demonstrated that the loss was primarily due to the tantala layers. No significant loss could be attributed to the coating interfaces. Measurements conducted at LMA on

the Advanced LIGO coating initially indicated some excess loss associated with the coating interfaces, but that conclusion was not borne out in subsequent measurements. Nevertheless, for any proposed change in coating material a study of interface losses should be conducted. [139].

### 5.1.5 Coating Research: Optical Properties

**Optical Absorption** Absorption of light in the coatings will result in thermoelastic distortion of the optics and will ultimately limit the circulating light power in the interferometer. When coupled to the bulk absorption in the input test masses, this leads to significant surface deformation of the test masses and bulk thermal lensing in the input test masses. Coating absorptions as low as 0.3 ppm have been reported in undoped silica/tantala coatings, while titania-doped tantala and silica-doped titania coatings have been shown to have absorption at or below 0.5 ppm. The coatings used in aLIGO made of silica/titania doped tantala have an absorption of  $< 0.4$  ppm [140]. Further improvements beyond this level will make thermal compensation easier for A+ LIGO, and detailed studies of absorption are essential for any coating materials considered for use in current or future detectors.

Studies aimed at understanding and improving coating mechanical loss may involve working with coatings with relatively high absorption during a research phase e.g. to understand why a particular dopant affects the mechanical loss. Further research into the absorption of anti-reflection coatings is also required, as these coatings consistently have a higher absorption (up to 10 ppm) than high-reflectivity coatings (typically below 1 ppm).

**Optical Loss from Scattering** In order to maintain the highest optical power in future detectors, it is important to minimize the optical scatter. Scatterometer measurements should be conducted for proposed coatings and new coating materials. Studies of the dependence of scatter on coating materials and manufacturing are important in determining the lowest possible scatter.

Realizing more sophisticated quantum non-demolition (QND) topologies also requires extreme low-loss optical systems as is explained in section 4.2.1 for the case of filter cavities. One of the important sources of optical loss is scattering from mirror-surface aberrations. These are traditionally investigated by measuring the angular distribution of scattered light (i.e. measurements of the bidirectional scattering distribution function (BSDF)), or scanning the surface with lasers and integrating the scattered light in spheres. As much as these measurements are important to link scattering from mirrors with losses in optical systems like cavities, they do not give direct information about the cause of scattering.

Scatter loss in (future) optical systems with sizes up to a few hundred of meters will likely be dominated by point-defect scattering as the quality of substrate polishing has advanced to a level that makes residual surface-roughness scatter loss negligible in most cases. These conclusions are based on numerous simulations and partially on scattering measurements in first-generation GW detectors. Even though it is believed that very low-loss systems can be realized in the near future with scatter loss around 10 ppm per mirror, the question is how much further loss can be decreased. Loss estimates play a major role when deciding between the various candidate QND configurations for future GW detectors. Whereas input

filters (see section 4.4.1) are relatively robust against optical loss, output filters (see section 4.4.2) that can potentially eliminate all back-action quantum noise are known to be highly susceptible to loss. A few ppm loss per mirror typically destroys the entire advantage that output filters have over input filters (eventually making them even worse in performance). Similar problems are encountered with alternative QND schemes.

Assuming that point-like defects residing in the mirror coatings are the dominant source of scatter loss, one has to investigate individual defects for their material compositions, morphologies, and structures. The answers can be used to understand the origin of the defects with the goal to improve the coating process. Various analysis methods are available. Defect morphology can be studied optically or with force microscopy depending on defect size. Defect materials can be investigated spectroscopically. The analyses should progress from larger to smaller defects since the larger defects dominate the point-defect scatter even if they are significantly less numerous.

### 5.1.6 Coating Research: Other Coating Properties

**Thermo-optic Noise** Coating Thermo-Optic (TO) noise is the apparent motion of the surface of a mirror due to stochastic temperature fluctuations. These temperature fluctuations drive two separate but correlated effects, commonly known as coating Thermo-Elastic (TE) and coating Thermo-Refractive (TR) noise [141, 142]. The power spectral density of these temperature fluctuations has been calculated in various places [141, 143].

The first effect (TE) is physical motion of the surface of the mirror due to the thermal expansion of the coating. This effect requires knowledge of the effective Coefficient of Thermal Expansion (CTE) for the coating materials. We say “effective” because the coatings are attached to large substrates, therefore not free to expand laterally; Poisson ratio effects and substrate behavior need to be accounted for. They are also under great stress so that any change in the Youngs modulus with temperature could conceivably couple into out-of-plane expansion of the layers.

The second effect (TR) is apparent motion of the surface of the mirror as measured by a laser beam, due to changes in the mirror’s complex reflection coefficient. Since the mirrors used are generally high reflectors, they magnitude of the reflectivity doesn’t change when the temperature changes (varying the optical path length of the individual layers and thus the superposition of the reflected light field), but the phase picked up on reflection does change to first order in the temperature change. For historical reasons this effect is referred to as TR, but physical expansion of the layers also plays a role here.

It is expected that these two effects will partially cancel [37], and there is evidence that this is indeed the case [144, 145]. In [37], it was argued that with nominal values the cancellation would put the level of TO noise roughly an order of magnitude below coating Brownian noise. No noise floor data out of the LIGO interferometers has contradicted this estimate.

However, the relevant thermal / thermo-optic parameters involved (specifically the CTE and  $dn/dT$  for the thin-film high- and low-index materials in the coating) are known to differ from bulk values, and measurements of these values exhibit wide variation (up to an order of magnitude, see e.g. [146–150]). The situation is complicated by the fact that the parameters

could conceivably depend on deposition technology, doping, and layer structure. This is particularly unfortunate for two reasons.

First, it makes predicting this noise level challenging. In order for TE and TR effects to cancel so that the sum is of order 10% of either effect alone, it is necessary that the magnitudes of the TE and TR effects are within 10% of each other to begin with. Since these parameters are not known at the 10% level, as we move forward it will be important to check individual coating TO responses. We currently have capabilities for making such a measurement in the collaboration (using the technique described in [144], but are interested in other methods of making the measurement as well.

Second, it makes design of future coatings with greater degrees of cancellation (or design of coatings to jointly minimize say TO noise and Brownian noise) effectively impossible until we have better measurements of the individual parameters themselves. A method for extracting individual parameters from measurements on a set of differing coating structures was shown in a proof-of-principle manner in [144], but we still need trusted measurements for these parameters as they will appear in our coatings. To the extent possible, coating material, chemistry, deposition process, substrate, and coating layer structure should be held as close as possible to the materials, processes, and designs used in the LIGO mirror coatings until these various factors impacts on the TO parameters of the coating materials can be teased out. Studies of these effects (both experimental and theoretical) that help us understand their impacts on coating TO parameters are also of serious interest as they can then help us in the selection and design of future coatings.

**Young’s Modulus and Stress** The Young’s modulus of a coating is required both for the analysis of mechanical loss measurements and for calculations of the level of coating thermal noise. It is therefore important to obtain accurate values of Young’s modulus for every coating, and post-deposition coating treatment, studied. Residual stress in coatings is likely to be an important property, and there is interesting evidence suggesting that stress can alter mechanical loss of coatings, particularly in silicon nitride. Therefore studies of the effects of residual stress on the loss and of methods of altering the stress in particular coatings are of interest. The use of several measurement techniques can be beneficial in these studies, as each technique has different systematic errors and, for example, different sensitivity to the properties of the coating substrate material. In addition to measuring these properties at room temperature, where possible the capability to measure the temperature dependence of these properties should be developed.

A Young’s modulus and dissipation measurement method with sub nanometer spatial and depth resolution developed by Konrad Samwer shows that the Young’ modulus in glasses has a position dependent spread as wide as 30% (and the local loss factor is also poorly defined), that the spread is reduced with annealing, while crystals have constant Young’s modulus everywhere. It has also been shown that fused silica, which is the glass with the lowest known mechanical loss, has a substantially narrower Young’s modulus spread than other glasses. The method can either explore small shallow volumes, or wider and deeper volumes, up to several hundreds of atomic spacings in dimension. The capability of this method to scan the Young’s modulus with sub nanometric resolution offers a new way to explore the uniformity of our coatings, as a function of annealing, and perhaps shine some

light on some loss mechanisms.

**Uniformity** It has also proved challenging to maintain coating thickness uniformity across the large face of Advanced LIGO optics. Non-uniformity in thicknesses leads to non-uniformity of transmission and scatter of light out of the cavity mode to other optical modes. This also leads to limits on optical power and squeezing, as the higher spatial frequency scatter discussed in the previous paragraph. Even larger optics are possible in future detectors, making maintaining coating uniformity even more challenging and thus an important research topic. The limitation on obtainable uniformity can come from metrology limitations, so improved metrology is an important research direction. An additional research direction is to explore corrective coatings, which place additional coating material onto a coated optic after uniformity measurements have been made.

### 5.1.7 Coating Research: Coating Deposition Parameters

Variations in the loss of nominally identical coatings from different vendors have been observed, suggesting that the precise deposition parameters may be important in determining the loss. Thus more detailed measurements of the effects of parameters such as ion energy, sputtering ion, oxygen pressure and thermal treatment may be valuable. While ion beam sputtering produces the lowest optical loss coatings, the mechanical loss of coatings deposited by other techniques has not been extensively studied. Studies of coatings deposited by different techniques (e.g. magnetron sputtering, e-beam evaporation, atomic layer deposition) may enhance understanding of the relationship between loss and structure in these materials.

Coating layer thickness may also prove an important variable in determining amorphous material properties, especially mechanical loss. There are suggestions that making traditional quarter-wavelength coating layers from a sub-structure of much thinner layers of two materials may enable improvements in mechanical loss and may allow higher treatment temperatures to be realized without crystallization occurring [P1800102] [151]. If this effect can be confirmed and understood, it may be possible to design high reflective coatings with much lower thermal noise using known materials like tantalum and silica (see **nm-Layered Composites** in Sec. 5.1.10).

### 5.1.8 Coating Research: Structural Studies

The mechanical losses are fundamentally connected to the atomic structure of the amorphous film, and therefore experimental and theoretical tools for characterizing the atomic properties form an important area of research. Until recently the experimental scattering methods have primarily yielded information on short-range order ( $< 0.5$  nm). Recent theoretical insights show that structural motifs on several nm spatial scales are responsible for the elastic losses at frequencies and temperatures of interest here. Thus, theoretical and experimental methods for characterization of medium range order ( $0.5 \sim 5$  nm) are essential for understanding and predicting elastic losses in these films.



**Experimental Structure Characterization** Experimental structure determinations are based on x-ray or electron scattering methods, both of which have been employed. The data collected can be converted to statistical averages of local atomic environments, often described in terms of pair distribution functions (PDFs). Electron PDF (e-PDF) and extended x-ray absorption fine structure (EXAFS) have been used to gain an understanding of the short range order in doped and heat-treated tantala [152–155]. In more recent work, grazing incidence pair distribution functions (GI-PDF) has been used to reach higher resolution on unprocessed titania and zirconia doped tantala thin films [156]. Fluctuation electron microscopy (FEM), which uses a focused electron beam to enable characterization of the variance in local environments, is a measurement that is more sensitive to medium range order than conventional e-PDFs [157].

**Atomic Modeling** Molecular dynamics methods allow numerical simulation of the atomic structures of amorphous solids obtained by cooling the corresponding liquid phase [158]. The calculated energy landscape obtained through suitable processing of the structures computed from multiple such numerical experiments allows calculation of the elastic loss. In addition to comparing the calculated losses to experimental data, the simulations can also be tested by comparison with the PDFs of these model structures with those obtained from experiment. For titania-doped tantala, these calculated losses correctly captured the trends vs Ti doping and the absolute values at cryogenic temperatures, but were less accurate for room temperature data. Similarly, the simulated PDFs were close to experimental data for short range order, but were less successful for capturing the medium range order [159].

To address these issues, reverse Monte Carlo (RMC) methods are being used to compute the atomic structure from experimental scattering data. By including the measured variance and/or the high resolution data from our novel GI-PDF measurements as the constraints on the RMC, one can expect to obtain a better representation of the medium range order than was previously possible. It will then be possible to use these improved structures as a basis for modeling the effects of dopants on the structure and elastic loss of the films. Empirically, it has been seen that doped films often have lower losses than pure oxides, e.g. Ti in tantala [36]. Thus, an important goal of this modeling effort is to predict promising material/dopant combinations to synthesize and explore experimentally. To implement this plan, it will be necessary to have appropriate inter-atomic potentials for the coating materials of interest. Currently we have inter-atomic potentials for silica, tantala, titania and zirconia; we plan to develop atomic potentials for other oxides, including hafnia and alumina to enable high throughput simulations of mixed oxides.

**Deposition Simulation** Another key theoretical approach is the direct MD simulation of vapor deposition processes. Such a tool would be invaluable to reduce the large parameter space of substrate temperatures, deposition rates, materials and dopants to a promising subspace that could be synthesized and characterized efficiently. Recent reports of such simulations Refs. [160] and [161] using simple model potentials showed promise for this approach. Such simulations will be implemented using the accurate atomic potentials, to obtain realistic predictions to guide experimental work.

### 5.1.9 Coating Research: Direct Thermal Noise Measurements

Direct measurements of thermal noise are of interest to compare with the predictions obtained from mechanical loss measurements and to test the improving theories of coating thermal noise [162].

**Fixed Spacer Cavities** Coating thermal noise measurements are now being carried out using fixed spacer cavities, which are limited by coating thermal noise over a wide frequency band and are significantly easier to operate than a suspended mirror system. These systems will use standard size (1 inch diameter) substrates and can provide a convenient test-bed for the development of low thermal noise optics.

**AEI 10 meter prototype** The 10 meter prototype interferometer at the Albert Einstein Institute in Hannover Germany will be able to directly measure coating thermal noise. In addition to testing Khalili cavities as a means of reducing coating thermal noise, it will have the ability to change the size of the laser spot on the mirrors. This will allow for a direct test of spot size dependence, which is an important driver of the desire for larger optics in future detectors.

### 5.1.10 Coating Research: Coating Design

**nm-Layered Composites** Planar layered composites consisting of nm-scale alternating films of Titania and Silica are increasingly stable against crystallization, as the (Titania) layers thickness is reduced [151]. These composites behave as homogeneous materials as regards their optical and viscoelastic properties, for which simple and accurate modeling is available. Crystallization inhibition up to very high annealing temperature has been also observed in nm-layered Hafnia-Alumina composites [137].

**Optimized Coatings** Since the thermal noise in the coatings typically scales as the total thickness of the more lossy material, reducing this thickness while maintaining the optical properties will reduce thermal noise. Constrained numerical optimization codes have been shown to produce high reflectivity coatings while reducing the volume of high index materials by as much as 20%. Thermo-optic noise from thermoelastic and thermorefractive effects is included in this optimization. The mechanical loss of the low index (silica) material takes on a larger role for thickness optimized coatings, as optimization typically makes the high index (titania- tantala) contribution equal to the low index. Such an optimized design was proposed for use in Advanced LIGO [163, 164]. Greater understanding of mechanical loss in thin film silica and/or other low index materials is crucial to exploiting the full potential of this optimization. Thickness optimization should be generalized to incorporate the dopant concentrations (which affect both the optical and viscoelastic properties of the materials) in the parameters to be optimized, in the perspective of designing minimal noise coatings based on doped (or nm-layered) mixtures.

Designing coatings that take advantage of different loss angles for bulk and torsional motion will be important once numerical values are found. This optimization will need to be done

while including thickness effects on Brownian thermal noise as well as thermo-optic noise.

**Multi-material coatings** The use of multi-material coatings to take advantage of the properties of different materials has been proposed [165, 166]. In particular, it may be possible to exploit the fact that most of the incident light intensity is reflected by the first few bi-layers of a coating, potentially allowing coating materials with higher optical absorption, but lower mechanical loss, to be used in the lower layers of a coating stack without significantly increasing the total absorption of the coating stack. It should be noted that multi-material coating designs, coupled with the low optical absorption observed in amorphous silicon coatings deposited using a novel ion-beam sputtering process at the University of the West of Scotland, may potentially allow the use of amorphous silicon layers in mirror coatings for use at 1064 nm for A+. Amorphous silicon is an attractive coating material due to its low mechanical loss and high refractive index: see section 5.3.1 for more details.

#### 5.1.11 Mirror substrate research

**Fused Silica** Experiments to measure mechanical loss in silica versus annealing parameters, including ramp down and dwell times have led to improvements in the substrate thermal noise. In order for the fused silica thermal noise to pose a problem in the future, the thermal noise of the coating would have to be reduced by more than an order of magnitude. This makes silica substrate mechanical loss studies a lower priority than coating mechanical loss.

#### 5.1.12 Mirror Substrate Research: Parametric Instabilities

The build-up of parametric instabilities in the arm cavities due to the high laser power levels are a potential problem in Advanced LIGO and beyond. These undesirable effects result from the exchange of energy between the light stored in the optical cavities and the acoustic modes of the mirror which define the cavities. At high optical powers, the radiation pressure force of scattered high order optical cavity modes can couple strongly to the mechanical modes of the test masses, resulting in a parametric instability. High excitation of the mirror's acoustic modes can result in difficulties in the controls engineering and at very high amplitudes can lead to lock loss. Unfortunately, the requirements for high sensitivity are commensurate with the conditions under which parametric instability occurs. These include high optical power and low mechanical loss materials in the mirrors.

Using finite element methods, it is possible to start developing a quantitative understanding of this problem by modeling the modes and parametric gain for different test mass configurations, as well as investigate methods for mitigating the instabilities. In order to make a realistic estimate for the parametric gain, it is necessary to also include the full field calculations of the dual-recycled interferometer [167].

Measurements on suspended test masses are needed to obtain realistic, as-built, test mass  $Q$  values to establish the net gain for the instabilities. Adding tuned mass dampers to the barrel of the test masses (LASTI-MIT) and/or using feedback to the electro-static drive also show promise for controlling parametric instability. In addition, spatially-resolved radiation pressure feedback on the mirror surfaces is being contemplated (Gingin-UWA). Outstanding

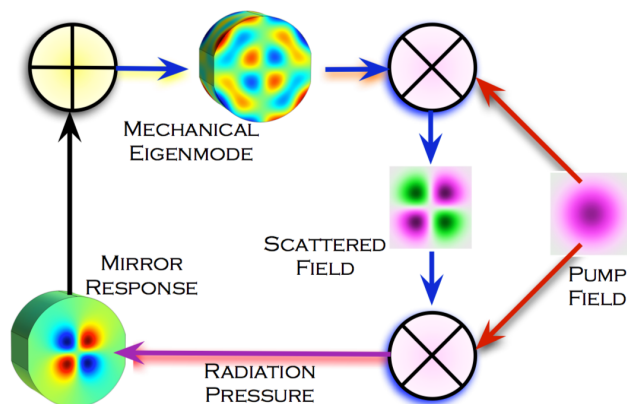


Figure 25: Schematic diagram of the parametric instability mechanism [15].

questions include whether these approaches are compatible with high sensitivity, including low shot noise, low thermal noise, and realizable controls.

Small scale suspended test mass experiments are also underway to study parametric instability. The larger scale experiments will be useful in testing noise performance of various parametric instability suppression schemes, the small scale experiments will be useful to validate the theory of parametric instability and for proof-of-principle test of suppression schemes. The small scale experiments using millimeter scale test masses in a tabletop configuration using coupled cavities and/or specially designed near-concentric cavities can be done using a standard NPRO laser and standard 1 inch optics. The small scale resonator experiment at UWA has reached the threshold of parametric instability. These experiments are valuable but because of its applicability to Advanced LIGO and A+ LIGO, the emphasis of parametric instability research should be on suppression schemes.

Modeling [168–170] already showed that 3MI signal gain had extreme sensitivity to metrology errors. This result carried the message that 3MI monitoring could be very sensitive, but its usefulness was not recognized because the main concern at the time was to suppress the few interactions for which the parametric gain could exceed unity as these represent parametric instability. Using selected modes with gain between  $10e-1$  and  $10e-3$  that have been shown experimentally to be easily observed with excellent signal to noise ratio, we would be able to monitor the interferometer global state variations. Because of the complexity of the information that the 3MI provides, further efforts in modeling and experimental research is required to make the 3MI to be a practical monitoring tool.

### 5.1.13 Mirror Substrate Research: Charging on the Optics

Surface charge may build up on the test masses through a variety of mechanisms, including contact with dust (particularly during pump down) and/or the earthquake limit stops, removal of First Contact used to keep the optic clean during transport and handling, as well as cosmic ray showers.

One noise mechanism is static charge distribution on either the optic or the earthquake stop

will couple motion of the earthquake stop into forces on the optic [171]. Another mechanism is the noise caused by time-varying charge distributions on the optic (or the earthquake stop) resulting in time-varying forces on the optic. Gaussian noise from this mechanism can be described by a Markov process [172]. The result depends on the magnitude of the deposited charge and the correlation time of the deposited charge, with a smaller actuating noise for correlation times far from the reciprocal of the frequency at which the noise is being measured. These correlation times are being measured using scanning Kelvin probes operated in vacuum which measure the magnitude and distribution of surface charges and their rate of motion across a sample.

Charge may also interact with the electro-static drive causing noise or reduced effectiveness of the drive. Modeling has been started to study this, and experimental work at LASTI has begun to better understand the role of charge with the electro-static drive. There have also been two experimental verifications of a charging contribution from dielectric polarization of the fused silica [173, 174]. Further experimental and theoretical work is planned on polarization noise. Calculations have also been carried out to estimate the force noise that might be expected from Coulomb interactions between charge accumulations on the test mass and various components in the suspension system. The earthquake stops being the closest to the test mass surfaces are of greatest concern for most issues with charge on the optic [171].

Ongoing work will focus on different cleaning and handling methods (including ways of applying and removing First Contact), discharging the optics using ionized gas, and conductive coatings. Understanding what sensitivity limits might come from charging is crucial for A+.

**Discharging using ionized gas** The primary solution to mitigating charge in the advanced detectors involves blowing nitrogen gas across needles at 4kV AC, which ionizes the gas externally to the vacuum chamber [175–177]. The nitrogen ions, which comprise both polarities, travel through an aperture and into the vacuum tank. This is currently installed and used in aLIGO, and will likely remain in A+.

**Conductive coatings** Developing and testing finite conductivity coatings is also an important area of research. Here, the influence of charge on the coating surface will be reduced by having a lightly conductive coating under the dielectric stack, which will support a compensating image charge plane. This will mitigate the effect of surface charges interacting with nearby support structures, particularly the earthquake stops. The conductive coating can then be effectively grounded by UV photoemission conduction, between the optic and support structures. Work is ongoing on conductive ion beam sputtered layers composed of alumina doped zinc oxide (AZO), and measuring the relationship between electrical conductivity, optical absorption, and mechanical loss. In addition, atomic layer deposition (ALD) of zinc oxide is also under investigation, and show promise for being conductive with extremely thin (10 nm) layers. To discharge the optics with conductive coatings, a "UV electron photoemission wireless conduction" system has been developed, and tests verify that it can ground the test mass to less than a 10 V potential. The UV source will consist of UV GaAs LED's and photoelectrons will be generated – implemented on the barrels of the optics, or possibly the earthquake stops.

**Charge noise measurements** It would also be useful to directly measure noise from charging, to confirm both the Weiss Markov-process noise model and the parameters found from the Kelvin probe work. Torsion balances, which have been used for laboratory gravity experiments and to test noise models of LISA, offer another possibility to verify Markov noise from charges. Torsion balances are well suited to this since they reach their highest sensitivity at frequencies where Markov charge noise is expected to be large. For charging studies, the torsion pendulum will need to be made entirely of an insulator, likely fused silica, which is a departure from previous experience. The LSC group at the University of Washington, which has experience with torsion pendulums through LISA and other research programs, has performed studies of charging noise [G1000367]. Additional torsion balance experiments are also being developed at University of Glasgow [G1100714] and Moscow State University. In Glasgow, a torsion bob comprising fused silica discs is being utilized to study charge motion on the surface of fused silica and the level of charge deposition when fused silica surfaces come into contact. A Kelvin probe located within the vacuum tank also allows the correlation time to be measured. Initial results suggest that the Weiss theory does give an accurate estimation to the level of charge noise [G1300213, P1300078]. Further tests are planned with zinc oxide coatings which have the possibility of providing well defined characteristic correlation times, thus allowing the level of charge noise to be varied in a reproducible way. Research at Moscow State University is focusing on the exploration of noise associated with dielectric polarization induced when an electrostatic drive (ESD) is operated near a fused silica test mass [G1200166]. The experimental setup includes a monolithic torsion oscillator with frequency of 63 Hz fabricated entirely from fused silica.

## 5.2 Advanced LIGO Upgrades Beyond A+

Beyond the A+ detectors, it may be possible to develop improved technologies that would enable increased performance at room temperature. For the detector optics, this means further reductions in thermal noise.

Larger fused silica test masses in the range of  $\sim 80 - 200$  kg could be employed as a way to reduce thermal noise. This would require improved substrate surface figure error over the area of the larger test masses, while managing the residual substrate fixed lens and elastic distortion of the figure error when the mirror is suspended. There would also need to be improved coating thickness uniformity over the larger area of the larger test masses.

Further improvements to amorphous coatings would aim for beyond a factor two decrease in thermal noise targeted by A+. Similar research techniques to those described in section 5.1 would enable the discovery of new coating materials and approaches. In addition to amorphous coatings, development of large area AlGaAs/GaAs crystalline coatings and the transfer process necessary to apply them to silica substrates may also be possible, although is currently thought to be cost prohibitive [178] (see Section 5.3.1, which describes development for Voyager crystalline coatings, but is also applicable to room temperature coatings).

### 5.3 LIGO Voyager

The proposed use of silicon mirrors operating at 120 K and operating wavelength of 1550 nm or possibly 2000 nm will require further research into silicon as an optical material, continued development of suitable low thermal noise coatings and measurement of a wide range of properties of these coatings at low temperature. Many of the research areas and techniques mentioned in the previous section are relevant to LIGO Voyager, but research on different materials and with many of the measurements carried out at low temperature.

#### 5.3.1 Mirror coating research: Mechanical Loss

The thermal noise of a mechanical system depends on the temperature and mechanical dissipation. For many optical coating materials, the mechanical loss increases at cryogenic temperatures [179, 180]. Thus using coatings designed for room temperature at cryogenic temperatures will see less improvement in the thermal noise than one might naïvely assume. Thus the shift to cryogenic operation will require the development of new mirror coatings.

**Crystalline Coatings** Single-crystalline coatings grown by molecular beam epitaxy are of significant interest as a possible alternative to current amorphous coatings. This is likely to be of particular relevance for low-temperature detectors, as the cryogenic dissipation peaks observed in current silica and tantala coatings are thought to be related to the amorphous structure of the materials [159]. There are multiple crystalline coatings under investigation within the LSC; among them are aluminum gallium arsenide, AlGaAs, and aluminum gallium phosphide, AlGaP. General factors which require investigation for crystalline coatings include adhesion to substrates at cryogenic temperatures, light scattering, and fabrication on curved mirror substrates.

Aluminum Gallium Arsenide (GaAs:AlGaAs) coatings [38] have been studied as free-standing micromechanical resonators for use in quantum optomechanical experiments, and have been shown to have very low mechanical losses (as low as  $4.5 \times 10^{-6}$ ) at cryogenic temperature [181]. However, these coatings are grown on GaAs substrates, and must be transferred and bonded onto appropriate mirror substrates for use in gravitational wave detectors. Current ongoing experiments with AEI Hannover are exploring larger area coatings at 48 mm in diameter, with larger coatings over 200 mm diameter planned for the future [178]. The application of this technique to curved mirrors will require development effort, and it is essential to study these coatings after transfer to appropriate substrates to evaluate any additional mechanical loss and light scatter which may be associated with the bonding process.

Aluminum Gallium Phosphide (GaP:AlGaP) coatings are lattice-matched to silicon, allowing a reflective coating to be grown directly on to a silicon mirror substrate, eliminating the need for coating transfer and bonding [182]. Initial measurements of the mechanical loss at room temperature were limited by thermoelastic damping in the silicon substrate: however, these measurements did allow an upper limit of approximately  $< 2 \times 10^{-4}$  to be placed on the coating loss at room temperature [183]. Continuing to characterize the loss and optical properties of these coatings at cryogenic temperatures is of high priority. These coatings are more typically grown on GaP substrates, and thus further development of the techniques for growing these coatings on silicon substrates will be necessary [182].

**Amorphous Coatings** In addition to the proposed design of Voyager operating at 120 K planned future detectors such as the Einstein Telescope in Europe and KAGRA in Japan will operate at cryogenic temperature to reduce thermal noise. It is therefore essential to fully characterize the performance of coating materials at cryogenic temperatures. In particular, the mechanical loss can be a strong function of temperature and low temperature loss peaks have been observed in silica, tantala and titania-doped tantala coatings [179, 180, 184], there is still some benefit to operating these coatings at cryogenic temperature. For example cooling to 20 K would provide a reduction in coating thermal noise by a factor of about 2, rather than the factor of 4 improvement which would be expected if the coating loss was constant with temperature [179].

Cryogenic mechanical loss measurements of coating materials are a valuable tool for exploring the microscopic processes responsible for energy dissipation. Identification and analysis of Debye-like loss peaks allows key parameters of the dissipation mechanisms to be calculated and, coupled with atomic modeling and structural measurements, may allow the association of loss peaks with particular types of atomic motion within the coating structure.

There has recently been significant progress in understanding the loss mechanisms in tantala coatings, with analysis of cryogenic loss peaks providing information about the dissipation mechanisms, and structural studies and atomic modeling revealing correlations between structure, doping level and loss. The level of loss in tantala below 100 K is strongly dependent on heat-treatment and doping level, and continued studies to optimize coating composition and post-deposition annealing may yield further improvements in coating thermal noise. The results are consistent with a model in which transitions of atoms between energetically stable positions are responsible for the loss, and suggest that the atomic structure of the coating is a key factor in determining the loss.

To enable further reductions in coating thermal noise there is an ongoing effort to identify coatings with a lower mechanical loss at cryogenic temperatures. A number of research paths are being pursued, including further improvement of current silica/tantala coatings, the use of alternative coating materials, particularly amorphous silicon high-index layers, tantala doped titania, silica doped hafnia and titania, nm-layered silica/hafnia and silica/titania composites.

For cryogenic coatings, amorphous silicon is a leading candidate for a high index material. Amorphous silicon (a-Si) coatings can have a particularly low mechanical loss, with recent measurements placing a conservative upper limit of  $5 \times 10^{-5}$  on the loss angle of ion-beam sputtered a-Si below 50 K [185]. In addition, the high refractive index of silicon allows thinner coatings with fewer layer pairs. The use of a silicon/silica coating could potentially reduce coating thermal noise by a factor of 2.4 at 20 K compared to a silica/tantala coating. However, the first measurements of optical absorption in these coatings suggest that significant efforts to understand and reduce the optical absorption may be required. However, the effect of the high absorption could be significantly reduced through the use of ‘multi-material’ coatings 5.3.7 (see Sec. 5.1.10), and the use of heat-treatment has been shown to be effective in significantly reducing the absorption of a-Si films. Recent work carried out within the LSC has shown that amorphous silicon films deposited using an ion-beam sputtering system with a novel electron cyclotron resonance ion source can have optical absorption a factor of up to 50 lower than other, commercially available ion-beam sputtered a-Si coatings [P1800148 - in



preparation]. Finally, the absorption of a-Si can be a factor of  $\sim 7$  lower at 2000 nm than at 1550 nm [186], so a move to a  $\sim 2000$  nm laser could also be necessary to enable the use of a-Si coatings.

It has been shown that infusing hydrogen in the silicon coating can significantly reduce the optical absorption, thought to be related to passivating dangling bonds. Studies to test this for ion-beam sputtered coatings and to evaluate the effect of hydrogenation on the optical absorption should therefore be carried out. However, the hydrogen can diffuse out of the silicon, reversing this effect. Nevertheless, the usefulness of dopants to improve the mechanical loss is clear. Recent work also suggests that deposition of a-Si films at elevated temperature (400°C) can result in significant reduction of the mechanical loss due to an increase in the structural order of the film. The effect of deposition at elevated temperatures on the optical absorption is therefore an important area of study.

Studies of different deposition methods for a-Si are also of great interest. In particular, chemical vapor deposition is a more mature technology for a-Si deposition than ion-beam sputtering (used, for example, for solar cells) and allows for relatively straightforward tuning of the deposition process including control of doping and stress.

Measurements of hafnia coatings indicate that, even when in a partially poly-crystalline form, this material has a lower loss than tantala at temperatures below 50 K. Preliminary results [187] seem to indicate that amorphous  $\text{TiO}_2$  may also be almost exempt from a cryogenic loss peak. Experience with tantala suggests that poly-crystalline structure may significantly increase the loss. One method of preventing crystallization of hafnia films is doping with silica and it has been shown that this does not significantly increase the loss at room temperature. Silica doping is also effective in stabilizing Titania against crystallization [134]. Layered nm-scale silica-titania (and alumina-hafnia) can also be annealed at high temperatures. Low temperature loss measurements of silica-doped hafnia coatings are underway. Cryogenic loss measurements on tantala and Silica doped Titania, and nm-layered hafnia/silica and Titania/Silica composites are also being planned.

There have been reports on the particularly low cryogenic mechanical loss of stressed amorphous silicon nitride films [188, 189]. The composition of this material is highly process-dependent and is possible to make both high and low refractive index  $\text{SiN}_x$  films by varying the composition, leading to the possibility of an entirely CVD-grown,  $\text{SiN}_x$  based HR coating.

Diamond-like carbon (DLC) coatings may also be of interest for further study, as there is evidence in the literature that the loss of this material is very low, with some films having a lower loss than amorphous silicon.

Studies of other possible alternative amorphous coating materials should continue, and where possible the choice of material (or treatment regime e.g. dopant, doping level, heat-treatment) will be informed by the results of structural measurements and modeling. While most of the effort to date has focused on developing alternative high-index coating materials for use at low temperature, it should be noted that silica coatings also have a cryogenic loss peak of a similar magnitude to that observed in tantala. Thus more studies of possible alternative low-index materials are required. In this connection, silica doped Hafnia (or nm-layered Hafnia-Silica composites) could be an interesting candidate low-index material, with a refractive index of  $\approx 1.5$ . This could be used with Tantala doped titania (or silica doped

titania, or nm-layered titania-silica composites) for the high-index material (with refractive index  $\approx 2.15$ ), to achieve a contrast comparable to that of the silica/tantala HR coatings presently in use, hopefully featuring much better cryogenic behavior.

**Shear and bulk loss angles** Understanding the mechanical loss angles associated with bulk and shear motion is of interest for any candidate coating for LIGO Voyager, and suitable finite element modeling will be required to support experimental investigations. Extension of this work to crystalline coatings, where separate loss angles are expected to be associated with each crystal axis, is required.

### 5.3.2 Coating Research: Optical Properties

**Optical absorption** Measurements of coating optical absorption at cryogenic temperatures will be necessary to provide the information necessary to select coatings for the proposed Voyager design. Just as with mechanical loss, the optical absorption can be strongly dependent on doping and annealing. These dependencies will need to be investigated for any proposed coating in order to minimize both mechanical loss and optical absorption. As noted above, aSi is one of the most promising low mechanical loss coating materials for use at low temperature, however, the absorption of aSi coatings is currently significantly higher than will be required for Voyager and further development is required.

**Optical Loss from Scattering** In order to maintain the highest optical power in future detectors, it is important to minimize the amount of scattered light. Scatterometer measurements should be conducted for proposed coatings and new coating materials. Studies of the dependence of scatter on coating materials and deposition parameters will be important in determining the lowest possible scattered light levels. Realizing more sophisticated quantum non-demolition (QND) topologies also requires extremely low-loss optical systems as is explained in section 4.2.1 for the case of filter cavities. One of the important sources of optical loss is light scattering from mirror-surface aberrations. These are traditionally investigated by measuring the angular distribution of scattered light (i.e. measurements of the bidirectional scattering distribution function (BSDF)), or scanning the surface with lasers and collecting the scattered light with integrating spheres. As much as these measurements are important to link scattering from mirrors with losses in optical cavities, they do not give direct information about the cause of scattering.

Assuming that point-like defects residing in the mirror coatings are the dominant source of scatter loss, one has to investigate individual defects for their material compositions, morphologies, and structures. The answers can be used to understand the origin of the defects with the goal to improve the coating deposition process. Various analysis methods are available. Defect morphology can be studied optically or with force microscopy depending on defect size. Defect materials can be investigated spectroscopically. The analyses should progress from larger to smaller defects since the larger defects dominate the point-defect scatter even if they are significantly less numerous.

### 5.3.3 Coating Research: Other Coating Properties

**Young’s Modulus and Stress** The Young’s modulus of a coating is required both for the analysis of mechanical loss measurements and for calculations of the level of coating thermal noise. It is therefore important to obtain accurate values of Young’s modulus for every coating including post-deposition annealing. Measuring the temperature dependence of Young’s modulus will be of particular importance for LIGO Voyager, and this capability should be developed.

Residual stress in coatings is likely to be an important property, and there is interesting evidence suggesting that stress can alter mechanical loss of coatings, particularly in silicon nitride. Therefore studies of the effects of residual stress on the mechanical loss and of methods of altering the stress in coatings are important. The use of several measurement techniques can be beneficial in these studies, as each technique has different systematic errors and, for example, different sensitivity to the properties of the coating substrate material.

**Uniformity** As discussed in 5.1.6, coating uniformity is important to avoid the scattering of light out of the cavity mode and into other optical modes, which leads to limits on the optical power and squeezing. With the larger optics proposed for Voyager, maintaining coating uniformity will be even more challenging and thus an important research topic. Moreover, the limitation on obtainable uniformity can come from metrology limitations, so improved metrology is an important research direction. An additional research direction is to explore ion beam mill and corrective coatings, which remove or place additional coating material onto a coated optic after uniformity measurements.

### 5.3.4 Coating Research: Coating Deposition Parameters

Variations in the loss of nominally identical coatings from different vendors have been observed, suggesting that the precise deposition parameters may be important in determining the loss. Thus more detailed measurements of the effects of parameters such as ion energy, sputtering ion, oxygen pressure and thermal treatment may be valuable. While ion beam sputtering produces the lowest optical loss coatings, the mechanical loss of coatings deposited by other techniques has not been extensively studied. Studies of coatings deposited by different techniques (e.g. magnetron sputtering, e-beam evaporation, atomic layer deposition) may enhance understanding of the relationship between loss and structure in these materials.

### 5.3.5 Coating Research: Structural Studies

There is currently little understanding of the mechanical loss mechanisms in amorphous or crystalline coatings, although it seems likely that in the amorphous case the loss mechanism is related to the local atomic structure and may involve transitions of two-level systems as is believed to be the case in fused silica. Studies of the structure of amorphous coatings, aimed at understanding links between structure, composition, loss and optical properties, are a critical part of the research required to develop coatings suitable for use in LIGO Voyager. The suite of structural investigation and atomic modeling techniques discussed in 5.1.8 will also be applied to interesting candidate materials for Voyager, and used to inform

the mechanical loss studies outlined above. Understanding structural defects in crystalline coatings is also likely to be an important area of research.

### 5.3.6 Coating Research: Direct Thermal Noise Measurements

Direct measurements of thermal noise, as described in 5.1.9, are also of interest on coatings suitable for use in LIGO Voyager. In cases (e.g. crystalline coatings) where thermo-optic noise cancellation is required to get the thermo-elastic or thermo-refractive noise below the Brownian noise, Q measurements alone cannot truly probe the thermal noise limits.

### 5.3.7 Coating Research: Coating Design

**nm-Layered Composites** Planar layered composites consisting of nm-scale alternating films of titania and silica can be more and more resistant to crystallization, as the (titania) layers thickness is reduced [151]. These composites behave as homogeneous materials as regards their optical and viscoelastic properties, for which simple and accurate modeling is available [190]. Crystallization inhibition up to very high annealing temperature has also been observed in nm-layered hafnia-alumina composites [137]. Recent work has shown that nano-layer structures of silica and titania can suppress the characteristic loss peaks found in silica films [P1800102].

**Optimized Layer-Thickness Coatings** Since the thermal noise in the coatings is dominated by the total thickness of the more lossy tantala layers, reducing this thickness while maintaining the optical and mechanical properties will reduce thermal noise. Constrained numerical optimization codes have been shown to produce high reflectivity coatings while reducing the volume of high index materials by as much as 20% [191]. Using this optimization method, a coating thermal noise reduction of 9% has been experimentally verified through direct thermal noise measurements [192].

**Multi-material coatings** The use of multi-material coatings to take advantage of the properties of different materials has been proposed [165,166]. In particular, it may be possible to exploit the fact that most of the incident light intensity is reflected by the first few bi-layers of a coating, potentially allowing coating materials with higher optical absorption, but lower mechanical loss, to be used in the lower layers of a coating stack without significantly increasing the total absorption of the coating stack. This type of design may allow the use of aSi in the lower part of a coating stack, taking advantage of both the low mechanical loss and high refractive index of this material.

### 5.3.8 Reduced Coating and Coating Free Optics

Several ideas have been proposed to reduce the mirror thermal noise by reducing the required coating thickness or removing the coating altogether, including:

- Corner cube style retro reflectors

- Brewster angle prism retroreflectors
- Khalili cavities as end mirrors
- Diffraction gratings

While some of these techniques may be more appropriate for consideration for use in LIGO Cosmic Explorer, some of them may provide possible alternative solutions to the coatings for LIGO Voyager. Corner reflectors and Brewster angle mirrors would allow for no coatings to be needed and Khalili cavities would allow for much thinner coatings than conventional mirrors. Experimental work is needed to test some of these concepts for practical limitations. A bench experiment has been done forming a cavity with one Brewster angle mirror and one conventional mirror on fixed suspensions to see if a high finesse cavity can be formed. Follow on work with suspended mirrors will be necessary to evaluate the mechanical stability of such a system. The new prototype interferometer at the AEI in Hannover is slated to test Khalili cavities as a way of reducing coating thermal noise.

**Diffraction gratings** All-reflective interferometers using diffraction gratings as optics avoid problems associated with the transmission of large laser powers through optical substrates. Moderately high finesse optical cavities have been demonstrated using small gratings. The challenge will be to scale up the optical aperture to what is required for a large detector. In addition, absorption by the grating surface can distort its surface profile, possibly resulting in changes in the beam profile as well as power-dependent changes in the diffracted beam shape and efficiency. Modeling has been done along with sample produced, but these effects need to be investigated more in depth. Investigations of mechanical loss in gratings are needed to verify thermal noise levels as are direct thermal noise measurements. Demonstration of high reflectance values is also important.

### 5.3.9 Mirror substrate research

**Silicon** The OWG is investigating alternative materials to fused silica for use as test mass substrates and which can be used in low temperature detectors. Both silicon and sapphire potentially offer superior performance at cryogenic temperatures and/or specific frequency ranges. Different substrate materials, operating temperatures, and laser wavelengths may also require and/or allow for different coatings and suspension connection techniques.

Previous research efforts on silicon have largely focused on acquiring and fabricating cylindrical test specimens and investigating their mechanical properties as a function of doping. Studies of silicon properties, including mechanical loss for predicting thermal noise, of different crystal orientations are valuable. In addition, silicon cantilever micro-resonators with resonant frequencies in the sub-kHz range have been fabricated to explore dissipation mechanisms in a regime where thermoelastic effects are significant. Surface loss effects are also emphasized by the large surface-area to volume ratio of the micro-resonators. Preliminary experiments measuring the dissipation have been carried out and reveal disagreement with theoretically predicted loss.

Understanding the optical loss of silicon at 1550 - 2000 nm is also an active area of research. The high thermal conductivity of silicon could significantly reduce the effects of thermal loading of transmissive components if the optical loss is low enough. Understanding the temperature dependence of the optical absorption along with all other thermo-optic and thermo-physical properties is important. Research will be required to develop suitable components if a change in wavelength is considered. Silicon mirrors and suspension elements have an advantage of being conductive thus control of charging effects may be easier to implement. Nonetheless, charging will need to be investigated since doping and especially coatings can influence the charging properties.

Investigations of the wavelength dependence of the absorption of silicon is of interest, as it may be beneficial to consider operating at a longer wavelength to reduce or eliminate two-photon absorption effects.

**Sapphire** Recent efforts have yielded information about the mechanical and optical properties of sapphire, methods for growing and processing large sapphire blanks, and ways to achieve high homogeneity, low absorption sapphire. Studies on annealing for improved optical absorption have been extended to elucidate further details of the kinetics of the out-diffusion process. Gathering experimental data at low temperature is important to predict the performance of cryogenic sapphire test masses. The KAGRA project is pioneering this effort. Room temperature sapphire is also a potential mirror substrate for detectors optimized at higher frequencies. Measurements of mechanical properties including mechanical loss as a function of crystal orientation are also important for predicting substrate and coating thermal noise.

### 5.3.10 Mirror Substrate Research: Parametric Instabilities

As discussed in 5.1.12, parametric instabilities are a potential problem arising from high laser power. These undesirable effects result from exchange of energy between light stored in cavities and acoustic modes of the cavity mirrors. Experimental research and modeling of parametric instabilities, and possible suppression techniques, in silicon mirrors at high power levels will be required.

### 5.3.11 Mirror Substrate Research: Composite Masses

Increasing the mass of the test masses reduces the influence of both classical and quantum radiation pressure noise. Beyond a certain size, however, it is impractical to fabricate monolithic masses. Using large masses made as a composite of multiple, smaller pieces could circumvent this problem. Non-cylindrical mass distributions could also be used to increase the total mass and total angular moment of inertia without increasing the optical path length within the substrate. The larger translational and angular moments of inertia would reduce the radiation pressure noise and the influence of the Sidles-Sigg instability. Thermal noise issues related to mechanical loss from the interfaces will have to be understood.

### 5.3.12 Mirror Substrate Research: Charging

The use of semi-conductor silicon optics will require several new areas of research into charging.

- Silicon resistivity at low temperature. Determination of the resistivity and time constants of charge motion on silicon at cryogenic temperatures. At these low temperatures the free carriers begin to freeze out. Research should assess the charge mobility on high purity silicon substrates with resistivity at the level few thousand ohm-m using Kelvin probe, 4-terminal measurements and electromechanical oscillators.
- Coating charging. R&D is necessary to assess the time constant of charge motion on the surface of coatings necessary for cryogenic detectors. Materials include amorphous silicon and crystalline coatings based on AlGaAs and AlGaP. Pushing the performance of future detectors to lower frequencies (1-10Hz) will require the assessment of charge noise mechanisms. These include surface charge with low mobility on the surface of coatings, charge motion in silicon surface/substrates and the interaction of these charges with the electric field of the ESD and nearby grounded hardware.
- Study of discharge mechanisms including (i) utilising He gas rather than N<sub>2</sub>, (ii) mono/multi layers of gas adsorbed on surfaces (iii) pumping time at low temperature (iv) charging mechanisms.

While optics research targeted at LIGO Cosmic Explorer is rather more speculative and long term, it is clear that research on optical components for future ground-based interferometers must begin well in advance of any complete conceptual design. One important consideration will be the operating temperature of LIGO Cosmic Explorer, and research into the temperature dependence of both the thermal noise and the optical properties of coatings and mirror substrates will be critical in informing this choice. Thus all of the research avenues detailed in the previous sections for substrates and coatings for use at room temperature and at cryogenic temperature are potentially of relevance for Cosmic Explorer.

It should be noted that some of the promising technologies discussed in the LIGO Voyager section, in particular AlGaP and AlGaAs crystalline coatings, have not yet been demonstrated at the size required for future gravitational wave detector mirrors. Technical restrictions, such as the maximum available diameter of GaAs substrates on which AlGaAs can be grown, may potentially restrict or prevent the use of these coatings in Voyager. However, on the time-scale of Cosmic Explorer, it seems more likely that some of these technical issues may be overcome. In addition to this, research into improved amorphous coatings should continue to be pursued. As discussed above, some amorphous coatings (e.g. aSi) have already been demonstrated to have mechanical loss factors that are equivalent to or better than those measured for crystalline coatings, and further progress with modeling and theoretical understanding of the properties of these materials may allow both mechanical loss and optical absorption to be optimized. Finally, other concepts for making mirrors without coatings (e.g. waveguide mirrors, consisting of coating-free structured surfaces), or with a much reduced coating thickness, are also of interest and further research into the application of these methods to large-scale interferometer mirrors is of interest.

## 6 Suspensions and Seismic Isolation Systems

The research of the Suspension and Seismic Isolation Working Group (SWG) is aimed at providing the necessary isolation, alignment, and control of the interferometer optics from seismic and mechanical disturbances while simultaneously ensuring that the displacement due to thermal noise of the suspended systems is at a suitably low level. To first order we can divide the research into two broad subdivisions, suspensions and isolation, both of which involve mechanical and control aspects. Suspension research involves study of the mechanical design of the suspensions, the thermo-mechanical properties of the suspension materials and suitable techniques for damping suspension resonances and applying signals for interferometer control. Isolation system research involves mechanical design and active control for isolation and alignment. The overall isolation of the optics comes from the product of the two systems.

The isolation and suspension system for the most sensitive optics in Advanced LIGO is comprised of three sub-systems: the hydraulic external pre-isolator (HEPI) for low frequency alignment and control, a two-stage hybrid active & passive isolation platform designed to give a factor of  $\sim 1000$  attenuation at 10 Hz, and a quadruple pendulum suspension system that provides passive isolation above a few Hz. The final stage of the suspension consists of a 40 kg silica mirror suspended on fused silica fibers to reduce suspension thermal noise.

The R&D for baseline Advanced LIGO isolation and suspension sub-systems is complete. These systems are in operation at the LIGO facilities, and were successful during Advanced LIGO's first two observing runs. Upgrades to the Suspension and Seismic Isolation and Alignment System are grouped into four broad categories: In section 6.1 we describe ongoing work on immediate improvements and risk reduction for the baseline Advanced LIGO detector. In section 6.2, we describe the A+ (and beyond) improvements which can improve the baseline sensitivity without substantially altering the existing equipment. In section 6.3 we describe the Voyager R&D which will make more substantial changes and allow us to operate with cooled optics, and finally in section 6.4 we describe the long term work on suspension and isolation systems which build the foundation for the ultimate ground-based gravitational wave detectors.

### 6.1 Vibration Isolation and Control R&D for Incremental Upgrades to Advanced LIGO

There is ongoing R&D work to provide incremental improvements to Advanced LIGO to improve performance of the aLIGO vibration isolation and suspensions, and improving the stable operation of the observatories by making the instruments more capable of withstanding high winds and teleseismic earthquakes. This type of work is designed to be easily incorporated with the existing detector systems with minimal disruption of the Observatories. These upgrades include work to add additional environmental sensors and incorporate them into the controls, adding more sophisticated control algorithms to improve performance during unusual environmental conditions, small mechanical changes to damp vibration modes, and relatively modest changes to the facilities.



### 6.1.1 Tilt/horizontal coupling and advanced seismometers

One of the limits to the performance of seismic isolation systems is the coupling between ground tilt and horizontal motion of the isolation platforms. This is fundamentally caused by the inability of a horizontal sensor (or a passive horizontal isolation stage) to distinguish between horizontal accelerations and tilts in a gravitational field.

This tilt-horizontal coupling causes a variety of problems and is a basic limit to the performance of the isolation systems at low frequencies (below  $\sim 0.3$  Hz) [193]. Seismic motion in the 30-300 mHz region is governed by the microseism, wind and large earthquakes. Excess motion in this band affects the locking of the interferometer and consequently the duty cycle of the observatories. It may also affect noise in the GW band by frequency up-conversion through non-linearities in the interferometer control. The aLIGO HEPI and ISI systems are now limited by the tilt-horizontal coupling at low frequencies, and an external sensor could be easily integrated into the system to reduce amplification of low frequency ground motion.

We have developed sensors to measure the rotational acceleration of the ground or of stages of the seismic isolation system in vacuum, which could be used to remove the rotational component, creating a purely translational horizontal sensor and/ or reduce the rotation of the stages. Several rotational sensors have been investigated in the past [194–197] and two approaches are currently being pursued. These include:

1. A tiltmeter based on a low-frequency flexure-beam-balance with an optical readout has been developed which looks very promising [198]. Two of these instruments, known as beam rotation sensors, (BRS), were installed at the LIGO Hanford Observatory and have been shown to reduce tilt-noise in a co-located seismometer by factors of  $\sim 10$  below 0.1 Hz under windy conditions. Four additional BRSs are currently being installed at LLO. Recent studies suggest that incorporating these sensors into the isolation control system is beneficial to the interferometer, enabling locking under higher wind speeds than before. A compact and UHV-compatible version of this instrument with an interferometric readout is also being developed which is suitable to be mounted directly on the ISI. Such an instrument has the potential to significantly improve the angular control of the ISI and further reduce the low-frequency differential motion of the platforms.
2. Suspending a horizontal seismometer. This approach is distinct from the others in that the seismometer is made to passively reject tilt noise, thus producing a tilt-free horizontal sensor [197].
3. There has also been work on a variety of other technologies for making gyroscopes. Work on passive ring gyros [196] has concluded, and some investigation into Atomic Spin Gyroscopes has been undertaken [199, 200].

### 6.1.2 Wind Mitigation

The largest driver of tilt on the technical slabs at the LIGO Observatories is the wind. In addition to measuring the tilt as described in section 6.1.1, work has recently begun to reduce the wind driven tilt of the buildings. The tilt effects are seen most clearly at the

End Stations where the End Test Masses are located. The Hanford detector currently has difficulties when the wind speed exceeds about 10 meters per second. The wind causes tilting of the slab between at least 10 mHz to 1 Hz. It has been suggested that wind breaks could reduce the force loading on the buildings and thereby reduce the typical tilt of the slabs during windy times [201]. Preliminary work has begun to design and test ways to install wind breaks at the Hanford End Stations. Since the amplitude of the spectrum of ground tilt seems to scale with the square of the wind speed, even modest reductions of the wind speed at the building could potentially have large pay-offs in improved duty cycles of the instrument. This work would be completely complementary to the work on better measurement of the ground tilts. Transforming our large, unmeasured tilts into medium sized and well measured tilts would allow us to reduce the excess motion of the isolation platforms below 100 mHz.

### 6.1.3 Pier Motion Control

The motion at the top of the HEPI piers is significantly higher than the ground motion, especially in the horizontal direction. A set of solutions have been suggested to address the problem [202]. We have recently shown that the amplification is due to motion of the bending piers and the attachment of the payload to the pier top [203]. The FE model is being updated with new experimental results. These results shed considerable light on an old problem, and enable research on how to address the problem. Solutions to mitigate this motion amplification will be proposed and studied. Both passive (reinforcement, truss, cross beams, cables) and active solutions (feedback shaker on the chamber, active tendons) should be considered.

### 6.1.4 Control System Enhancements for the Existing System

Advanced LIGO has an impressive array of sensors and a flexible control system. Most of the baseline Advanced LIGO control schemes use local information to control the seismic platforms and Interferometer readouts to control the interferometer lengths and angles. As system integration proceeds, studies need to be conducted to investigate optimal ways to combine all the sensor information to achieve the best interferometer performance. For example, using feedforward to directly cancel the contribution of ground motion signals which appear in the interferometer signal was demonstrated in Enhanced LIGO [204]. Studies on the coherence between various channels from the ground, up through the seismic platforms and suspension systems to the various interferometer readouts should be conducted. When coherent contributions are found, they could be used to inform the development of advanced control techniques to be layered onto the existing controls. Other control improvements have also been suggested such as optimally distributing the control authority between the isolation stages [205] to improve system locking and robustness, and changing the control laws in real time in response to changing environmental conditions [205, 206]. We often experience variations in the local environment arising from many different sources such as wind at the Hanford Observatory, storms in the gulf of Mexico which cause large microseismic motions, logging near the Livingston Observatory, and teleseismic earthquakes (those far from the Observatories). We can now accurately predict the arrival time of teleseismic waves, and we can predict that amplitude within a factor of about 2.5 [207, 208]. We are now implementing an updated control scheme to better compensate for the motions [209].

Another question which needs to be answered is “what in particular limits the upper unity gain frequencies of the isolation control loops for the Advanced LIGO seismic isolation platforms?” A clear understanding of the practical limits could be used to inform a campaign to modify the platforms to achieve better isolation performance, particularly in the 5-40 Hz band. For example, a numerical study has shown that a high frequency blend with a force sensor (even a virtual force sensor) may help to increase significantly the controller bandwidth of the active isolation stages, without compromising the isolation [210, 211]. This technique will be further studied on more elaborate models.

An interesting risk reduction idea would be to develop a fail-operate control system for the seismic isolation platforms. These systems use position sensors, geophones and seismometers in their control loops. While these instruments have very low failure rates, we cannot exclude the possibility that one of them could stop working during operations, thus compromising the functioning of the platform and the interferometer. It is therefore important to study methods to monitor performance of the sensors, actuators, and the mechanical plant to help mark when the behavior of a component begins to degrade, and to identify which component is malfunctioning. A set of basic Matlab scripts have been developed to perform this task [212] but those functions are difficult to use in the observatory environment and can only be used while the interferometer is offline. A realtime estimator system could be developed to continuously compare the expected and actual performance of the system so that diagnostics could be run continuously in the background. Development of realtime state-space tools to allow this to be implemented in the LIGO realtime computers is now being pursued jointly by the Controls working and is described in section 8.2.6. In the event of a problem, it is also important to study “emergency control schemes” that would allow the platform to keep operating with a malfunctioning sensor. Feedforward techniques to either compensate for the loss of performance, or possibly reconstruct the lost signal are good candidates. Control schemes accounting for the failure of each type of sensors should be proposed and studied.

There is ongoing research on the controls for the Suspension system. We have recently demonstrated a technique called ‘Global Damping’ to reduce the sensor noise from the Suspension which couples into the Interferometer [213]. Global damping also reduces the interdependence of damping and IFO length control. This technique could be incorporated into the control system at the Observatory and studies should be done to investigate its utility for the full-scale interferometer.

We have also modeled the benefit of implementing feedforward control from the ISI table to the top stage of the suspension at frequencies near the microseism. This control will need to apply both force in the longitudinal direction and torque in the pitch DOF to generate the necessary control force for the pendulum, so the controllers must be individually tuned [214]. Implementation at the sites will allow us to ascertain the long term benefit and performance of this system.

Testing of adaptive control schemes have been done [205] to automatically adjust the trade-off between damping strength and feedthrough of sensor noise in the GW band. Tests are also planned to study damping methods of the bounce and roll modes within the monolithic section. As described above, studies of how to distribute the control authority (hierarchical control) are underway, both for control of a pendulum chain, and also for offloading pendulum

control to the seismic isolation system.

### 6.1.5 Suspension Design Updates

Upgrades to specific suspension designs are being considered as potential near-term improvements for Advanced LIGO [215]. First, the current size of the beamsplitter optic (37 cm diameter) was established early in the Advanced LIGO design, and is the limiting optical aperture. Retrofitting the detector with a larger beamsplitter optic would enhance optical performance [216]. This will require design changes to the beamsplitter triple suspension itself and its supporting structure [217, 218]. A larger beamsplitter is now included in the currently proposed design for A+ as discussed in section 6.2.

Also under consideration is the addition of actuation on the beamsplitter optic. Currently global control signals are applied to actuators at the middle mass of the beamsplitter triple suspension. Actuation directly on the optic would allow wider bandwidth operation to improve lock acquisition.

Secondly, noise in the signal recycling cavity (SRC) length coupling into the GW readout may be a greater noise source than we would like. In particular the highest vertical and roll modes (around 28 and 41 Hz) in the small and large HAM triple suspensions (HSTS and HLTS) potentially add noise in the GW band. This noise could be addressed by adding a third stage of cantilever springs at the middle masses of the HSTS and HLTSs, taking the vertical mode below 10 Hz and increasing the overall vertical isolation. Preliminary design work has begun on this possible modification [219]. However given that the HAM-ISI isolation performance is better than the design requirements, this proposed modification is not being taken further at present [220]. An alternative is the addition of passive dampers which would reduce the peaks at the highest vertical and roll modes at the expense of modest increase in thermal noise performance on the wings of the peaks. A design based on the bounce and roll dampers (BRDs) developed for the quad suspensions, discussed below, has been developed and tested on both an HSTS and an HLTS [221, 222]. They could provide a reduction in Q of a factor of about 10. At present further work on this is on hold.

Thirdly some improvements to the test mass quadruple suspensions are also underway. Work by colleagues in the LISA area and subsequent follow-up by LSC groups has shown that enhanced gas damping in small gaps could lead to excess noise in aLIGO suspensions [223–225]. In the current suspensions there is a small 5 mm gap between the end test mass and end reaction mass to allow sufficient actuation force for cavity locking using electrostatic drive (ESD). A new design of reaction mass, called the annular end reaction mass (AERM), has been developed which is donut shaped to reduce the squeezed film damping associated with this small gap. [226]. The outer dimensions and material are the same as the present mass, with a central hole  $\sim 220$  mm in diameter. Since this mass will be lighter than the currently suspended reaction masses, the penultimate reaction mass design requires modification to increase its mass to compensate, so that the existing blades are still correctly loaded. AERMs have now been procured and it is planned to install them on both detectors after the end of O2. This activity is currently underway.

Two other developments for the quad suspensions are also underway associated with damping high Q modes. We have developed a non-magnetic damper for the upper intermediate mass

blades to suppress the thermal excited motion at the internal modes of the blade [227]. We have also developed passive dampers, so-called BRDs (bounce and roll dampers), to reduce the size of the bounce and roll mode peaks around 10 and 13 Hz associated with compression and extension of the silica fibers between the test mass and penultimate mass [228]. Both types of dampers have been installed at LLO and are due for installation at LHO after O2. The benefit derived from the use of BRDs at LLO has already been observed. They are able to reduce the quality factor of the bounce and roll modes from 500,000 to less than 10,000, reducing the time taken for the modes to damp down from several hours to a few minutes or less [229].

BRDs are also under development for possible application to the beamsplitter suspension [230], where it has been suggested that damping their high frequency bounce and roll modes could lead to better plant inversion for better MICH feedback.

A further addition to the quad suspensions is under consideration: the use of passive violin mode dampers (VMDs) ([231,232]). The violin modes of the quad suspensions (fundamental plus higher harmonics) have very high Qs (of order  $10^9$ ) which can become significantly excited by earthquakes or other disturbances. Methods to actively damp them have been developed and are in use. However the feedback loops require manual tuning and this can take time and can be challenging when the modes are close in frequency. The use of passive tuned mass dampers attached to the PUM could alleviate the situation. For example reducing the Qs down to  $\sim 10^7$  would reduce decay time of excited modes to  $\sim$ hours.

## 6.2 Research and Development for LIGO A+ and beyond

A joint preliminary construction proposal to the U.S. National Science Foundation and the U.K. Science and Technology Facilities Council, with participation of the Australian Research Council has been submitted in 2018, to support the LIGO A+ upgrade. A+ seeks to expand the volume of the universe observable by Advanced LIGO by factor of approximately 6.6 and 4.1, for  $1.4 M_{\odot}$  neutron star binaries and  $30 M_{\odot}$  binary black holes, respectively. With funding in early 2020, the upgrade could be complete by early 2023 and enter observations in 2024. The details of the A+ submitted design can be found in [233,234]. The A+ baseline for suspensions includes a larger beamsplitter suspension, new compact triple suspensions for relaying the balanced homodyne signals, and the possibility of actuation changes for the beamsplitter and thinner fused silica fibres.

More generally, there are also activities to develop technologies that go beyond the A+ baseline, using the same LIGO infrastructure, but which must be under development to ensure sufficient time for developing engineered solutions. These will require more substantial changes to the Suspension, Isolation, and Alignment subsystems. For example, suspensions may be replaced to take heavier test masses. In addition to the improved isolation from ground motion, improved rejection of thermal noise is also being studied, although the system will remain at room temperature. The seismic isolation systems will remain largely unchanged, but could be improved with higher performance sensors. The only adjustment would be to the pier mountings, to allow more space for proposed longer suspensions. The total mass of the QUAD suspension would remain unchanged, but the relative masses of different components would change to cater for a larger test mass (e.g. 80 kg).

### 6.2.1 New design of compact triple suspensions: required for A+

The implementation of balanced homodyne readout in A+ will require a new design of triple suspension for the relay mirrors in the revised optical layout. The design can be based on existing triple suspensions such as the HAM Small Triple Suspensions (HSTS), suitably reduced in size to accommodate smaller optics and to meet tight space and height constraints, while still meeting the required low noise performance.

### 6.2.2 Larger Main Optics: beyond A+

Studies are underway to explore the design of improved suspensions for the 4 main optics (the ‘Test Masses’) of Advanced LIGO. It has been shown that by increasing the mass from 40 kg to 80 kg, and increasing length of the final suspension fiber to 120 cm, the thermal motion of the optic can be lowered by a factor of 2 from the current design [235]. By using 160 kg test masses and 120 cm long final suspensions, the thermal noise of the suspension can be improved so that it is 3 times better than the current design. However, work needs to be done to optimize the entire suspension design [236] so that it has good overall seismic isolation. Work needs to be done to show how to fit this larger, heavier suspension onto the existing Seismic Isolation Platforms. One possibility is mount the longer suspension to the top of the final stage of the seismic isolation system, rather than the bottom, and interleave the suspension through the isolation system. A second option would be to raise the entire seismic isolation system up on its existing support structure. Since the optical table would need to move up substantially, it is not clear what the performance impact would be and it requires investigation.

### 6.2.3 Alternative Control Approaches for Larger Suspensions: A+ and beyond

The design and installation of larger optic suspensions gives an opportunity to revisit the design of the caging, sensing, and control of the pendulum system. One alternative approach is to merge a portion of the suspension cage and the reaction chain [237]. This is a major design change, but could result in improved access and alignment and it would allow all the pendulum DOFs to be sensed and controlled. The reduced mass of the combined cage and reaction structure would allow a more massive main optic.

### 6.2.4 Improved OSEMS: A+ and beyond

The Suspension sensors should be improved. A set of Optical Sensor/ Electromagnetic Motors (OSEMS) are used to provide actuation and local sensing for the Advanced LIGO suspensions. Improved sensors could be used to increase the damping of pendulum modes without compromising the performance, but making practical, low-cost units which have better performance than the existing OSEMS is a challenge. Studies are underway to provide better thermal stability and optical modifications to improve the signal to noise by a factor of two by adding a displacement-doubling prism to the optical system [238].

A promising alternative is to replace, or add in parallel, compact interferometers such as EUCLID [239] or the Homodyne Quadrature Interferometer demonstrated in [240]. These

devices have approximately 4 orders lower noise than BOSEMS, but further development is required to integrate them into suspensions and ensure they meet LIGO vacuum requirements. Improved sensors will be necessary for achieving the performance goals of future upgrades at low frequency.

### 6.2.5 Seismic Platform Interferometer: A+ and beyond

It is also possible to improve the performance below 1 Hz with an auxiliary system which reduces the differential motion and tilt of the various optical tables in the detector. This type of approach has been discussed for many years, and is traditionally called a ‘Suspension Point Interferometer’ (SPI), i.e., an interferometric sensor which measures between the points which suspend the arm mirrors [241].

The systems under investigation are slightly different; the method involves controlling the relative motion of the optical tables, and hence an alternative name is Seismic Platform Interferometer. The relative motion of the tables for this system will need to be measured in at least 3 degrees of freedom, namely length, pitch, and yaw. This will allow the detectors to be mounted securely to the table, and will also allow the benefits to be shared by multiple suspensions on the same table, a common situation on the HAM optical tables.

A prototype system has been demonstrated at Stanford [242, 243] using a fiber coupled 1.5 micron laser, which used a Mach Zender interferometer to measure the inter-platform length motion and an optical lever to measure differential angles. At the AEI 10m prototype, another SPI prototype is being used successfully. It is based on a set of Mach-Zender interferometers, and uses a LISA-pathfinder style phasemeter for readout [244]. The rotational degrees of freedom (DOFs) are sensed via differential wave-front sensing. The target sensitivity is 100 pm/ $\sqrt{\text{Hz}}$  and 10 nrad/ $\sqrt{\text{Hz}}$  at 10 mHz for displacement and rotational DOFs respectively. In addition to the SPI, optical lever to improve the resolution of the angle DOFs are investigated.

Considerable work remains to adapt either of these systems for use with Advanced LIGO (e.g. stable mechanical coupling to stage 1 of the HAM-ISI or stage 0 of the BSC-ISI, reliable UHV compatible fiber coupling, control integration). In addition, were this system to be used for the 4 km arms, then considerable work would be required to achieve the necessary laser frequency stability required to realize a beneficial system.

It should be noted that improved rotational sensing described in Section 6.1.1 and the SPI are complementary approaches to the low-frequency noise issue. It is also important to realize that since the optical tables for Advanced LIGO are controlled in all 6 degrees of freedom, once new SPI or tilt sensors become available, they can be incorporated into the existing control system easily, because the seismic tables will not require modification.

### 6.2.6 Improved Seismic Sensors: A+ and beyond

Future improvement of the seismic isolation performance will require the use of a very-low-noise inertial instrument, with noise performance roughly a factor of 100 better than today’s GS-13. Such an instrument remains to be designed, built and tested. The requirements for the design of this instrument will include: very low sensitivity to temperature gradients,

very low thermal noise, very low sensitivity to environmental and actuator electromagnetic fields, and very low readout sensor noise. Both horizontal and vertical instruments will be necessary for future seismic isolation upgrades. Several improved inertial sensors are now being studied [245, 246].

### 6.2.7 Sensor Development at Low Frequency: A+ and beyond

Some of the technical noises in Advanced LIGO are currently more than an order of magnitude above the instrument fundamental limits at 10 Hz [247]. Angular controls noise and scattered light noise are examples of such noises. As a result, intermediate-mass black holes with masses  $\simeq 10^2 - 10^4$  solar masses lie outside of the most sensitive band of the instrument. Therefore, improving the low-frequency sensitivity of the GW detectors plays a crucial role in extending both the mass and spatial range of detectability. A recent proposal [248] shows a path towards detecting gravitational waves at 10 Hz and at lower frequencies.

One of the most important elements required for such an upgrade is the 6D isolation system [249]. In the core of this instrument is an inertial mass whose position is monitored relative to the optical bench using interferometric readout in all six degrees of freedom. The mass is suspended using fused silica fibers and the whole assembly is quasi-monolithic similar to LIGO test mass suspensions. This design is motivated by low thermal suspension noise and readout noise and allows for a significant improvement of the ISI motion below 10 Hz.

### 6.2.8 Studies of the monolithic final stage: A+ and beyond

Extensive characterization (strength, dimensions, mechanical loss) of fused silica fibers as suspension elements [250, 251], produced using both oxy-hydrogen and laser-based pulling techniques [252], has been done. Welding techniques and silicate bonding techniques including characterization of associated losses [253] has been done, along with extensive exploration on the ear shape and fiber shape.

Several monolithic suspensions (4 per detector since O1) have been installed in Advanced LIGO, with the end test masses currently undergoing re-installation in the downtime between O2 and O3, to enable test masses with improved coatings and annular reaction masses. Integration of the suspensions into the Advanced LIGO interferometers has significantly reduced the low frequency thermal noise, to allow access to the 10 Hz to 40 Hz range of the interferometer. As the interferometer noise is improved, participation of experts will be critical to understanding the interferometer performance.

Several areas of research could yield enhancements to Advanced LIGO suspensions. Further understanding and characterizing of losses in silica fibers including investigations of non-linear thermoelastic noise and of surface losses could lead to improvements. Changes in fiber neck shape including shorter neck and thicker stock could lead to enhanced thermal noise performance. Research is also underway to further understand the role of weld loss in addition to techniques to observe and ameliorate stress in the weld regions. Furthermore, an increase in strength of the fibers could allow reduction in cross-section and in vertical bounce frequency, enhancing isolation. Investigations of the silicate bond mechanical loss and strength as a function of time and following temperature treatments are underway to reduce



further the loss contribution and optimize ear design. As noted in 3.1.4, increased fiber stress can move the bounce/roll modes below 10 Hz which reduces the requirements on passive or active damping of these modes. A test suspension has been hanging at 1 GPa for over 3 years at LIGO Hanford, and a further series of tests at stresses up to 1.5 GPa will verify the robustness of an engineered suspension. Such an upgrade can easily be incorporated into a 60 cm aLIGO QUAD system operating at higher fibre stress (to lower the vertical bounce mode out of band and increase the violin modes) such as in A+, or an upgraded longer/heavier system for beyond A+.

Suitable prototype long fibres manufactured from thicker diameter stock material have been demonstrated, and the necessary prototype upgrades to fibre pulling and characterisation equipment have been undertaken. Laser welding, which has been proven on larger test ear horns, has been tested and provisionally shown to work. The research has led on to preparation for demonstration of the first four fibre 1.2m long heavy mass (120-200kg) suspension. Long term fibre stress corrosion experiments are ongoing to demonstrate and characterise the longevity of fibres supporting a higher working stress, in the stress regime of 1-5 GPa [254].

### 6.2.9 Violin Mode Analysis

Extensive analysis of the violin modes of the monolithic suspensions, with Q's in the order of a billion, is underway. Long term tracking of their amplitude, phase and frequencies for the fundamental and higher order harmonics, together with more accurate FEA models of the actual installed suspensions, will contribute to the characterisation of the monolithic suspensions and long term monitoring of their state [254,255]. This work has already proven to provide temperature sensitivity of the suspensions which has been used to associate violin modes harmonics with individual suspensions and fibres, in a non-invasive way. We are enhancing our understanding of the effects of the complex geometries of the ends of the suspensions fibres, horns and ears with violin mode phenomena such as inharmonicity, frequency spread, which are well understood, together with frequency splitting which is an area of ongoing research. Future research will include monitoring violin mode Q values over time (believed to increase) to refine our understanding of the mechanical loss mechanisms of the suspension fibres fused silica material. Also searches for non-Gaussian events on the monolithic suspension are to be carried out. This regime of research will ultimately provide tighter bounds on the contributing mechanical loss terms for characterising thermal noise of each individual installed suspension.

### 6.2.10 Mechanical Upconversion: Crackling Noise

Some sources of gravitational waves produce short, impulsive events in an extremely large body of data, and so characterization and reducing “background” transients of technical origin is important. Investigations of non-thermal noise originating in the fused silica fibers has been carried out with no non-thermal noise being seen at modest sensitivity (insufficient to exclude it as a significant noise source for aLIGO). Work has been done to study the noise associated with the violin modes of the silica suspensions in GEO 600 [256]. Further work has started to extend these studies by modeling and analyzing data from O1 and future engineering or science runs to put upper limits on this noise component in Advanced LIGO.

Direct experiments to characterize the level of and/or put upper limits at a meaningful sensitivity level to potential non-Gaussian transient events associated with the Advanced LIGO suspension system are challenging. However new ideas for carrying out such experiments are encouraged.

One approach which is being pursued to observe impulsive releases of energy or acoustic emissions (“creak effect”) is to strain the element statically while also driving the element through a large amplitude motion at low frequency below the measurement band, while interferometrically measuring the element at high sensitivity in band (above 10 Hz). By large amplitude motion we mean much larger (100~1000 times) than the out of band motions estimated through modeling. We will drive the large amplitude low frequency motions in a common mode fashion between two identical devices under test while measuring the noise which will be uncorrelated between the two elements. Experiments of this type are underway to measure or put upper limits on noise from maraging steel cantilever blades [257–260] and separately for silicate bonds.

Furthermore, these techniques can be applied to study low frequency mechanical noise in glassy metals, potential materials for use in the suspensions of future low-frequency interferometers.

### 6.2.11 Maraging flexure robustness studies

Maraging steel is the material used for the construction of crucial passive isolation components in the ground-based detectors. In VIRGO and KAGRA the operating stresses are high and in the region of 1-1.5 GPa. Most of the aLIGO suspension blades are stressed to around 1 GPa [261], and the aLIGO seismic system has stress levels which are roughly 50% lower. There have been occurrences of cantilever blades failing in Virgo and KAGRA. Initial studies point to hydrogen embrittlement [262] in the maraging steel which can lower the ultimate tensile stress. Furthermore, it appears that the hydrogen can migrate towards regions of highest stress. It is of interest to study the reasons of these failures both to avoid further damage to present and future suspensions.

## 6.3 LIGO Voyager

The LIGO Voyager design represents a major shift in the LIGO design with adoption of cryogenic test masses. Investigations of moving to cryogenic temperatures have shown that they can provide significant improvements in the thermal noise, even with test mass temperatures as high as 120 K to 130 K. The LIGO Voyager design is based on the ideas presented in the ‘Blue Team’ design and incorporates a 150 to 200 kg silicon optic which is radiatively cooled to about 123 K with a cold-shield held at around 77 K. Increasing the mass of the interferometer mirrors will linearly reduce the displacement noise due to radiation pressure noise. Changing the size, mass, temperature, and material for the optic requires many changes across the detector, and especially close collaboration between the Optics Working group (see Section 5) and the Suspension and Isolation Working Group.

This design requires test masses and heat links (either to the mass or to the cooling shield) with excellent thermal and vibration properties. These are likely copper, but might also be

silicon or sapphire. Any cold system will require an ultimate heat sink. Current cryocoolers, even those designed for gravitational wave detectors [263] are not free from vibration, and the heat transport properties of heat links are limited. Thus, it is essential that studies of systems with suspension elements of suitable design and dimensions to provide an efficient path for required heat conduction while still maintaining good thermal noise and mechanical isolation performance be carried out, and followed with experimental demonstrations. A possible way to reduce the requirements of the heat links is to shorten them by having flowing liquid nitrogen in pipes within the main vacuum system. Chilling the nitrogen could further reduce the requirements on these heat links.

### 6.3.1 Experimental Demonstrations of Cooled Optics

Cooling the optics without compromising their vibration performance is a significant challenge, and experimental demonstrations are critical to the long-term engineering success of the program. These experiments are now underway [264–268]. It has recently been realized that the scattered light from the optic will be a significant heat load for the cryogenic shield. It is critical that we demonstrate technology capable of removing the heat from both the absorbed laser power in the optic and the scattered power from the optic which is incident on the cryoshield – while maintaining vibration levels small enough to prevent upconversion of the scattered light returning from the shield back into the interferometer beam.

Another of the challenges to implementing a cryogenic silicon suspension is extracting the required heat in a reasonable amount of time. A purely radiative cooled suspension with a 150 to 200 kg silicon test mass will require cool down periods to 120 K on the order of weeks at best to months at worst. Research is currently underway to investigate the benefits and feasibility of introducing a cool down period which incorporates aggressive cooling technology that would not be permitted during Observing Runs. Current research is focussing on the methods of convectively cooling the suspension in a dry atmosphere and contacting the suspension stages with a movable cold link [264].

Methods of maintaining room temperature detector infrastructure near the cold suspension stages should be pursued. This issue is particularly relevant for Advanced LIGO hardware that will be reused in a future cryogenic system. Notable components that might be influenced are suspension springs, electronics, and sensors and actuators.

**Materials** Investigations of materials suitable for construction of elements of the isolation and suspension systems with good properties for use at cryogenic temperatures should be studied ,e.g. silicon carbide which has excellent stiffness to weight ratio (specific stiffness) and low thermal expansion constant and silicon, which has excellent thermal conduction properties and high specific stiffness.

Passive damping of structures at low temperatures is also an area of concern. Viton is quite lossy at room temperature, and is vacuum compatible, and so is used extensively in Advanced LIGO to help control structural vibrations above 80 Hz. Unfortunately, handbook values [269] indicate that it is not effective at cryogenic temperatures, so investigations into replacements should be undertaken.

**Integrated Control of Cryogenic Suspensions** The design and installation of new optic suspensions should be accompanied by a careful design of the caging, sensing, and control of the suspension and final optic. One alternative approach to the Advanced LIGO design is to split the suspension cage into a warm and cold section and to combine the structures of the cold portion of the cage with the cryogenic shield [267].

**Cooling Options** During the detector operation the test mass needs to be kept at a temperature of about 123 K. Therefore, it will be necessary to extract up to 10W of heat absorbed by the test mass from the laser beam under high vacuum conditions. This can be done by means of radiative cooling. However, the thermal emissivity of the silicon test mass and the optical coatings deposited upon its face is too small to provide the necessary rate of radiative cooling. The emissivity of a barrel of a cylindrical test mass can be increased by using a high-emissivity coating such as DLC (Diamond-Like Carbon) coating [270], Acktar Black coating [271] or the carbon nanotube NASA Goddard black coating [272]. However, any coating of the test mass introduces additional mechanical loss and additional thermal noise associated with this loss. It was found that Acktar Black coating results in up to 10% increase of the total strain noise of LIGO Voyager design [273]. Carbon nanotube black coating gives significantly lower noise level but it can be damaged under mechanical influence [274]. Choosing a promising coating with low mechanical loss and high resistance to external influences remains an important task.

### 6.3.2 Low noise cantilever blade springs and improved suspension thermal noise

There are a variety of techniques being explored which could improve the room temperature thermal noise of the suspensions [275]. It is possible to improve dissipation dilution by increasing suspension length or thickening fiber ends to enhance energy distribution (e.g. 5 mm stock rather than 3 mm stock).

Studies are also underway to understand how to lower the first ‘bounce’ mode of the test mass. Development of fused silica or silicon blade springs which could be incorporated in the final monolithic stage for improved vertical isolation compatible with lower thermal noise is an attractive option to explore for possible upgrades to Advanced LIGO and future interferometers. Sapphire is also a possible material choice. Experiments are already underway to investigate the breaking stress for such materials when used as blades. In addition to robustness tests of silicon/sapphire as a spring material, work on protective coatings, mechanical loss and thermal conductivity also need to be pursued.

### 6.3.3 Silicon and Sapphire Suspensions

Silicon has attractive thermal and thermo-mechanical properties making it a strong candidate for the suspension elements in future detectors possibly operating at cryogenic temperatures to reduce thermal noise. It is also conductive which may have advantages for controlling charging effects (discussed elsewhere). Development and measurement of suitable suspension flexure elements, including studies of the optimum material, thermal noise properties, and the geometry and assembly of elements including methods of bonding to test masses are being pursued [235]. Analysis techniques include the use of FEA to study the various

contributions to thermal noise such as surface loss and bond loss. Investigation of fabrication techniques, properties of silicon-silicon bonds such as strength and thermal conductivity and thermo-mechanical properties of silicon, for example as a function of doping, are examples of areas which can be addressed. Sapphire should also be studied as this is a potential backup/alternative material, having the benefit of low thermal noise but allowing the use of 1064 nm transmissive optics. This is the material of choice in KAGRA.

Suspension elements are being fabricated via two methods (i) laser heated pedestal growth, and (ii) mechanical fabrication from wafers. The laser heated pedestal method is being pursued in both sapphire and silicon. With this technique surface tension is used as the “crucible”, resulting in the purest fibers. It is expected that the strength and thermal conductivity of these pristine fibers should be high, and this will be tested once samples are available. Sapphire is an easier material in principle to grow as it absorbs efficiently at 10.6 $\mu$ m, from the CO<sub>2</sub> heating laser. Silicon undergoes a “metallic” phase transition resulting in a change in emissivity during melting, and this might prove an interesting control problem to maintain the melt temperature. Mechanical etching/machining silicon ribbons has already been used to make short (5 cm) long suspension fibers. This technique has been shown to reduce fiber strength by approximately one order of magnitude (300 MPa) compared to the pristine material (4 GPa). This is likely to be caused by a damage layer. Oxidizing/ wet etching the samples has been shown to improve the strength by a factor of 2, possibly due to healing these damaged layers.

As an alternative protective coating, diamond-like carbon (DLC) has been investigated alongside thermally-grown oxide coatings and magnetron sputtered silica coating. In addition, the effect of edge polishing the silicon flexures and argon plasma pre-deposition treatment was investigated [276]. It is likely that silicon surface quality is the dominant factor in determining both its tensile and flexural strength. However, it was shown that application of a multilayer DLC coating combined with an extended pre-deposition argon etching process can increase average tensile strength by around 80%, with some outlying data points suggesting yet greater improvement is possible. Edge polishing appears to slightly degrade the strength of silicon, likely due to creation of edge defects; however, application of a 3 $\mu$  DLC multilayer coating can more than compensate for these effects, possibly due to the conformal nature of the coating process and the coating filling in cracks or chips introduced by polishing. Additionally, DLC coatings on silicon were shown to provide a high level of protection from abrasion, which may facilitate easier handling / assembly of silicon suspension components, making them generally more robust. Similar coatings could also be considered for relevant silica and sapphire components. Thermal noise calculations showed that replacing the current maraging steel blades of aLIGO with DLC-coated silicon blade springs could offer a 6.4-fold reduction in vertical thermal noise at 10 Hz, with significant improvement across all frequencies above 3.5 Hz

There is currently some activities underway in both the UK and Australia to develop small scale silicon prototypes. In the UK the proposal is to develop a 1 kg suspensions which can be cooled in a cryogen-free cryocooler, while in Australia the plan is to suspend a small scale optic from a torsion fibre. Both of these activities are complimentary and a very useful first step towards understanding the challenges faced when building up suspensions from crystalline materials.

The collaboration should also look at opportunities to operate larger prototypes at 123 K. For example, there has been some discussion about the possibility of turning the Caltech 40 m prototype into such a testbed.

**Attachment techniques** A slightly modified version of silicate bonding for silicon-silicon attachment for, e.g., the attachment of interface pieces to silicon test masses is well underway. Strength measurements at both cryogenic and room temperatures of these bonds has shown it is a viable attachment technique and investigations are ongoing to further understand the influence of different parameters on strength like the nature and thickness of the oxide layer required. Bond loss measurements on silicon test masses have also started producing initial results.

To interconnect the silicon suspension elements the technique of hydroxide catalysis (or silicate) bonding can be used. It is essential to quantify the mechanical loss of the bond region created between the silicon (or sapphire) substrates in order to calculate thermal noise associated with the bonds. The loss of the bonds can be estimated from the measurement of the Q-factor of a mechanical oscillator, fabricated from silicon elements using hydroxide catalysis bonding. Groups from Moscow and Glasgow fabricated tuning forks from silicon ribbons, which were then silicate bonded. They have found the upper limit of the bond loss at the temperature of 123 K [277]. Further research is aimed at improving the technology of silicon ribbon fabrication and their bonding for creation of prototypes of quasi-monolithic silicon suspension.

## 6.4 R&D for LIGO Cosmic Explorer

The baseline concept for Cosmic Explorer is to use the Voyager Technology for the isolation of the optics. The technology envisioned for LIGO Voyager should be of such performance that the direct seismic coupling is below the noise of a cryogenic mass, and that the motion at 10 Hz is dominated by Newtonian Noise, even though local seismic arrays are used to eliminate at least 90% of that noise.

Sensitivity studies are now well underway to understand some of the design implications for future 3rd generation detectors. These include the European Einstein Telescope [278] and Cosmic Explorer [279]. The Einstein Telescope proposes a xylophone configuration, operating both low temperature and room temperature instruments in the same facility, in a triangular configuration. At a recent Einstein Telescope meeting in 2018, the science team is also looking at the trade-offs of operating the instrument in a single 123 K configuration. Thus there is good alignment between this design and that of Cosmic Explorer.

In the previous sections the main technologies for dealing with the 3rd generation technologies have been identified, and below we refer to the areas which require active research. It is important to note that technologies typically take 10-15 years to reach maturity, from lab scale demonstration to an engineered solution which can be installed at a site, and thus for a >2030 timeframe this R&D must be developed now.

### 6.4.1 Seismic Noise

As the coupling of vertical motion to the sensitive direction of the GW detector increases linearly with detector length (due to the curvature of the Earth), the GW strain resulting from a fixed vertical displacement noise level is insensitive to detector length [279]. The active seismic isolation concepts and systems developed for Advanced LIGO will be adequate to support these new suspensions, though inertial sensors and tilt sensors with lower noise will be necessary if the suspension modes were reduced to lower frequencies. The capacity of the inertial seismic isolation systems and suspension must be adequate to support any increase optic mass and care should be taken to ensure low frequency resonances of the structure do not pollute the detection band.

The implementation of beam rotation sensors, to monitor tilt and to correct the horizontal inertial sensors, will be an important ongoing development effort.

### 6.4.2 Suspension Thermal Noise

Current research into test-mass suspensions is focused on supporting larger masses to reduce the challenges of quantum noise, and longer suspensions for reduced thermal and seismic noise both in the horizontal and vertical directions. Vertical thermal noise can be further reduced by lowering the vertical resonance frequency of the last stage of the suspension, possibly by introducing monolithic blade springs into the suspension designs. Research needs to focus on robust techniques to engineer these larger suspensions. This includes techniques to handle masses of the order 150-300 kg (see Table 2), robustly bond and fabricate suspension fibres (both fused silica and crystalline silicon/sapphire), and build a suitable protocol for mating/installation to the inertial seismic isolation system.

### 6.4.3 Material Characterisation

There is much effort in the collaboration devoted to measuring and characterising the thermo-mechanical properties of fused silica, sapphire and silicon. This effort needs to continue as the most robust thermal noise models require a full understanding of the suspension, including thermal conductivity, optical absorption, mechanical loss and tensile strength. This is particularly the case for the test mass optics, suspension fibres and potential spring materials which may move away from maraging steel (due to its higher loss angle of  $10^{-4}$ ).

### 6.4.4 Cooling Strategies

There is significant work that needs to be put into techniques to cool the next generation mirror suspensions. While this is an issue that must be approached for LIGO Voyager, the larger scale of the test mass may add some additional complexity. There are ongoing activities in both the US [280] and Brazil [281] to understand the technical challenges of cooling to both liquid nitrogen and liquid helium in a seismically quiet environment.

### 6.4.5 Environmental Noise

While there is still uncertainty about the location and fabrication strategy for a Cosmic Explorer type instrument, it is important to note that the collaboration should be working on mitigating environmental noise as a high priority. This includes the choice of seismically quiet sites, sites with low wind noise (if systems are located above ground), and low microseismic and anthropogenic noise.



## 7 Lasers and Auxiliary Systems

The Lasers and Auxiliary Systems working group (LAWG) developed out of the Lasers working group. In addition to all types of *classical* lasers (squeezing is part of the Quantum Noise WG), this group now includes auxiliary systems which encompasses all technologies which are not part of any of the other working groups.

The following sections are organized in terms of the needs and requirements for lasers and auxiliary systems according to the road-map (section 2) It lists the identified technologies and research areas with respect to aLIGO, A+, LIGO Voyager and LIGO Cosmic Explorer and closes with research areas and R&D that is relevant to all envisioned detectors.

### 7.1 Achieving Advanced LIGO Design Sensitivity

#### 7.1.1 Advanced LIGO PSL

The aLIGO pre-stabilized laser system (PSL) was designed to use the following components [282]. A four stage Nd:YVO amplifier system is used to increase the 2 W power of a Nd:YAG non-planar ring-oscillator (NPRO) to 35 W [283]. An injection locked Nd:YAG end-pumped rod system was chosen as the high power oscillator. It was shown that an output power of more than 200 W in a linear polarized single spatial and frequency mode with such an aLIGO laser system [284] can be emitted. The laser has been developed and built by the GEO group in Hannover (Laser Zentrum Hannover (LZH) and Max-Planck-Institut für Gravitationsphysik / Albert-Einstein-Institut AEI). All PSL systems have been installed and were running nominally in the beginning, but the high-power oscillators were replaced in Fall 2017 at LLO and Spring 2018 at LHO with a second amplifier stage because an unexpected high beam-jitter coupling of the Hanford interferometer and due to problems to repair the LLO high power stage after a human interaction caused significant damage to the system. The switch to the second amplifier stage reduced the maximum output from 200 W to 70 W

The need remains to reduce the noise of the PSL and to investigate noise sources. Beam jitter fluctuations and resulting amplitude fluctuation can be traced to couple into the GW channel.

Continuing technical support for the PSL system, to address problems that arise from continuous running over long times and to further monitor the long term stability and operating parameters of the laser is necessary. This includes maintaining and improving the noise of the PSL based on the second stage 70 W amplifier, as well as working towards a fully operational 200 W stage based on the high-power oscillator. In parallel alternative amplifier based options to generate a low noise 200 W laser need to be developed in case the beam jitter and repair problems of the HPO can not be solved. Promising candidates are high power fiber amplifiers or coherently combined solid state amplifiers.

## 7.2 A+

Small upgrades to aLIGO are possible between major observing runs, given sufficient technical readiness and scientific payoff. The following section lists R&D areas that need continuous efforts for the A+ time frame.

### 7.2.1 PSL upgrades

The need remains to reduce the noise of the PSL and investigate noise sources. Beam jitter fluctuations and resulting amplitude fluctuation can be traced to couple into the GW channel.

In the ongoing process of commissioning the aLIGO PSL it became apparent that maintaining a stable, but one of a kind, PSL is a demanding task. Alternative laser designs with lower noise might and/or better longevity are required for A+. Solid-state amplifier systems are possible laser candidates with potential better noise characteristics. Such systems need to be investigated.

The task and goals for A+ are to provide continuing technical support for the PSL system, to address problems that arise from continuous running over long times and to further monitor the long term stability and operating parameters of the laser.

### 7.2.2 Faraday Isolator in Squeezing Systems

Faraday isolators are required to separate the counter-propagating beam from the incoming beam. The aLIGO Faraday isolators use TGG as the Faraday material. The TGG rotates the polarization angle by an amount proportional to the length of the crystal, the Verdet constant, and to the applied magnetic field. The main issues with the Faraday isolator are beam distortion due to laser heating and subsequent thermal lensing, a reduction of the optical isolation due to depolarization and changes in the temperature dependent Verdet constant. This is further complicated by the fact that the FI is usually placed inside the vacuum chamber following the suspended input optic mode cleaner.

The power handling capabilities of the output Faraday isolator are far less critical. However, the optical losses inside the Faraday would currently limit the amount of usable squeezing. Since squeezing is one of the leading ideas to 3<sup>rd</sup> generation detectors, any improvement in the optical losses could directly improve the range of these detectors.

Figure 26 shows how the losses in the squeezed beam path are strongly affecting the amount of squeezing detectable in the interferometer and therefore the improvement in the sensitivity. With 10 dB of squeezing injected, for instance, the losses need to be less than 20% in order to be able to detect at least 6 dB of squeezing.

In a typical layout for injecting squeezing in the interferometer the squeezed beam needs to pass through not only the main output Faraday of the IFO (twice), but also at least one additional Faraday to isolate the interferometer from the squeezer and mitigate noise from back scattered light. The losses of the new Advanced LIGO Faraday have been measured to be  $\sim 4\%$ . Similar measurements in initial LIGO Faradays give losses of 4 - 6%. With the current losses the Faradays themselves will account for about 15% of the losses in the

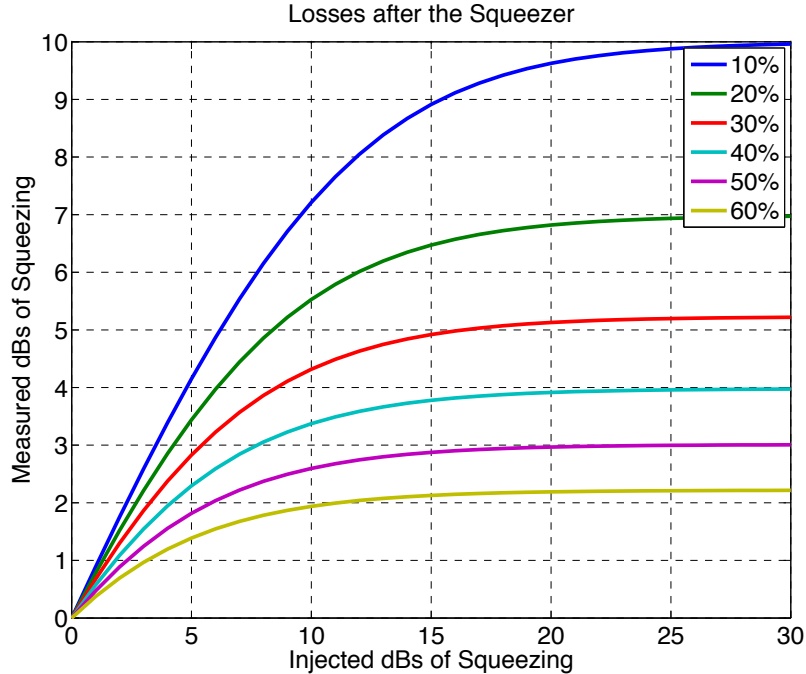


Figure 26: Measured squeezing as function of the injected squeezing for different levels of losses in the squeezed path.

squeezed path. In order to maximize the benefit from the injection of squeezed light, it is important to reduce the losses of a single Faraday to 1 - 2%.

For the main dark port Faraday, it is also important to reduce the amount of light which leaks into the squeezed beam path, as this is a source of noise once it is back scattered back into the IFO from the squeezer source. In the new Advanced LIGO Faraday about 0.5% of the light which passes through the Faraday is reflected by the thin film polarizer. By improving the mechanical design of the Faraday one can hope to reduce this percentage down to  $\sim 0.01\%$ .

### 7.2.3 Laser Stabilization

Power stabilization will probably be the most demanding laser stabilization task in future gravitational wave detectors. Technical power noise on the laser can couple via many paths into the gravitational wave channel. Advanced LIGO requires a relative intensity noise (RIN) of around  $10^{-9}/\sqrt{\text{Hz}}$  in the interferometer input beam. The accurate sensing of the needed 500 mW laser power at that location is difficult and the signal is still contaminated by pointing, polarization, and potentially even frequency noise. Ongoing research is needed to understand these couplings and reach the required stabilities.

## 7.3 LIGO Voyager

The LIGO Voyager design mitigates limiting noises of aLIGO by allowing moderate changes to the LIGO detectors without affecting the vacuum envelop but with possibly including

cryogenic temperatures of 120 K. The following section describes the R&D related to those improvements.

### 7.3.1 PSL LIGO Voyager

The road map for the general description in section 2. The designs call for laser wavelengths of  $1.5\ \mu\text{m}$  and longer to be compatible with silicon optics at cryogenic temperatures. As the final sensitivity depends on the power inside the interferometer the input power can be traded against power recycling gain without affecting the fundamental noise limits. However, the power recycling gain is limited by the optical losses inside the arm cavities as well as the contrast defect at the beam splitter. Typical power levels for the lasers required for LIGO Voyager are envisioned to be about 200 W.

**High power concepts: 1550-1650 nm** Erbium doped fiber lasers and Er:YAG lasers emit between 1550 and 1650 nm where the absorption in silicon is expected to be very low. Commercially available erbium fiber systems include a master laser and a fiber amplifier and achieve output powers of 10 W in single mode, single frequency operation and higher power levels are expected in the near future.

The GEO group achieved more than 100 W of output power by an off-resonant pumping scheme. The Adelaide group has completed a high-resolution spectroscopic investigation of Er:YAG at cryogenic temperatures, where it is a 4-level gain medium, and demonstrated a 5 W cryogenic Er:YAG laser at 1617 nm. Power scalability as well as the free running noise performance and the robustness of this system needs to be investigated. Further stabilization concepts for these kind of amplifier systems have to be developed.

**High power concepts -  $2\ \mu\text{m}$  and longer** Another interesting region for gravitational wave detector wavelengths, especially to avoid 2-photon interaction in silicon, is  $2\ \mu\text{m}$  or longer. Ho:YAG and Tm:YAG lasers are possible candidates but more research is needed to find the most suitable laser medium and to demonstrate the feasibility of a  $2\ \mu\text{m}$  PSL.

**Additional Requirements** Many different applications drive the laser development worldwide and it appears that nearly every year improved laser systems become commercially available. However, there is currently no application which has similar stringent requirements on the temporal and spatial stability as gravitational wave detectors. Hence a laser development programs for 3<sup>rd</sup> generation detectors is needed to design and build a reliable laser with sufficiently low free-running noise, an appropriate spatial beam profile and good controllability. Programs that have been initiated are going to require continued support.

### 7.3.2 Photodiodes

Advanced LIGO is currently using four in-vacuum photodiodes in parallel to measure the required 500 mW of light [285]. This is a sub-optimal arrangement for several reasons including reliability and alignment issues. To get a quantum limited measurement of the

power fluctuation of 500 mW of light, new photodetectors need to be developed with sufficient power handling capability, spatial uniformity and quantum efficiency. First experiments showed that back-illuminated InGaAs diodes show promising features. However neither the spatial uniformity nor a sufficiently high quantum efficiency has been demonstrated so far. Furthermore current power stabilization experiments seem to be limited by  $1/f$  electronic noise in photodiodes. The origin of this noise needs to be better understood and either the noise source has to be reduced or easily applicable selection criteria need to be found to get the best devices from the available vendors.

High efficiency photodiodes are especially important when using squeezed light, as described in section 4.3.1 of the quantum noise research. Further R&D in close collaboration between the material and device experts, electrical engineers and groups that can test the photodiodes is needed to develop better photodiodes for LIGO Voyager and later generation gravitational wave detectors, especially with respect to the longer  $1.5\ \mu\text{m}$  to  $2\ \mu\text{m}$  wavelength range where quantum efficiency is currently inferior to  $1\ \mu\text{m}$  systems.

### 7.3.3 Electro-Optic Modulators

Length and alignment sensing schemes rely heavily on the generation of optical sidebands which co-propagate with the carrier field into the interferometer. These sidebands are currently generated by RTP-based electro-optic modulators which withstand several 100 W of continuous laser power without degrading the beam profile. LIGO Voyager and LIGO Cosmic Explorer detectors are likely to work at different wavelengths and/or at higher power levels for which suitable electro-optic modulators are not yet available or have not been tested. The main problems encountered in high power applications are photo refractive damage and variations in optical path length across the beam profile caused by the residual absorption of the laser beam.

Photo refractive damage has a fairly well defined threshold in specific nonlinear crystals and can be increased by doping the crystal. The most promising family of crystals in the near infrared region are crystals belonging to the  $MTiOXO_4$ -family such as RTP;  $M$  is an alkaline metal such as K, Rb, or Cs, and  $X$  is either P or As. These crystals have fairly large electro-optical coefficients, good thermal properties, and, in principle, very low optical absorption coefficients between  $1$  and  $2\ \mu\text{m}$  laser wavelength. Optical absorption in the  $MTiOXO_4$ -family increases at lower wavelength and potentially limits the laser power to a 10's of watts for visible lasers.

$\beta$ -barium borate (BBO) and its derivatives are often used in the visible and near-UV region of the spectrum. BBO is uniaxial and has a very high damage threshold. Values larger than  $3\ \text{kW}/\text{cm}^2$  for cw-light have been quoted by multiple vendors. It's negative thermo-optical coefficient prevents self-focusing. However, the electro-optical coefficient is low compared to other electro-optical crystals and BBO appears to be of limited value unless the laser wavelengths is reduced to well below  $500\ \text{nm}$ .

Magnesium-oxide doped lithium niobate ( $\text{MgO}:\text{LiNbO}_3$ ) might be a potential alternative for IR lasers. The doping increases the photo refractive damage threshold significantly [286] and the crystal has a  $\sim 10\%$  larger modulation coefficient than RTP. However, going to laser powers beyond  $1\ \text{kW}$  and laser wavelength below  $1\ \mu\text{m}$  requires significant testing of

electro-optical materials to ensure that electro-optical modulators will be available for LIGO Voyager and future detector generations.

### 7.3.4 Faraday Isolators

The aLIGO Faraday isolator uses two TGG crystals, a quartz rotator, and a waveplate to compensate the depolarization inside the TGG crystals. It is optimized to suppress thermal effects that would distort the laser mode. Scaling this to kW-class power levels requires further reductions in the optical absorption in TGG while the large number of components already increases the number of ghost beams significantly. Other options are to increase the Verdet constant by cooling the crystal, to use stronger magnetic fields, and to use other magneto-optical materials such as GGG, CdMnTe.

One recent approach used to increase the magneto-optical effect is doping TGG single crystals and ceramics, or other garnets with rare-earth elements such as  $\text{Ce}^{3+}$ ,  $\text{Dy}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Pr}^{3+}$ . This method improves the Verdet constant by 20%-70%, keeping absorption levels similar to TGG single crystals. Terbium-doped  $\text{Y}_2\text{O}_3$  ceramics investigated in the 380-1750 nm range showed more than 3 times higher Verdet constant than TGG crystals and ceramics. High power studies are needed for a complete evaluation of these materials.

Increased absorption in TGG will likely prevent the use at wavelength longer than  $1.3 \mu\text{m}$  and in addition make the crystals very long. At  $2 \mu\text{m}$  the crystals would need to be twice as long as the 12 mm crystals used in aLIGO, but the telecommunication sector developed ferromagnetic rare earth iron garnets (RIGs) such as yttrium iron garnet (YIG) and more recently  $\{\text{BiRE}\}_3(\text{FeGaAl})_5\text{O}_{12}$  to rotate the polarization. Unlike paramagnetic Faraday materials, these ferromagnetic materials can be magnetized such that they don't require any external magnetic field. These materials are typically grown in sub-mm thick films on lattice-matched substrates such as GGG for 45 deg rotation. However, the absorption is still in the 1-10 ppm/cm range at interesting wavelengths which prohibits high power laser operation. At shorter wavelength, optical absorption increases in all materials pronouncing thermal effects.

Faraday isolators using potassium dihydrogen phosphate (KDP) and its isomorphs have been developed for the UV but the absorption is still fairly high. Very promising are Ce:YIG films grown on GGG substrates with different crystallographic orientations, showing high Verdet constant of -5800 deg/cm at 1550 nm, and low absorption coefficient at this wavelength and longer. Another alternative are thin ferromagnetic films such as iron or cobalt. They show a strong Faraday effect in the visible and near infrared, but currently the strong optical absorption of these type of materials needs further optimization.

A targeted research program to study these materials at higher power levels at all interesting wavelength ( $1.06 \mu\text{m}$ ,  $1.5 \mu\text{m}$  and  $2 \mu\text{m}$ ), is required to develop Faraday isolators for the next generation of gravitational wave detectors.

### 7.3.5 Cryogenics

It has long been known that cryogenically cooled test masses can have much improved material parameters which lead to significant reductions in thermal noise. However, operating at

cryogenic temperatures presents multiple new challenges which need to be addressed. The most pressing is to find ways to cool the temperature, to isolate the mirrors from their hot surroundings and to constantly extract the deposited laser heat without short-circuiting the suspension and seismic isolation system. Detailed thermal models have to be developed and tested to maximize radiative and conductive cooling paths.

An additional challenge is the strong possibility of contamination through condensation on the surfaces. Methods will need to be developed to (i) mitigate the level of contamination in cryogenic mirrors, (ii) quantify the magnitude and type of contaminants, and (iii) if necessary, clean contaminated mirrors in situ. One idea which should be pursued is to use fs laser spectroscopy to identify the contaminant and potentially also to remove it.

**Cryogenics: Radiative Cooling** Operation at cryogenic temperatures poses formidable challenges including heat extraction from the cooled test masses, required both under steady state operation and for cooling from room temperature in a reasonable time. The system needs to work without adding noise or short-circuiting the mechanical isolation. In the steady state, the circulating power may be in the range 0.5 to 3 MW, and with anticipated coating losses of 0.5 to 1 ppm, power loss in the arm coatings is of order 0.3 to 3 W per optic. For cooling, a reasonable estimate is between 2 and 100 W of heat conduction from the test masses to the cold environment.

Studies are underway of a novel method of heat removal: near-field radiative coupling between two objects: one hot and one cold. The basic idea is that many thermal fluctuations in the hot object do not couple to radiation; instead, they produce evanescent fields outside the object. If a cold object with appropriate properties is introduced into this evanescent field region, energy is transferred, cooling the hot object. This approach is potentially capable of removing more than 200 W from a test mass. The heat transfer can be greatly enhanced using a small gap but this is accompanied by force coupling and this effect needs to be taken into account. Room-temperature experiments to explore this method of heat transfer have been observed and are characterizing in detail the heat transfer in the near-field regime. Cryogenic experiments are planned, as are measurements to determine the effects of coatings on the heat transfer, and to attempt to optimize the coatings for maximum transfer with spacing around 0.1-1  $\mu\text{m}$ .

## 7.4 LIGO Cosmic Explorer

The current LIGO facilities, while extraordinary in their capabilities, present significant limitations to potential long term improvement. The Ultimate stage will allow for a complete redesign and major expansion of the detector. The following topics are future developments within this project phase.

### 7.4.1 PSL for LIGO Cosmic Explorer

Most ideas regarding the LIGO Cosmic Explorer detectors target the low frequency performance where the sensitivity is limited by radiation pressure noise, thermal noise and Newtonian noise and not by shot noise or the available laser power. The laser needs for

silicon test masses for these detectors can probably be met by LIGO Voyager PSLs assuming that the lasers will be developed in time, except for an increase in output power.

**High power concepts: 1550 nm to 2  $\mu$ m and longer** The main difference between Voyager and Cosmic Explorer is that the typical power levels for the lasers required will go up from several 100 W up to 1 kW. To further explore the high frequency region of the GW-spectrum, further increases of the laser power will be required and the laser concepts as outline in the Voyager section will need to be brought to the 1 kW level.

If new test mass materials other than silicon become available changes towards a shorter wavelength might be advantageous. Shorter wavelengths generally require thinner coating layers which could reduce coating thermal noise; however, it should be kept in mind that shorter wavelength also generate smaller beam sizes for identical cavity g-parameters. The general scaling of coating thermal noise with thickness of coating layers and beamsizes is in first order independent of the wavelength (for identical cavity g-parameters). The most obvious way to reduce the wavelength is via frequency doubling a powerful 1064 or 1030 nm laser but many other laser concepts are currently being explored by industry and many scientific institutions.

**High power concepts - 1030-1064 nm** In the future, if kilowatt class lasers become necessary for frequency doubling, Yb doped YAG, which operates at 1030 nm, could replace the Nd system because of its higher efficiency, lower quantum defect, better thermal management and potentially longer-lived laser diode pumps. Its main disadvantages are that it is a quasi-3-level system at room temperature and thus more sensitive to increased temperatures within the gain medium, and that it has a much lower pump absorption coefficient. However, at cryogenic temperatures, the quasi-3-level system turns into an efficient 4 level system. The Adelaide group demonstrated already 200 W, with diffraction limited beam quality for all power levels, requiring no thermal compensation, at 1030 nm from a cryogenic Yb:YAG laser using a single zig-zag slab [287]. There is a substantial commercial interest driving the development of both Yb solid state and fiber lasers, amplifiers and their pump diodes for very high power applications.

Another option is the coherent beam combination of two or more 200W class lasers. Either solid state high power systems or Yb:doped fiber amplifiers with identical or independent seed laser sources are potential candidates to generate the 200 W output power. The characterization and stabilization of coherently combined systems at intermediate and high power levels should be investigated. The beam combination option is a possibility for all wavelengths that are mentioned in this section.

**High power concepts - spatial mode filtering and adaptive optics** To convert distorted laser beam profiles into the target eigenmode of the interferometer either static or dynamic wave front correction systems or passive filtering will be required. Advanced LIGO uses optical cavities (the mode cleaner) to filter the fundamental 00-mode and suppress all higher order modes. This technology is slowly reaching its limits as the high power build up inside these optical cavities starts to distort the filter cavities themselves. Other spatial



modes such as Laguerre Gauss 33-modes or Mesa beams require modified filter cavities which are resonant only for these specific spatial modes.

For higher power levels intrinsic problems are expected with the filtering method and hence dynamic adaptive beam correction methods should be designed. These could be based on well known Shack-Hartmann detectors and adaptive optic techniques currently employed in astronomical telescopes. These techniques have also significant commercial potential for many other high power laser applications.

## 7.5 General R&D

The R&D projects listed in the following section require constant improvements and are relevant to all envisioned detector generations, from A+ to LIGO Cosmic Explorer.

### 7.5.1 Auxiliary Lasers

Auxiliary lasers serve several functions in interferometric gravitational wave detectors.

- CO<sub>2</sub> lasers at 10  $\mu\text{m}$  are used to write a heating pattern into the compensation plate placed next to the ITM.
- Diode lasers at various wavelengths are used together with Hartmann sensors to sense thermal deformations in the test masses and the compensation plate.
- Frequency doubled Nd:YAG lasers are injected at the end stations for lock acquisition of interferometer length degrees of freedom.

The status and planned R&D on these laser types is described as part of the subsystems they are used in: CO<sub>2</sub> lasers as part of the TCS actuation in Section 7.5.2, diode lasers for Hartmann sensors as part of TCS sensing and control in the AIC working group.

### 7.5.2 Thermal Correction System

The goal of the Thermal Correction System is to optimize the spatial mode inside the interferometer. This spatial mode can be degraded by imperfections in the mirrors caused by radii of curvature or surface figure errors as well as non-homogeneous heating of the optics by the science beam. Untreated, this will reduce the mode matching between the two arm cavities and the recycling cavities and change the beam size inside the interferometer. Advanced LIGO uses ring heaters to optimize the radii of curvatures of the ITMs and ETMs and CO<sub>2</sub> lasers to compensate the thermal lens in the ITM substrate by acting on the compensation plate [288].

**Ring Heater** The ring heater has to meet several requirements. It has to generate a homogeneous heating profile with minimal heating of the suspension structure. Its location very close to the test masses requires that it has to meet very stringent cleanliness requirements. Advanced LIGO is currently using a ringheater which uses nichrome wire wound around a

bent glass rod, an alternative design developed by the UF group sandwiches the nichrome wire between two alumina coated aluminum surfaces [289]. Both heaters are embedded inside a gold coated thermal shield to maximize the heat transferred to the mirror and minimize the radiative heat transfer into the suspension.

The heat loss at the end points due to thermal conductivity generates an asymmetric heating profile which will also distort the mirror surface and the thermal lens inside the substrate. At higher heating powers, these asymmetries will start to reduce the optical build-up inside the arm cavities and increase the contrast defect at the beam splitter. The development of ring heaters which produce more symmetric heating profiles could already be crucial for high power operation in Advanced LIGO.

One way to reduce the heating by the ringheater or even eliminate it is to coat the barrel of the optic with a thin layer (a few microns) or an IR reflecting metal such as gold [288]. This would reduce the radial heat flow and homogenize the temperature distribution inside the substrate. Adding a gold barrel coating to the optics would have implication for other aspects of the design, notably thermal noise, charge mitigation, and parametric instabilities. Measurements of the mechanical loss of a thin gold coating indicate that the gold coating can be applied without adversely affecting thermal noise. Gold coating applied to the barrel for thermal compensation purposes might not reduce the optics modal Q's enough to cause significant improvement in parametric instability performance. Tests of a gold coatings interaction with possible charge mitigation schemes, including UV, should be explored. Results of these tests might require follow-ups with other materials and/or coating methods or with additional modeling. This technique may be ready for use in a 3<sup>rd</sup> generation detector.

**CO<sub>2</sub> laser** Power fluctuations of the CO<sub>2</sub> laser are one of the dominant noise sources associated with the TCS. Commercially available CO<sub>2</sub> lasers do not meet the stringent requirements on power stability for Advanced LIGO during high power operation (125 W input power in the science beam) and R&D has started to develop better CO<sub>2</sub> lasers for Advanced LIGO.

A scanning (or, more generally, a directed-beam) thermal compensation system that can vary the compensation profile in real time without injecting noise into the signal band would be very valuable to correct non-radial symmetric beam distortions. Such a system could already be important for high power operation in 3<sup>rd</sup> generation detectors. This will require research on carbon dioxide or other potential heating lasers, to reduce noise and possibly boost power, and potentially on measurement and control issues. In addition, by moving to shorter wavelengths it might be possible to develop MEMS or other technology based spatial light modulators to allow a programmable heating beam profile.

### 7.5.3 Active Wavefront Control

Several technologies, notably squeezing, assume improved mode-matching (see table of losses in 9) is available for these systems. The limits assumed in these scenarios are beyond static polishing and placement uncertainties for mode-matching optics. Therefore, active wavefront control (AWC) is required to maximize mode-overlap between cavities. This control is in addition to the Thermal Correction System. Although there is some overlap between the

two systems, TCS is predominantly needed for correcting dynamic changes wavefront errors and AWC is used for correcting static wavefront errors. Future adaptive control research is needed to address:

- A possible redesign of SRC/PRC configuration to improve mode-matching actuator orthogonality.
- Adaptive mode-matching on output optics addressing thermally adaptive optics for low spatial frequency curvature errors (astigmatism and defocus) and higher spatial frequency (point sources and polishing errors). Extreme mode-matching targets of larger than 99.5% may require correction of polishing errors beyond spherical correction.
- Other adaptive optic techniques capable of macroscopic motion of the order of tens of centimeters within the output chain.
- Improved sensing techniques for SRC/IFO, SRC/OMC and FC/IFO mode-matching.
- General adaptive optics for in vacuum use. A generic, low noise, high range small adaptive optic element, either mechanical or thermal based, is needed.

#### 7.5.4 Beam Shaping

Mirror thermal noise is one of the fundamental factors limiting the sensitivity of gravitational wave detectors. A Gaussian beam profile is not the best shape to average over thermal fluctuations and different, carefully chosen shapes allow for sensitivity improvements. Non-spherical mirrors, shaped to support flat intensity mesa profile beams, have been designed and fabricated using specialized coating techniques. These mirrors are being tested on a dedicated interferometer to assess ease of mode-matching and locking. Recent efforts have shown that the tilt sensitivity of the fundamental mesa mode agrees with expectations. It is possible to extend this study, producing useful alignment correction signals via the wavefront sensing technique. The Sidles-Sigg tilt instabilities must also be examined. In addition, continued modeling needs to examine how thermal effects alter the mode profile in a detector arm cavity and help develop thermal compensation strategies. One option involves depositing a static thermal compensation profile to mitigate these effects.

Modeling and experimental work is being carried out on Laguerre-Gauss and other optical modes that show promise for reducing thermal noise. Laguerre-Gauss modes may avoid some of the instability issues that cause concern with mesa beams. There are, however, questions about the strict requirements on the figure and polish of the optics necessary for these higher order modes. There has been modeling of the effects of different beam shapes on parametric instabilities. Further modeling and experimental testing will be necessary to truly evaluate the potential and limitations of these beam shapes.

## 8 Control Systems

The Control Systems Working Group (CSWG) covers fundamental and applied research in control systems as it relates to GW interferometers, including:

- system identification
- modeling
- synthesis
- analysis
- optimization
- performance assessment
- hardware and software implementation

The role of the CSWG is unique within the LSCs instrument science working groups. The use of control systems is pervasive within, and enabling to, the work of many of the other instrument science working groups. In addition to supporting its own fundamental research in cutting-edge control system techniques, the CSWG should support and enable the research of other LSC WGs. In particular, as next generation detectors are researched, control system considerations should be included at the very beginning as an integral aspect of the system.

The following sections describe the control system challenges in implementing the roadmap, and then R&D focus areas for the group which are relevant to all of the anticipated detectors. Although the discussion concentrates on Interferometer sensing and control, there are similar control challenges in many of the other interferometer subsystems as well, such as our principal "actuator" systems - SEI, SUS and TCS.

### 8.1 Interferometer Sensing and Control (ISC)

The Interferometer Sensing and Control (ISC) consists of four areas (LSC, ASC, CAL, and CDS). The dominant issue of ISC is to maintain IFO stability while not introducing control noise. The following is an overview of the five subsystems:

**Length Sensing and Control (LSC):** Length sensing and control refers to management of resonance conditions of the fundamental degrees of freedom in the suspended interferometer while maintaining a high sensitivity to gravitational waves. The longitudinal distances between the mirrors (or the lengths of resonant cavities) are controlled to minimize length changes.

**Alignment Sensing and Control (ASC):** The ASC primarily controls the 1st-order optical modes in the interferometer by adjusting the angular orientation of the mirrors. In addition to the LSC, this system is also necessary to maintain interferometer sensitivity. Although the angular motion of the mirrors are locally stabilized by the vibration isolation and the suspension systems, the global control is necessary in order to realize stable interferometer operation.

**Calibration (CAL):** The calibration of the detectors' LSC degree of freedom most sensitive to gravitational waves, the differential arm length (DARM), is the essential interface between the instrument and the consumers of its astrophysical output. This precision engineering involves detailed, highly precise, and accurate characterization of each component of the loop from aspects as complex as the detector's opto-mechanical response to differential arm length changes, down to the simplest electronics inside anti-aliasing chassis. This must be done over the entire, many kHz, frequency band of the detector and its validity must be maintained over  $\mathcal{O}(\sim 1 \text{ yr})$  duration observation runs. At its core, this requires development of an extremely stable, precise, and accurate set of references from which the calibration is derived.

**Controls and Data Systems (CDS)** The computational system involved in sensor data receiving and control distribution is an critical part of any modern control system. In this context, subjects covered below refer to research and development that includes improving the data acquisition system for specialized needs and advanced control topologies as well as taking full advantage of the existing capabilities of the digital infrastructure.

### 8.1.1 Length Sensing and Control (LSC)

The Advanced LIGO interferometers consist of five coupled cavities that must resonate simultaneously to reach the operating point. The LSC receives the length sensing signals from the photodiodes and sends commands to the actuators on the suspensions so that those five degrees of freedom stay at the operating point. Between the sensing and the actuation, the signals are processed by servo filters so that stable feedback control of each loop is established. The LSC system for the Voyager and Cosmic Explorer detectors will also rest on the above technologies.

The aLIGO LSC employs a combination of RF heterodyne detection and DC readout [290]. This RF sensing scheme comprises the Pound-Drever-Hall technique, Schnupp modulation, and third harmonic demodulation technique [291] as a baseline design. The DC readout scheme provides sensing of the GW signal with shot noise performance superior to that of RF detection, along with more immunity to laser noises when combined with an Output Mode Cleaner (OMC). DC readout also clears the way for the use of squeezed light injection to reduce quantum noise.

These sensing techniques have been independently demonstrated by the Enhanced LIGO [30, 292] and GEO600 [84, 293] interferometers, and other prototype interferometers [294–296]. The task for the aLIGO interferometer commissioning team is to integrate these well known techniques and find any hidden issues. As a scale model of the full aLIGO interferometers,

the Caltech 40m prototype plays an important role in this effort by giving us a first look at potential problems.

An example of an upcoming ISC challenge is the filter cavity which will be required for filtering squeezed light before it enters the interferometer. This cavity will require length control via a carrier frequency which is not the same as the carrier light circulating in the interferometer (to avoid polluting the squeezed state with scattered photons). The filter cavity will also present an alignment control problem similar to the input mode-cleaner. Similarly, the phase of squeezed light source itself must be controlled with respect to the interferometer.

For LIGO Cosmic Explorer, the idea of the Suspension Point Interferometer [241] may also be investigated as it could be a practical solution to mitigate the vibrational noise caused by the heat link of the cryogenic cooling system.

**Lock acquisition** A new feature of the aLIGO LSC is the use of independent laser sources in the Arm Length Stabilization (ALS) system [297–299]. This system enable us to lock the long arm cavities in a deterministic way with the guidance of auxiliary beams injected from the end mirrors.

Advanced LIGO achieves lock using two single arm interferometers to control the arm cavity mirrors independently of the Michelson and signal recycling cavities, known as Arm Length Stabilization (ALS). The arm cavity locking uses frequency-doubled lasers at each end station, with each laser frequency referenced to the master laser in the vertex.

It may be possible to significantly improve the robustness of the ALS system by injecting the frequency doubled laser beams from the vertex.

### 8.1.2 Alignment Sensing and Control (ASC)

**Sidles-Sigg instability** High circulating power in a kilometer-scale Fabry-Perot cavity can lead to an alignment instability in which a radiation pressure driven anti-spring overcomes the mechanical restoring force of the cavity’s suspended optics (Sidles-Sigg instability) [300]. The Advanced LIGO interferometers were designed to avoid instability, but this is still a potential threat for the Voyager/Cosmic Explorer detectors. Essentially, higher power leads to greater instability which can only be suppressed with high-bandwidth active alignment control. As the bandwidth of the ASC loops approaches the gravitational-wave band of the detector, noise from the alignment control signals is introduced into the detection band and can spoil the interferometer sensitivity. The investigations to mitigate this instability should be carried out in both practical and innovative aspects: reduction of the control noise level of the ASC system and exploration of any new scheme to directly suppress the instability (e.g., optical trapping of alignment degrees of freedom).

**Towards ASC of Voyager/Cosmic Explorer detectors** As the optical system incorporates more components, the Alignment Control System gets increasingly complicated. Already in current interferometers, we will be required to develop new techniques to control

the alignment of the OMC [301], the squeezed light source, and the filter cavities, in addition to the already complicated ASC of the main interferometer.

The higher power operation of the interferometer in future detectors will involve coupling of the higher order modes into the field content at the dark port. This coupling will depend on the angular motion of the mirrors and will make the ASC requirements more stringent. To address the issue, deep investigation of the interferometer behavior with simulation is required in the design phase.

Historically, the control noise of the ASC system has consistently been one of the main limiting technical noises in the low frequency band for all interferometric gravitational wave detectors. As the band of gravitational-wave detection moves down, and the bandwidth of alignment control increases, both for suppressing higher-order modes and optical instability, the ASC problem becomes increasingly difficult.

### 8.1.3 Calibration (CAL)

As the detectors improve sensitivity, the signal-to-noise ratio increases for known sources and the possibility for exciting new astrophysics increases. With this improvement comes the need for increased precision on individual detector calibration. [302] As the gravitational wave network adds new members, the need for accuracy between observatories becomes heightened. And finally, as previous detections have revealed, there is ever more a push to improve the calibration out to the highest possible frequencies. This section describes the R&D projects that must take place for us to move beyond the current capabilities of the calibration hardware.

**Upgrades the Photon Calibrator System** The primary reference for measuring the global sensing and actuation upon the detector’s DARM degree of freedom in Advanced LIGO has been the Photon Calibrator [303]. While this system has allowed us to achieve 2-3% magnitude, 1-2 deg phase accuracy and precision in detector response over the instrument’s most sensitive, 20-1000 Hz, band [304], its life-span is finite and maintaining the high-level of precision requires a rigorous maintenance schedule [305]. Improvements to the system, in-order to reduce implementation uncertainty such as spot-position accuracy or mitigating transmitted spot power imbalance are being considered.

**Calibration Precision/Accuracy beyond “1%”** The photon calibrator is primarily limited by (a) it’s absolute calibration of the system’s reference received from NIST, and (b) its alignment into the interferometer, to a level of 0.76%. The above paragraph discusses improvements to the system that might allow the photon calibrator to \*only\* be limited by its reference uncertainty. However, propagating that uncertainty through several combined measurements of the interferometer (to-date) adds up to the current uncertainty in detector response to DARM changes of 2-3% in magnitude. Much of the astrophysical information from SNR  $\sim 100$  events – expected to be prevalent at the design sensitivity of 2G detectors, and even more so in 3G detectors – would be limited by such  $\sim 1\%$  precision and accuracy. Thus, we must explore improvements to the current reference, and/or new ways to supplement that reference with others in order to harness the full capability of improved

sensitivity, for example by using direct gravitational interaction between a LIGO test mass and an externally modulated gravitational reference [306, 307].

**Detector Calibration Above 1 kHz** Virtually all force-actuator references are / will-be limited in actuation strength by the  $f^{-2}$  displacement response of the test mass suspensions. However, gravitational waves from supernovae and colliding neutron stars are expected to yield exciting information above  $\sim 1$  kHz. The two-fold inhibitors – high noise and low actuation strength – mean obtaining even sparse confirmational measurements at these frequencies takes days to weeks. The results are further confused by generally negligible aspects of the detector such as nonstationarity of the interferometer’s response, impacts of the full Fabry-Perot response to length changes, i.e. [308], which is yet to be updated for RSE, Fabry-Perot Dual-Recycled Michelson Interferometers (FPDRMI), near-Nyquist distortions of the digital-to-analog conversion system, and non-rigid body deformation of the test masses. Research to account for such effects in the model, conducting multi-week measurements, and considering modifications to the characterization hardware are on-going.

**Measure and Characterize the O3 Interferometer** As the detectors evolve in order to achieve design sensitivity, components that are important to the DARM control loop must change. Some changes are major, such as switching the ITMs and SRM optics between O2 and O3 which changes the interferometric response. Some changes are more minor such as replacing the electrostatic driver electronics on the ETMs (planned for the near future). Like the photon calibrator, maintaining high precision and accuracy of the detector’s response requires constant vigilance, repeated characterization measurements, and updating components to the loop model.

#### 8.1.4 Control and Data Acquisition System (CDS)

**Adaptive Noise Cancellation** Adaptive noise cancellation techniques will be widely used in the Advanced LIGO interferometers and the future detectors. Basically, noise cancellation based on Wiener filters eliminates the noise in a target signal correlated to witness channels.

This family of techniques has a wide area of application such as seismic noise reduction, Newtonian noise subtraction, and magnetic/acoustic noise cancellation. Residual coupling of auxiliary LSC degrees of freedom to the main GW signals, which presented a sensitivity limit in the initial and Enhanced LIGO interferometers, may also be reduced by adaptive noise cancellation.

Up to now, a basic scheme with linear adaptive noise cancellation is being tested on the CDS system at the 40 m prototype [309]. This technique has also been applied to seismic noise subtraction during Enhanced LIGO [309]. In the future, in addition to the linear noise subtraction, subtraction of bilinear noise should also be investigated as it may prove useful in removing a variety of noise sources from the GW signal. Machine learning (artificial neural networks, deep learning) may be able to learn non-linear and non-stationary couplings in interferometer noise for the purpose of subtracting this noise from signals [310].

Many of the noise subtraction applications to date have been performed off-line. The focus for



the controls working group will be applications that benefit from on-line, real-time, adaptive subtraction (filtering).

**Automatic Optimization** As the interferometer configuration gets more complicated, the difficulty of optimizing system parameters increases significantly. For example, we have to deal with the different range and frequency dependence of the multiple sensors and actuators at each stage of the suspension and isolation system, as opposed to the single pendulum suspensions use in iLIGO. Optimization of hierarchical control will depend on the noise level of the interferometer, which may change in time, making the traditional approach of by-hand optimization untenable.

Many of the optimization procedures in the future interferometers will need to be automated within the interferometer control system, and some may need to be dynamically adjusted. This approach should be applied in Advanced LIGO and expanded in the future detectors with continuous effort toward automatic optimization or “machine learning”. New involvement and collaboration with researchers of this particular field is encouraged.

**Modern control** Along with adaptive noise cancellation and automatic optimization, other applications of modern control theory to interferometer gravitational wave detectors should be considered.

Feedback control of interferometer gravitational wave detectors has mostly relied on the classical control theory up to now. The classical control theory, developed in the 1950s, is still effective as one of the approaches for control applications. However, control theory has continued over the intervening 60 years and significant improvements to the classical approach are available. Modern control theory, offers various possibilities to improve the interferometer control.

An example application of modern control theory is the modal-damping work done on the aLIGO suspensions at MIT [205]. The challenge faced in suspension control, and also elsewhere in aLIGO and similar detectors, is that of combining multiple sensor signals with a variety of noise levels and response functions to provide a controller which is effective in the control band while minimizing noise in the gravitational-wave detection band. The currently employed approach of classical controls requires experts to spend considerable time diagonalizing the system and constructing near-optimal (really just “good enough”) filters for each degree of freedom of the system.

These technologies can be first investigated by simulation and then applied to prototype facilities. Once effectiveness of these advanced control is demonstrated, it can be implemented to the aLIGO interferometers in order to accelerate the interferometer optimization.

**Virtual Interferometer** The flexibility of the Advanced LIGO CDS system, and recent changes to the RCG software allows, for the first time, the beginnings of the “Virtual Interferometer”. The recent RCG changes support a “virtual” IOP with the RCG code running in user space (instead of the kernel).

Future detectors should have a computer facility which enables us to login and run the control

system for the test of operating and diagnosing tests without having the actual detector equipment. At the beginning this can only be a tool for the commissioning and development, as seen in the “Simulated Plant” concept being tested at the 40m prototype [311], but eventually can provide various realtime tests of the interferometer behavior by cooperative interaction with the other simulation tools. This should include the test of the online data production for mock data challenges.

## 8.2 Research Areas

Our research efforts in the near term are focused on the following topical areas, each described in the subsections below:

- Application of Machine Learning to Controls
- Length-to-Angle (L2A) decoupling
- Feedback optimization
- System Identification
- Interferometer robust configuration for earthquakes
- State space control for Real-Time Code Generator (RCG) Software
- Angular Sensing & Control
- Calibration

### 8.2.1 Application of Machine Learning to Controls

Interferometer lock loss is a significant contributor to interferometer downtime and generally represents missed opportunities to collect science data. Lock losses occur for a variety of reasons, many of which are unknown or poorly understood. An automated system that can categorize lock losses or diagnose proximate causes can yield valuable information about where to focus commissioning efforts to improve detector stability. Considering the large amount of auxiliary data produced by the interferometers, machine learning and data mining methods may find use in the development of such an automated system. One example of an automated system to help explain lock losses is the Lockloss Monitor [312].

In the future, it would be desirable to have a system that could analyze the DARM channel along with a number of auxiliary channels to diagnose and predict lock-losses even before they occur; with enough forewarning, the interferometer controls configuration could be quickly adjusted to prevent the lock-loss from ever occurring, i.e. lock maintenance. Such a system is being designed specifically for earthquake events; see section 8.2.5.

The difficulty in automatically locking an optical cavity (especially a high finesse cavity) is that there is no signal to use for feedback until very near the resonance condition. When the mirrors are swinging uncontrolled, the system is effectively ‘blind’. This is motivation for the (lower finesse) Arm Length Stabilization (ALS) system, which uses light which is frequency

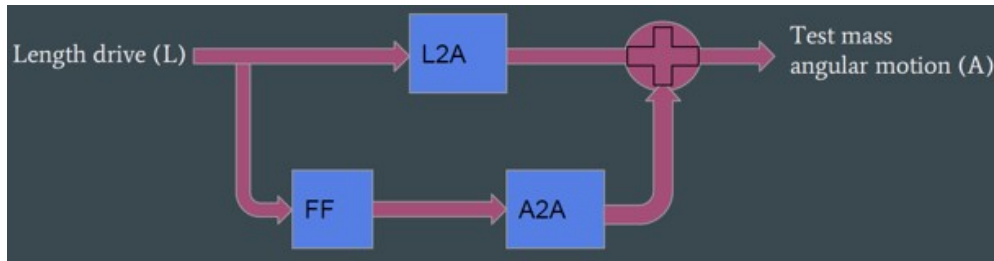


Figure 27: Local feedforward correction filter for length-to-angle decoupling.  
 $FF = -L2A/A2A$ .

doubled main laser light (i.e. green light), as a first step in getting the Fabry-Perot arm cavities close to the main laser light resonance condition. Machine learning techniques may be able to perform state estimation on these nonlinear plants (resonant cavities) and serve to provide guidance to a control system [310]. An experiment to train a Deep Learning (DL) system on simulated interferometer data and then apply it to the 40m interferometer is underway.

### 8.2.2 Length-to-Angle (L2A) decoupling

The suspension systems exhibit frequency dependent coupling between length commands and angular motion due to cross-coupling as well as imbalance between the coil actuators. Although the angular control loops mitigate this coupling, improved control can result from implementing length to angle (L2A) feedforward correction. Local feedforward correction at the suspension level can be used to correct for mechanical cross-coupling in the suspensions, as depicted in figure 27. In addition to local feedforward correction, the length control signal can be fed forward to the angular control system to compensate for any residual cross-coupling (e.g. associated with sensing).

The local L2A feedforward correction may permit improved angular sensing and control (ASC) stability and performance by reshaping the ASC loop design. The Power Recycling Mirror (PRM) length-to-pitch decoupling example in figure 28 resulted in improved robustness to disturbances. Global feed forward correction has yet to be attempted, but coherence studies indicate potentially significant performance gains. Various time domain and frequency domain approaches to feed forward filter synthesis are being explored.

If the beam is decentered on the test mass, angular noise couples to strain. If we're able to sense beam position relative to the center of the test mass (optic), then we can steer the beam to center. Using the fixed pattern of scatter points embedded in the optic coating, convolved/modulated with the jittering Gaussian-shaped beam, gives us the potential to determine the beam centroid, either with standard image processing techniques, or via machine learning [310]. An experiment to try this DL approach on track the beam position on an image stream of a test mass at LHO is planned.

### 8.2.3 Feedback optimization

Feedback optimization serves two goals:

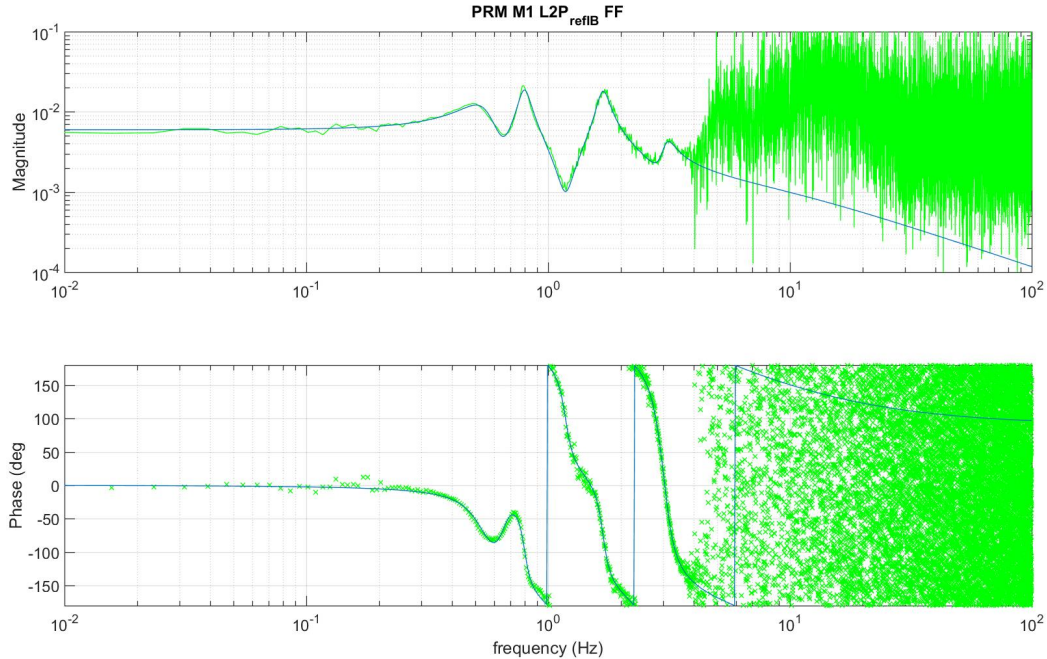


Figure 28: Local feedforward correction filter for length-to-pitch decoupling for the Power Recycling Mirror (PRM), M1 stage. Trend data indicates improved robustness to disturbances (wind, earthquakes) after implementing length-to-pitch (L2P) and length-to-yaw (L2Y) feedforward (FF) filters from the ReflB photodiode for the PRM.

1. to optimize our control loops directly in terms of the requirements of the interferometer, and
2. to automate the design process.

The first goal involves writing a cost function directly in terms of what matters to the interferometer, such as noise, control authority, and stability. The second goal requires utilization of a general purpose optimization algorithm, such as a particle swarm method that is capable of finding global minimums amongst numerous local minima. This algorithm should reliably find a close to optimal solution in a reasonably short amount of time, alleviating the need for commissioners to create time consuming designs using loop shaping and guess and check techniques that result in suboptimal solutions. Traditional optimal control approaches, such as LQG and H-infinity, have not proven more useful than loop shaping and guess and check techniques for many applications due to the fact that the cost functions do not directly translate to interferometer requirements. Consequently, we are working on more generalized cost functions and optimization algorithms to meet the goals of this feedback optimization focus area.

### 8.2.4 System Identification and Transfer Function Fitting

System identification is the process of developing or improving a mathematical representation of a physical system using experimental data. Techniques for system identification range from those which require physical insight to develop an underlying model structure to those which are model-independent fits of the experimental data.

All control implementations require some form of system identification, either explicitly or implicitly. As LIGO seeks to implement more robust control, better descriptions of the plant, and the associated plant model uncertainty, are required.

The basic system identification process steps, which incorporate physical insight, are as follows:

- Build a physics-based model
- Design the transfer function measurement
- Collect the transfer function data
- Select, filter, de-trend the data, as appropriate
- Fit the model
- Validate the model (compare prediction to data)
- Iterate until converged on a validated model

Optimal transfer function measurement [313, 314] is basically concentrating the injection to the frequencies where the uncertainty in the model is greatest. Given a coarse model of the system, this method delivers lower system parameter uncertainties for a fixed excitation RMS and measurement time. In many cases, we have  $< 50\%$  uncertainty in the roots, so this is less expensive in time than our usual technique of making a detailed swept-sine measurement with `diaggui`. As usual, we use `awgstream` or `awg.py` to inject this signal. Our next steps are to develop user friendly software for control room use - this software can also be prototyped using the Virtual Interferometer; see section 8.1.4.

Transfer function fitting in LIGO is often divorced from the physics-based model, i.e. we do not fit to model parameters, we instead curve fit the frequency domain transfer function response and do not validate the model and use it for subsequent control synthesis. Even in these scenarios, we can benefit from more accurate curve fits to the measured transfer functions, as indicated in figure 29.

System identification is a well established field rich with techniques suitable for various applications. We are working toward:

- determining the applicability of various system identification techniques to GW interferometer subsystems,
- defining a "road map" for application to our systems and further research,

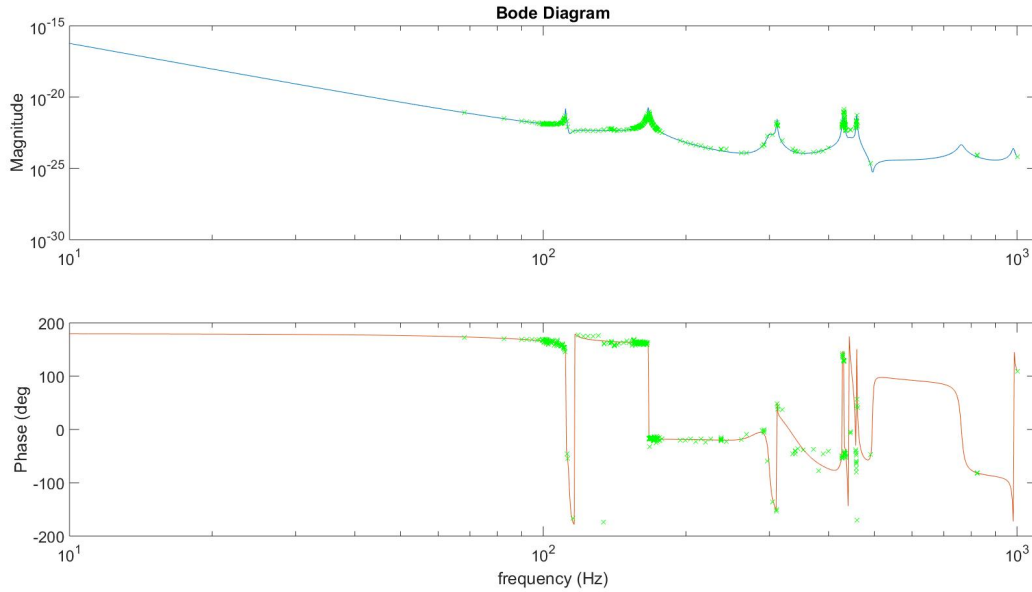


Figure 29: High Frequency Suspension Dynamics: A number of dynamical, resonant features have been found in the quadruple pendulum suspension (UIM/L1 actuation stage) in the frequency range of 100 to 500 Hz. Measurements are a combination of swept-sine (SS) and broad-band (BB) transfer function ratios. For system calibration, a pole-zero fit of this transfer function above 20 Hz, that can be used as a “correction” filter to a model that falls off as  $f^6$ , was needed. While this filter does not need to perfectly resolve the details of all of the high-Q features, but it must track the overall frequency dependence over the 20 - 500 Hz region well. The best pole:zero fit is the stable, non-minimum phase, zero delay, 22:28 order model shown (compatible with a Linear Time Invariant (LTI) representation).

- using frequency, time domain and state space techniques for model parameter identification, and
- developing methods to incorporate measurement uncertainty and underlying physical modeling to provide better fits to measurement data from which improved and robust control laws can be formulated.

Ultimately we seek a recursive, real-time parameter identification of our MIMO systems (in their operating states). Many techniques are available such as Generalized Least Squares and Maximum Likelihood Estimators, which are computationally simple. Virgo has been exploring the development of Observer/Kalman Filter identification techniques (time domain based) extended to identification of a closed loop effective controller/observer combination for their seismic isolation system. We think that State-Space Frequency Domain (SSFD) identification techniques are a natural extension to our traditional frequency domain approaches and lead naturally to the application of modern control approaches to our MIMO systems.

Improvements of the existing models may be required in order to account for cross coupling or complex geometry effects, such physical features not being captured by the current models.

We see finite element methods as good approach for this task, namely to model the full geometric complexity of the plant, as well as to cover its high frequency dynamical response. In addition, a reduced-order parametric model of the system, extracted from the the finite element model, would be tested against experimental characterizations, through comparison between modeled and measured transfer functions for instance, and the parameters would be adjusted consequently. Such a parameterization would help to identify and study the impact of key physical aspects of the suspension behavior. This modeling strategy, involving detailed physical modeling of the plant, a parametric model, the fitting and the verification against experimental data would not only be used to enhance the reliability of the model in the framework of robust control of the suspension, but will also be of great use in in the context of future detectors design.

### 8.2.5 Interferometer robust configuration for earthquakes

The detectors are susceptible to teleseismic events, which can significantly impact a detector’s duty cycle. Efforts by the SWG have produced an early-warning system, known as *Seismon* [315, 316], for gravitational-wave observatories, which relies on near real-time earthquake alerts provided by the U.S. Geological Survey (USGS) and the National Oceanic and Atmospheric Administration (NOAA). *Seismon* estimates the amplitude of the ground velocity at the site, and the time-of-arrival, to create warnings for the detectors.

Losing optical cavity lock leads to significant down time due to the time to reacquire lock and reestablish thermal equilibrium. Algorithms (including machine learning [312]) are being explored to set thresholds on predicted ground motion relative to probability of lock loss.

The goal of the early warning system is to allow time to reconfigure the control system if the probability of interferometer lock loss is high. The control system reconfiguration trades off sensitivity for added robustness in order to maintain lock. For this reason, it cannot be permanently activated and needs to be part of the Guardian switching system. Based on the information provided by *Seismon*, we could regain duty cycle otherwise lost during earthquake events. Advanced LIGO has an impressive array of sensors and a flexible control system, and different control strategies are being considered (see section 6.1.4). Studies are conducted to investigate optimal ways to reduce the impact of the extra-motion created by earthquakes at low frequencies while still providing isolation performance in the LIGO detection band. For example, tuning the seismic platform, control scheme to reduce the gain peaking at low frequencies has been studied [317]. Other control strategies are underway to increase the interferometer robustness via mechanical length to angle decoupling [318] and common mode rejection along the arms [319].

### 8.2.6 State space control for Real-Time Code Generator (RCG) Software

The current primary mechanism in the RCG environment for implementing MIMO control is to use filter matrices. While these filter matrices work well for some applications, they are often suboptimal for implementing state space models for applications such as simulated plants, fault detection or ‘Modern Controls’ techniques such as Kalman filtering and LQG. Since these applications begin in state space form, it is natural to implement them in state space form. Converting them to filter matrices introduces computational redundancies such

as numerous pole-zeros cancellations and/or unnecessary transformation steps in the implementation. The implementation of state space blocks in the RCG environment is envisioned to have a user interface similar to existing matrix blocks, with the added complexities that state spaces will consist of four matrices, the A, B, C, and D matrices, and the user will have to specify the number of inputs, outputs, and states. Since matrices are easy to update in real-time in the RCG environment, these state space tools may also open up new ways for implementing adaptive control techniques.

### 8.2.7 Angular Sensing & Control

As mentioned in section 8.1.2, the control noise of the ASC system has consistently been one of the main limiting technical noises in the low frequency band for all interferometric gravitational wave detectors. In fact angular control noise is dominant in the aLIGO spectrum below 25 Hz, due to a number of causes [320]. We plan to focus on the following research related to angular sensing & control:

- make the ‘plant’ appear to be power independent by digitally compensating for the power dependency [321]
- investigate means to reduce excess noise, and noise coupling, to the ASC sensors
- investigate new sensing schemes for the SRC (e.g. [322]) to overcome current sensing matrix limitations
- study means to mitigate initial alignment drift so that lock re-acquisition time is minimized

### 8.2.8 Calibration

As mentioned and motivated in section 8.1.3, we plan to focus on the following research:

- Improvements to the photon calibration system, in-order to reduce implementation uncertainty, such as spot-position accuracy and mitigating transmitted spot power imbalance.
- Improvements to the calibration reference, and/or new ways to supplement that reference.
- Research to account for calibration effects/limitations in the model at high frequencies (< 1 kHz).



## References

- [1] L. Barsotti, P. Fritschel, M. Evans, and S. Gras, “Updated Advanced LIGO Design Curve,” *LIGO Document: LIGO-T1800044*, 2018. <https://dcc.ligo.org/LIGO-T1800044>.
- [2] P. F. M. E. Lisa Barsotti, Lee McCullen, “The A+ design curve,” 2018. LIGO-T1800042, <https://dcc.ligo.org/LIGO-T1800042>.
- [3] P. Purdue and Y. Chen, “Practical speed meter designs for quantum nondemolition gravitational-wave interferometers,” *Phys. Rev. D*, **66**, 122004, 2002.
- [4] Y. Chen, “Sagnac interferometer as a speed-meter-type, quantum-nondemolition gravitational-wave detector,” *Phys. Rev. D*, **67**, 122004, 2003.
- [5] S. L. Danilishin, “Sensitivity limitations in optical speed meter topology of gravitational-wave antennas,” *Phys. Rev. D*, **69**, 102003, 2004.
- [6] H. Müller-Ebhardt, *On quantum effects in the dynamics of macroscopic test masses*. PhD thesis, Leibniz Universität Hannover, 2009.
- [7] M. Wang, C. Bond, D. Brown, F. Brückner, L. Carbone, R. Palmer, and A. Freise, “Realistic polarizing Sagnac topology with DC readout for the Einstein Telescope,” *Phys. Rev. D*, **87**, 096008, May 2013. <http://link.aps.org/doi/10.1103/PhysRevD.87.096008>.
- [8] A. Wade, K. Mckenzie, Y. Chen, D. Shaddock, J. Chow, and D. McClelland, “A Polarization Speed Meter for Gravitational-Wave Detection,” *LIGO Document: P1100205*, 2011. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=77471>.
- [9] S. P. Voronchev, N. V.; Tarabrin and S. L. Danilishin, “Broadband detuned Sagnac interferometer for future generation gravitational wave astronomy,” *arXiv:1503.01062*, 2015.
- [10] S. H. Huttner, S. L. Danilishin, B. W. Barr, A. S. Bell, C. Gr??f, J. S. Hennig, S. Hild, E. A. Houston, S. S. Leavey, D. Pascucci, B. Sorazu, A. P. Spencer, S. Steinlechner, J. L. Wright, T. Zhang, and K. A. Strain, “Candidates for a possible third-generation gravitational wave detector: comparison of ring-Sagnac and sloshing-Sagnac speed-meter interferometers,” *Classical and Quantum Gravity*, **34**, no. 2, 024001, 2017. <http://stacks.iop.org/0264-9381/34/i=2/a=024001>.
- [11] E. Knyazev, S. Danilishin, S. Hild, and F. Khalili, “Speedmeter scheme for gravitational-wave detectors based on EPR quantum entanglement,” *Physics Letters A*, 2017. <http://www.sciencedirect.com/science/article/pii/S0375960117302888>.
- [12] S. L. Danilishin, E. Knyazev, N. V. Voronchev, F. Y. Khalili, C. Gräf, S. Steinlechner, J.-S. Hennig, and S. Hild, “A new quantum speed-meter interferometer: Measuring speed to search for intermediate mass black holes,” *Light: Science & Applications*, 2018. <http://aap.nature-lsa.cn:8080/cms/accessory/files/AAP-lsa201815.pdf>.

- [13] S. Dwyer. personal communication, 2012.
- [14] T. Corbitt, Y. Chen, and N. Mavalvala, “Mathematical framework for simulation of quantum fields in complex interferometers using the two-photon formalism,” *Physical Review A*, **72**, pp. 013818+, July 2005. <http://dx.doi.org/10.1103/physreva.72.013818>.
- [15] R. X. Adhikari, “Gravitational Radiation Detection with Laser Interferometry,” *Reviews of Modern Physics*, 2013.
- [16] LIGO Scientific Collaboration and Virgo Collaboration, “Observation of Gravitational Waves from a Binary Black Hole Merger,” *Physical Review Letters*, **116**, 061102, 2016.
- [17] LIGO Scientific Collaboration and Virgo Collaboration, “Binary black hole mergers in the first advanced LIGO observing run,” *Physical Review X*, **6**, no. 4, 041015, 2016.
- [18] LIGO Scientific Collaboration and Virgo Collaboration, “GW170104: observation of a 50-solar-mass binary black hole coalescence at redshift 0.2,” *Physical Review Letters*, **118**, no. 22, 221101, 2017.
- [19] LIGO Scientific Collaboration and Virgo Collaboration, “GW170608: Observation of a 19 solar-mass binary black hole coalescence,” *The Astrophysical Journal Letters*, **851**, no. 2, L35, 2017.
- [20] LIGO Scientific Collaboration and Virgo Collaboration, “GW170814: A three-detector observation of gravitational waves from a binary black hole coalescence,” *Physical review letters*, **119**, no. 14, 141101, 2017.
- [21] B. Abbott, R. Abbott, R. Adhikari, A. Ananyeva, S. Anderson, S. Appert, K. Arai, M. Araya, J. Barayoga, B. Barish, *et al.*, “Multi-messenger observations of a binary neutron star merger,” *Astrophysical Journal Letters*, **848**, no. 2, L12, 2017.
- [22] J. McIver and A. Lundgren, “2018-2019 LIGO DetChar white paper draft,” *LIGO Document: LIGO-T1800167*, 2018. <https://dcc.ligo.org/LIGO-T1800167>.
- [23] P. Shawhan and G. A. Prodi, “The LSC-Virgo White Paper on Gravitational Wave Data Analysis and Astrophysics (2018-2019 edition),” *LIGO Document: LIGO-T1800058*, 2018. <https://dcc.ligo.org/LIGO-T1800058>.
- [24] Einstein Telescope Science Team, “Einstein gravitational wave Telescope conceptual design study,” *ET Document: ET-0106C-10*, 2011. <http://www.et-gw.eu/>.
- [25] C. D. Ott, Y. Chen, and R. X. Adhikari, “Astrophysical Motivations for the Third Generation LIGO Detectors,” *LIGO Document: LIGO-P1200099*, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=87934>.
- [26] KAGRA Collaboration, LIGO Scientific Collaboration and Virgo Collaboration, “Prospects for observing and localizing gravitational-wave transients with Advanced LIGO, Advanced Virgo and KAGRA,” *LIGO Document: LIGO-P1200087*, 2018. <https://dcc.ligo.org/LIGO-P1200087>.

- [27] L. Barsotti, L. McCuller, M. Evans, and P. Fritschel, “The A+ design curve,” *LIGO Document: LIGO-T1800042*, 2018. <https://dcc.ligo.org/LIGO-T1800042>.
- [28] J. Miller, L. Barsotti, P. Fritschel, and M. Evans, “Prospects for doubling the range of Advanced LIGO,” *LIGO Document: LIGO-T0900144*, 2014. <https://dcc.ligo.org/LIGO-P1400164>.
- [29] S. Chua, B. Slagmolen, D. Shaddock, and D. McClelland, “Quantum squeezed light in gravitational-wave detectors,” *Classical and Quantum Gravity*, **29**, 145015, 2012.
- [30] J. Abadie and et al., “Enhanced sensitivity of the LIGO gravitational wave detector by using squeezed states of light,” *Nature Photonics*, **7**, pp. 613–619, 2013.
- [31] H. Grote, K. Danzmann, K. Dooley, R. Schnabel, J. Slutsky, and H. Vahlbruch, “First Long-Term Application of Squeezed States of Light in a Gravitational-Wave Observatory,” *Phys. Rev. Lett.*, **110**, 2013.
- [32] H. J. Kimble, Y. Levin, A. B. Matsko, K. S. Thorne, and S. P. Vyatchanin, “Conversion of conventional gravitational-wave interferometers into quantum nondemolition interferometers by modifying their input and/or output optics,” *Phys. Rev. D*, **65**, 022002, 2001.
- [33] J. Harms and et al., “Squeezed-input, optical-spring, signal-recycled gravitational-wave detectors,” *Phys. Rev. D*, **68**, 042001, 2003.
- [34] G. M. Harry, A. M. Gretarsson, P. R. Saulson, S. E. Kittelberger, S. D. Penn, W. J. Startin, S. Rowan, M. M. Fejer, D. R. M. Crooks, G. Cagnoli, J. Hough, and N. Nakagawa, “Thermal noise in interferometric gravitational wave detectors due to dielectric optical coatings,” *Class. Quantum Grav.*, **19**, no. 5, 897, 2002.
- [35] editors G. Harry, T. Body, and R. de Salvo, “Optical Coatings and Thermal Noise in Precision Measurement ,” *CUP*, pp. ISBN-13: 978-1107003385, 2012.
- [36] G. M. Harry, M. R. Abernathy, A. E. Becerra-Toledo, H. Armandula, E. Black, K. Dooley, M. Eichenfield, C. Nwabugwu, A. Villar, D. R. M. Crooks, G. Cagnoli, J. Hough, C. R. How, I. MacLaren, P. Murray, S. Reid, S. Rowan, P. H. Sneddon, M. M. Fejer, R. Route, S. D. Penn, P. Ganau, J.-M. Mackowski, C. Michel, L. Pinard, and A. Remillieux, “Titania-doped tantala/silica coatings for gravitational-wave detection,” *Class. Quantum Grav.*, **24**, no. 2, 405, 2007.
- [37] M. Evans, S. Ballmer, M. Fejer, P. Fritschel, G. Harry, and G. Ogin, “Thermo-optic noise in coated mirrors for high-precision optical measurements,” *Phys. Rev. D*, **78**, 102003, 2008.
- [38] G. D. Cole, W. Zhang, M. J. Martin, J. Ye, and M. Aspelmeyer, “Tenfold reduction of Brownian noise in optical interferometry,” *ArXiv e-prints*, Feb. 2013.
- [39] A. Cumming and . others, “Measurement of the mechanical loss of prototype GaP/AlGaP crystalline coatings for future gravitational wave detectors,” *Classical and Quantum Gravity*, 2014. <https://dcc.ligo.org/LIGO-P1400059>.

- [40] A. Lin, R. B. S. Omar, *et al.*, “Epitaxial growth of GaP/AlGaP mirrors on Si for low thermal noise optical coatings,” *LIGO Document: LIGO-P1400236*, 2014. <https://dcc.ligo.org/LIGO-P1400236>.
- [41] N. Smith, A. Brooks, E. Gustafson, D. Coyne, R. Adhikari, L. Barsotti, and B. Shapiro, “LIGO Voyager Upgrade Concept,” *LIGO Document: LIGO-T1400226*, 2017. <https://dcc.ligo.org/LIGO-T1400226>.
- [42] M. Punturo, M. Abernathy, F. Acernese, B. Allen, N. Andersson, K. Arun, F. Barone, B. Barr, M. Barsuglia, M. Beker, *et al.*, “The Einstein Telescope: a third-generation gravitational wave observatory,” *Classical and Quantum Gravity*, **27**, no. 19, 194002, 2010.
- [43] B. P. Abbott, R. Abbott, T. Abbott, M. Abernathy, K. Ackley, C. Adams, P. Addesso, R. Adhikari, V. Adya, C. Affeldt, *et al.*, “Exploring the sensitivity of next generation gravitational wave detectors,” *Classical and Quantum Gravity*, **34**, no. 4, 044001, 2017.
- [44] N. Smith-Lefebvre and N. Mavalvala, “Modematching feedback control for interferometers with an output mode cleaner,” *LIGO Document: LIGO-P1200034*, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=89076>.
- [45] P. Willems, “Implications of ETM02 HR Coating Absorption for Thermal Compensation,” *LIGO Document, LIGO-T1100250-v2*, 2011. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=60664>.
- [46] A. F. Brooks and *et al.*, “Ultra-sensitive wavefront measurement using a Hartmann sensor,” *Opt. Express*, **15**, pp. 10370–10375, 2007. <http://www.opticsinfobase.org/oe/abstract.cfm?uri=oe-15-16-10370>.
- [47] A. F. Brooks, *Hartmann Wavefront Sensors for Advanced Gravitational Wave Interferometers*. PhD thesis, The University of Adelaide, 2007. <http://digital.library.adelaide.edu.au/dspace/handle/2440/57100>.
- [48] H. Yamamoto, “SIS: Stationary interferometer simulation.” <https://labcit.ligo.caltech.edu/~hiro/SIS/>. Accessed: 2018-05-18.
- [49] J. Degallaix, “OSCAR a Matlab based optical FFT code,” *Journal of Physics: Conference Series*, **228**, no. 1, 012021. Code is available at <https://www.mathworks.com/matlabcentral/fileexchange/20607-oscar> year=2010,, <http://stacks.iop.org/1742-6596/228/i=1/a=012021>.
- [50] H. Yamamoto, B. Bhawal, M. Evans, E. Maros, M. Rakhmanov, J. R. L. Savage, G. Cella, S. Klimenko, and S. Mohanty, “End to End Simulation Program for Gravitational-Wave Detectors,” , in *Gravitational Wave Detection II* (S. Kawamura and N. Mio, eds.), vol. 32 of *Frontiers Science Series*, pp. 331–336, Universal Academy Press, 2000.

- [51] H. Yamamoto, B. Bhawal, M. Evans, M. E. M. Rakhmanov, and R. L. J. Savage, “End to End Simulation Program for Gravitational-Wave Detectors,” , in *Gravitational Wave Detection II* (S. Kawamura and N. Mio, eds.), no. 32 in Frontier Science Series, Universal Academy Press, 2000.
- [52] M. Evans, “Optickle simulation software.” <https://github.com/Optickle/Optickle>, 2004. Accessed: 2018-05-18, <https://github.com/Optickle/Optickle>.
- [53] D. D. Brown and A. Freise, “Finesse.” <http://www.gwoptics.org/finesse>, May 2014. You can download the binaries and source code at <http://www.gwoptics.org/finesse>., <http://www.gwoptics.org/finesse>.
- [54] A. Freise, G. Heinzl, H. Lück, R. Schilling, B. Willke, and K. Danzmann, “Frequency-domain interferometer simulation with higher-order spatial modes,” *Classical and Quantum Gravity*, **21**, no. 5, S1067, 2004. <http://www.gwoptics.org/finesse/>.
- [55] G. Vajente, “Fast modal simulation of paraxial optical systems: the MIST open source toolbox,” *Classical and Quantum Gravity*, **30**, no. 7, 075014, 2013. <http://stacks.iop.org/0264-9381/30/i=7/a=075014>.
- [56] LIGO, “GWINC: Gravitational Wave Interferometer Noise Calculator.” <http://ilog.ligo-wa.caltech.edu:7285/advligo/GWINC>.
- [57] J. Harms, “Terrestrial Gravity Fluctuations,” *Living Reviews in Relativity*, **18**, no. 3, 2015. <http://www.livingreviews.org/lrr-2015-3>.
- [58] J. C. Driggers, J. Harms, and R. X. Adhikari, “Subtraction of Newtonian noise using optimized sensor arrays,” *Phys. Rev. D*, **86**, 102001, Nov 2012. <http://link.aps.org/doi/10.1103/PhysRevD.86.102001>.
- [59] M. Coughlin, N. Mukund, J. Harms, J. Driggers, R. Adhikari, and S. Mitra, “Towards a first design of a Newtonian-noise cancellation system for Advanced LIGO,” *Classical and Quantum Gravity*, **33**, no. 24, 244001, 2016. <http://stacks.iop.org/0264-9381/33/i=24/a=244001>.
- [60] M. Coughlin, J. Harms, N. Christensen, V. Dergachev, R. DeSalvo, S. Kandhasamy, and V. Mandic, “Wiener filtering with a seismic underground array at the Sanford Underground Research Facility,” *Classical and Quantum Gravity*, **31**, no. 21, 215003, 2014. <http://stacks.iop.org/0264-9381/31/i=21/a=215003>.
- [61] D. Fiorucci, J. Harms, M. Barsuglia, I. Fiori, and F. Paoletti, “Impact of infrasound atmospheric noise on gravity detectors used for astrophysical and geophysical applications,” *Phys. Rev. D*, **97**, 062003, Mar 2018. <https://link.aps.org/doi/10.1103/PhysRevD.97.062003>.
- [62] J. Harms and K. Venkateswara, “Newtonian-noise cancellation in large-scale interferometric GW detectors using seismic tiltmeters,” *Classical and Quantum Gravity*, **33**, no. 23, 234001, 2016. <http://stacks.iop.org/0264-9381/33/i=23/a=234001>.

- [63] G. Cella, “Off-Line Subtraction of Seismic Newtonian Noise,” in *Recent Developments in General Relativity* (B. Casciaro, D. Fortunato, M. Francaviglia, and A. Masiello, eds.), pp. 495–503, Springer Milan, 2000.
- [64] T. Corbitt and N. Mavalvala, “Review: Quantum noise in gravitational-wave interferometers,” *Journal of Optics B: Quantum and Semiclassical Optics*, **6**, no. 8, S675, 2004. <http://dx.doi.org/10.1088/1464-4266/6/8/008>.
- [65] D. McClelland, N. Mavalvala, Y. Chen, and R. Schnabel, “Advanced interferometry, quantum optics and optomechanics in gravitational wave detectors,” *Laser & Photonics Reviews*, **5**, 677, 2011.
- [66] S. L. Danilishin and F. Y. Khalili, “Quantum Measurement Theory in Gravitational-Wave Detectors,” *Living Reviews in Relativity*, **15**, 5, 2012. <http://www.livingreviews.org/lrr-2012-5>.
- [67] C. M. Caves, “Quantum-mechanical noise in an interferometer,” *Phys. Rev. D*, **23**, 1693, 1981.
- [68] M. T. Jaekel and S. Reynaud, “Quantum Limits in Interferometric Measurements,” *Europhys. Lett.*, **13**, 301, 1990.
- [69] R. Schnabel, N. Mavalvala, D. D. E. McClelland, and P. K. P. Lam, “Quantum metrology for gravitational wave astronomy,” *Nature communications*, **1**, 121, jan 2010. <http://www.ncbi.nlm.nih.gov/pubmed/21081919>.
- [70] R. Schnabel, “Squeezed states of light and their applications in laser interferometers,” *Physics Reports*, **684**, pp. 1–51, apr 2017. <http://linkinghub.elsevier.com/retrieve/pii/S0370157317300595>.
- [71] Y. Ma, H. Miao, B. H. Pang, M. Evans, C. Zhao, J. Harms, R. Schnabel, and Y. Chen, “Proposal for gravitational-wave detection beyond the standard quantum limit through EPR entanglement,” *Nature Physics*, may 2017. <http://www.nature.com/doi/10.1038/nphys4118>.
- [72] A. F. Pace, M. J. Collett, and D. F. Walls, “Quantum limits in interferometric detection of gravitational radiation,” *Phys. Rev. A*, **47**, 3173, 1993.
- [73] A. Wicht, K. Danzmann, M. Fleischhauer, M. Scully, G. Müller, R.-H. R. Rinkleff, and G. Müller, “White-light cavities, atomic phase coherence, and gravitational wave detectors,” *Opt. Commun.*, **134**, no. 1-6, 431, 1997.
- [74] H. J. Kimble, Y. Levin, A. B. Matsko, K. S. Thorne, and S. P. Vyatchanin, “Conversion of conventional gravitational-wave interferometers into quantum nondemolition interferometers by modifying their input and/or output optics,” *Phys. Rev. D*, **65**, 022002, 2001. <http://link.aps.org/doi/10.1103/PhysRevD.65.022002>.
- [75] H. Rehbein, J. Harms, R. Schnabel, and K. Danzmann, “Optical transfer functions of Kerr nonlinear cavities and interferometers,” *Phys. Rev. Lett.*, **95**, 193001, 2005.

- [76] M. Korobko, L. Kleybolte, S. Ast, H. Miao, Y. Chen, and R. Schnabel, “Beating the Standard Sensitivity-Bandwidth Limit of Cavity-Enhanced Interferometers with Internal Squeezed-Light Generation,” *Physical Review Letters*, **118**, 143601, apr 2017. <http://link.aps.org/doi/10.1103/PhysRevLett.118.143601>.
- [77] V. B. Braginsky, M. L. Gorodetsky, and F. Y. Khalili, “Optical bars in gravitational wave antennas,” *Phys. Lett. A*, **232**, 340, 1997.
- [78] A. Buonanno and Y. Chen, “Signal recycled laser-interferometer gravitational-wave detectors as optical springs,” *Phys. Rev. D*, **65**, 042001, 2002.
- [79] F. Y. Khalili and Y. Levin, “Speed meter as a quantum nondemolition measuring device for force,” *Phys. Rev. D*, **54**, pp. 4735–4737, 1996. <http://link.aps.org/doi/10.1103/PhysRevD.54.4735>.
- [80] P. Purdue, “Analysis of a quantum nondemolition speed-meter interferometer,” *Phys. Rev. D*, **66**, 022001, 2002. <http://link.aps.org/doi/10.1103/PhysRevD.66.022001>.
- [81] P. Purdue and Y. Chen, “Practical speed meter designs for quantum nondemolition gravitational-wave interferometers,” *Phys. Rev. D*, **66**, 122004, 2002. <http://link.aps.org/doi/10.1103/PhysRevD.66.122004>.
- [82] Y. Chen, “Sagnac interferometer as a speed-meter-type, quantum-nondemolition gravitational-wave detector,” *Phys. Rev. D*, **67**, 122004, 2003. <http://link.aps.org/doi/10.1103/PhysRevD.67.122004>.
- [83] F. Khalili and E. Polzik, “Overcoming the SQL in gravitational wave detectors using spin systems with negative effective mass,” *arXiv:1710.10405*, 2017.
- [84] The LIGO Scientific Collaboration, J. Abadie, B. P. Abbott, R. Abbott, T. D. Abbott, M. Abernathy, C. Adams, R. Adhikari, C. Affeldt, B. Allen, and et al., “A gravitational wave observatory operating beyond the quantum shot-noise limit,” *Nature Physics*, **7**, pp. 962–965, 2011. <http://dx.doi.org/10.1038/nphys2083>.
- [85] H. Vahlbruch, M. Mehmet, K. Danzmann, and R. Schnabel, “Detection of 15 dB Squeezed States of Light and their Application for the Absolute Calibration of Photoelectric Quantum Efficiency,” *Phys. Rev. Lett.*, **117**, 110801, Sep. 2016. <http://link.aps.org/doi/10.1103/PhysRevLett.117.110801>.
- [86] A. R. Wade, G. L. Mansell, S. S. Y. Chua, R. L. Ward, B. J. J. Slagmolen, D. A. Shaddock, and D. E. McClelland, “A squeezed light source operated under high vacuum,” *Scientific Reports*, **5**, 18052, Dec. 2015. <http://www.nature.com/articles/srep18052> \n <http://www.nature.com/articles/srep18052.pdf> <http://www.nature.com/articles/srep18052>.
- [87] T. Eberle, S. Steinlechner, J. Bauchrowitz, V. Händchen, H. Vahlbruch, M. Mehmet, H. Müller-Ebhardt, and R. Schnabel, “Quantum Enhancement of the Zero-Area Sagnac Interferometer Topology for Gravitational Wave Detection,” *Phys. Rev. Lett.*, **104**, 251102, June 2010. <http://link.aps.org/doi/10.1103/PhysRevLett.104.251102>.

- [88] M. Mehmet, S. Ast, T. Eberle, S. Steinlechner, H. Vahlbruch, and R. Schnabel, “Squeezed light at 1550 nm with a quantum noise reduction of 12.3 dB,” *Optics Express*, **19**, pp. 25763–72, Dec. 2011. <http://www.ncbi.nlm.nih.gov/pubmed/22273968>.
- [89] C. Affeldt, K. Danzmann, K. L. Dooley, H. Grote, M. Hewitson, S. Hild, J. Hough, J. Leong, H. Lück, M. Prijatelj, S. Rowan, A. Rüdiger, R. Schilling, R. Schnabel, E. Schreiber, B. Sorazu, K. A. Strain, H. Vahlbruch, B. Willke, W. Winkler, and H. Wittel, “Advanced techniques in GEO 600,” *Classical and Quantum Gravity*, **31**, no. 22, 224002, 2014. <http://stacks.iop.org/0264-9381/31/i=22/a=224002>.
- [90] K. L. Dooley, J. R. Leong, T. Adams, C. Affeldt, A. Bisht, C. Bogan, J. Degallaix, C. Gr??f, S. Hild, J. Hough, A. Khalaidovski, N. Lastzka, J. Lough, H. Lück, D. Macleod, L. Nuttall, M. Prijatelj, R. Schnabel, E. Schreiber, J. Slutsky, B. Sorazu, K. A. Strain, H. Vahlbruch, M. W??s, B. Willke, H. Wittel, K. Danzmann, and H. Grote, “GEO 600 and the GEO-HF upgrade program: successes and challenges,” *Classical and Quantum Gravity*, **33**, no. 7, 075009, 2016. <http://stacks.iop.org/0264-9381/33/i=7/a=075009>.
- [91] H. Billing, K. Maischberger, A. Rudiger, R. Schilling, L. Schnupp, and W. Winkler, “An argon laser interferometer for the detection of gravitational radiation,” *Journal of Physics E: Scientific Instruments*, **12**, no. 11, 1043, 1979. <http://stacks.iop.org/0022-3735/12/i=11/a=010>.
- [92] H. Vahlbruch, S. Chelkowski, K. Danzmann, and R. Schnabel, “Quantum engineering of squeezed states for quantum communication and metrology,” *New Journal of Physics*, **9**, pp. 371–371, Oct. 2007. <http://stacks.iop.org/1367-2630/9/i=10/a=371?key=crossref.8e9a9d0d78283cd8be6e00ef4a83539d>.
- [93] M. Meinders and R. Schnabel, “Sensitivity improvement of a laser interferometer limited by inelastic back-scattering, employing dual readout,” *Classical and Quantum Gravity*, **32**, 195004, oct 2015. <http://www.scopus.com/inward/record.url?eid=2-s2.0-84941756735&partnerID=tZOtx3y1> <http://stacks.iop.org/0264-9381/32/i=19/a=195004?key=crossref.d3a9bdf0d09ee0e2baeb0f48db59666b>.
- [94] S. Steinlechner, J. Bauchrowitz, M. Meinders, H. Müller-Ebhardt, K. Danzmann, and R. Schnabel, “Quantum-dense metrology,” *Nature Photonics*, **7**, pp. 626–630, June 2013. <http://www.nature.com/doifinder/10.1038/nphoton.2013.150>.
- [95] M. Ast, S. Steinlechner, and R. Schnabel, “Reduction of Classical Measurement Noise via Quantum-Dense Metrology,” *Phys. Rev. Lett.*, **117**, 180801, Oct. 2016. <http://arxiv.org/abs/1607.00130> <http://link.aps.org/doi/10.1103/PhysRevLett.117.180801>.
- [96] A. Furusawa, J. L. Sørensen, S. L. Braunstein, C. A. Fuchs, H. J. Kimble, and E. S. Polzik, “Unconditional quantum teleportation,” *Science*, **282**, pp. 706–9, Oct. 1998. <http://www.ncbi.nlm.nih.gov/pubmed/9784123>.
- [97] W. P. Bowen, R. Schnabel, and P. K. Lam, “Experimental Investigation of Criteria for Continuous Variable Entanglement,” *Phys. Rev. Lett.*, **90**, no. 4, 43601, 2003. <http://link.aps.org/doi/10.1103/PhysRevLett.90.043601>.



- [98] T. Eberle, V. Händchen, and R. Schnabel, “Stable Control of 10 dB Two-Mode Squeezed Vacuum States of Light,” *Optics Express*, 2013. <http://www.opticsinfobase.org/oe/abstract.cfm?uri=oe-21-9-11546>.
- [99] R. Adhikari, “Upgrades to the Advanced LIGO Interferometer,” *LIGO Document: LIGO-G1000524*, 2010. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=11613>.
- [100] G. Rempe, R. J. Thompson, H. J. Kimble, and R. Lalezari, “Measurement of Ultralow Losses in an Optical Interferometer,” *Opt. Lett.*, **17**, pp. 363–365, 1992. <http://www.opticsinfobase.org/abstract.cfm?id=11045>.
- [101] N. Uehara, A. Ueda, K. Ueda, H. Sekiguchi, T. Mitake, and et al., “Ultralow-loss mirror of the parts-in  $10^{-6}$  level at 1064 nm,” *Opt. Lett.*, **20**, pp. 530–532, 1995. <http://www.opticsinfobase.org/abstract.cfm?uri=ol-20-6-530>.
- [102] T. Zhang, H. Yu, V. Sudhir, C. Torrie, K. Strain, H. Grote, K. Arai, M. Zucker, R. Adhikari, S. Steinlechner, J. Briggs, P. Fritschel, and S. Hild, “Plans for the A+ Balanced Homodyne readout,” *LIGO Document: LIGO-G1800459*, 2018. <https://dcc.ligo.org/LIGO-G1800459>.
- [103] S. Steinlechner, B. W. Barr, A. S. Bell, S. L. Danilishin, A. Gläflke, C. Gräf, J.-S. Hennig, E. A. Houston, S. H. Huttner, S. S. Leavey, D. Pascucci, B. Sorazu, A. Spencer, K. A. Strain, J. Wright, and S. Hild, “Local-oscillator noise coupling in balanced homodyne readout for advanced gravitational wave detectors,” *Physical Review D*, **92**, 072009, oct 2015. <http://link.aps.org/doi/10.1103/PhysRevD.92.072009>.
- [104] B. Hage, A. Sambrowski, and R. Schnabel, “Towards Einstein-Podolsky-Rosen quantum channel multiplexing,” *Phys. Rev. A*, **81**, 62301, June 2010.
- [105] D. D. Brown, H. Miao, C. Collins, C. Mow-Lowry, D. Töyrä, and A. Freise, “Broadband sensitivity enhancement of detuned dual-recycled Michelson interferometers with EPR entanglement,” *Phys. Rev. D*, **96**, 062003, Sep 2017.
- [106] A. Thüring, R. Schnabel, H. Lück, and K. Danzmann, “Detuned Twin-Signal-Recycling for ultra-high precision interferometers,” *Opt. Lett.*, **32**, pp. 985–987, 2007.
- [107] A. Thuring, C. Graf, H. Vahlbruch, M. Mehmet, K. Danzmann, and R. Schnabel, “Broadband squeezing of quantum noise in a Michelson interferometer with Twin-Signal-Recycling,” *Opt. Lett.*, 181101, 2009.
- [108] V. B. Braginsky and F. Y. Khalili, *Quantum Measurement*. ed. K. S. Thorne, Cambridge University Press, 1992. <http://www.amazon.com/Quantum-Measurement-Vladimir-B-Braginsky/dp/052141928X>.
- [109] P. Purdue and Y. Chen, “Practical speed meter designs for quantum nondemolition gravitational-wave interferometers,” *Phys. Rev. D*, **66**, 122004, 2002.

- [110] H. Rehbein, H. Müller-Ebhardt, K. Somiya, C. Li, R. Schnabel, K. Danzmann, and Y. Chen, “Local readout enhancement for detuned signal-recycling interferometers,” *Phys. Rev. D*, **76**, 062002, 2007. <http://link.aps.org/doi/10.1103/PhysRevD.76.062002>.
- [111] J. Cripe, N. Aggarwal, R. Lanza, A. Libson, R. Singh, P. Heu, D. Follman, G. D. Cole, N. Mavalvala, and T. Corbitt, “Observation of a room-temperature oscillator’s motion dominated by quantum fluctuations over a broad audio-frequency band,” *ArXiv e-prints*, arXiv:1802.10069, Feb. 2018.
- [112] H. Miao, H. Yang, R. Adhikari, and Y. Chen, “Comparison of Quantum Noise in 3G Interferometer Configurations,” *LIGO Document: LIGO-T1200008-v3*, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=78229>.
- [113] N. A. Robertson, G. Cagnoli, D. R. M. Crooks, E. Elliffe, J. E. Faller, P. Fritschel, S. Goßler, A. Grant, A. Heptonstall, J. Hough, H. Lück, R. Mittleman, M. Perreurlloyd, M. V. Plissi, S. Rowan, D. H. Shoemaker, P. H. Sneddon, K. A. Strain, C. I. Torrie, H. Ward, and P. Willems, “Quadruple suspension design for Advanced LIGO,” *Class. Quantum Grav.*, **19**, pp. 4043–4058, 2002.
- [114] A. V. Cumming, A. S. Bell, L. Barsotti, M. A. Barton, G. Cagnoli, D. Cook, L. Cunningham, M. Evans, G. D. Hammond, G. M. Harry, A. Heptonstall, J. Hough, R. Jones, R. Kumar, R. Mittleman, N. A. Robertson, S. Rowan, B. Shapiro, K. A. Strain, K. Tokmakov, C. Torrie, and A. A. van Veggel, “Design and development of the advanced LIGO monolithic fused silica suspension,” *Class. Quantum Grav.*, **29**, no. 3, 035003, 2012. <http://stacks.iop.org/0264-9381/29/i=3/a=035003>.
- [115] N. A. Robertson for the LSC, “Seismic isolation and suspension systems for Advanced LIGO,” *Proceedings of SPIE*, **5500**, 81, 2004.
- [116] P. R. Saulson, “Terrestrial gravitational noise on a gravitational wave antenna,” *Phys. Rev. D*, **30**, pp. 732–736, 1984.
- [117] T. Hong, H. Yang, E. K. Gustafson, R. Adhikari, and Y. Chen, “Brownian Thermal Noise in Multilayer Coated Mirrors,” *in prep.*, 2012. in prep.
- [118] T. P. Purdy, P.-L. Yu, R. W. Peterson, N. S. Kampel, and C. A. Regal, “Strong Optomechanical Squeezing of Light,” *Phys. Rev. X*, **3**, 031012, Sept. 2013. <http://link.aps.org/doi/10.1103/PhysRevX.3.031012>.
- [119] T. Corbitt, Y. Chen, F. Khalili, D. Ottaway, S. Vyatchanin, S. Whitcomb, and N. Mavalvala, “Squeezed-state source using radiation-pressure-induced rigidity,” *Phys. Rev. A*, **73**, 023801, 2006. <http://link.aps.org/doi/10.1103/PhysRevA.73.023801>.
- [120] M. Korobko, F. Khalili, and R. Schnabel, “Engineering the optical spring via intra-cavity optical-parametric amplification,” *Physics Letters A*, 2017. <http://www.sciencedirect.com/science/article/pii/S0375960117303146>.

- [121] M. Zhou, Z. Zhou, and S. M. Shahriar, “Quantum noise limits in white-light-cavity-enhanced gravitational wave detectors,” *Phys. Rev. D*, **92**, 082002, Oct 2015. <http://link.aps.org/doi/10.1103/PhysRevD.92.082002>.
- [122] Y. Ma, H. Miao, C. Zhao, and Y. Chen, “Quantum noise of a white-light cavity using a double-pumped gain medium,” *Phys. Rev. A*, **92**, 023807, Aug 2015. <https://link.aps.org/doi/10.1103/PhysRevA.92.023807>.
- [123] M. Zhou, Z. Zhou, and S. M. Shahriar, “Catastrophic breakdown of the Caves model for quantum noise in some phase-insensitive linear amplifiers or attenuators based on atomic systems,” *Phys. Rev. A*, **93**, 033858, Mar 2016. <http://link.aps.org/doi/10.1103/PhysRevA.93.033858>.
- [124] H. Miao, Y. Ma, C. Zhao, and Y. Chen, “Enhancing the Bandwidth of Gravitational-Wave Detectors with Unstable Optomechanical Filters,” *Phys. Rev. Lett.*, **115**, 211104, Nov 2015. <https://link.aps.org/doi/10.1103/PhysRevLett.115.211104>.
- [125] Y. Ma, S. L. Danilishin, C. Zhao, H. Miao, W. Z. Korth, Y. Chen, R. L. Ward, and D. G. Blair, “Narrowing the Filter-Cavity Bandwidth in Gravitational-Wave Detectors via Optomechanical Interaction,” *Phys. Rev. Lett.*, **113**, 151102, Oct 2014. <https://link.aps.org/doi/10.1103/PhysRevLett.113.151102>.
- [126] E. E. Mikhailov, K. Goda, T. Corbitt, and N. Mavalvala, “Frequency-dependent squeeze-amplitude attenuation and squeeze-angle rotation by electromagnetically induced transparency for gravitational-wave interferometers,” *Phys. Rev. A*, **73**, 053810, 2006. <http://link.aps.org/doi/10.1103/PhysRevA.73.053810>.
- [127] M. Zhou, Z. Zhou, and S. M. Shahriar, “Realizing the GEIT System Using Zeeman Sublevels in Rb for Enhancing the Sensitivity-Bandwidth Product in Next Generation LIGO,” 2016. <https://dcc.ligo.org/G1600440>.
- [128] L.-M. Duan, J. I. Cirac, P. Zoller, and E. S. Polzik, “Quantum Communication between Atomic Ensembles Using Coherent Light,” *Phys. Rev. Lett.*, **85**, pp. 5643–5646, Dec 2000. <https://link.aps.org/doi/10.1103/PhysRevLett.85.5643>.
- [129] B. Julsgaard, A. Kozhekin, and E. S. Polzik, “Experimental long-lived entanglement of two macroscopic objects,” *Nature*, **413**, pp. 400–403, sep 2001.
- [130] C. B. Møller, R. A. Thomas, G. Vasilakis, E. Zeuthen, Y. Tsaturyan, M. Balabas, K. Jensen, A. Schliesser, K. Hammerer, and E. S. Polzik, “Quantum back-action-evading measurement of motion in a negative mass reference frame,” *Nature*, **547**, pp. 191–195, jul 2017.
- [131] Y. Levin, “Internal thermal noise in the LIGO test masses: A direct approach,” *Phys. Rev. D*, **57**, pp. 659–663, 1998. <http://link.aps.org/doi/10.1103/PhysRevD.57.659>.
- [132] W. A. Phillips, “Two-level states in glasses,” *Reports on Progress in Physics*, **50**, 1657, 1987. <http://dx.doi.org/10.1088/0034-4885/50/12/003>.

- [133] “Spectral Shape of Relaxations in Silica Glass,” *Physical Review Letters*, **84**, 2718, 2000.
- [134] S. Chao *et al.*, “Low Loss Dielectric Mirrors with Ion Beam Sputtered  $TiO_2 - SiO_2$  Mixed Films,” *Opt. Lett.*, **40**, 2177, 2001.
- [135] S. Ushakov *et al.*, “Crystalization in Hafnia and Zirconia Based Systems,” *Phys. Stat. Sol.*, **241**, 2268, 2004.
- [136] S. Chao *et al.*, “Thickness-dependent crystallization on thermal anneal for the titania/silica nano-layers deposited by ion-beam-sputter method.” LIGO Document: LIGO G1300921.
- [137] M. Liu *et al.*, “Microstructure and Interfacial Properties of  $HfO_2/Al_2O_3$  Nanolaminate Films,” *Appl. Surf. Sci.*, **252**, 6206, 2006.
- [138] S. D. Penn, P. H. Sneddon, H. Armandula, J. C. Betzwieser, G. Cagnoli, J. Camp, D. R. M. Crooks, M. M. Fejer, A. M. Gretarsson, G. M. Harry, J. Hough, S. E. Kittelberger, M. J. Mortonson, R. Route, S. Rowan, and C. C. Vassiliou, “Mechanical loss in tantala/silica dielectric mirror coatings,” *Class. Quantum Grav.*, **20**, no. 13, 2917, 2003. <http://dx.doi.org/10.1088/0264-9381/20/13/334>.
- [139] I. Pinto *et al.*, “Interdiffused coatings.” LIGO Document G-1200976.
- [140] L. Pinard, C. Michel, B. Sassolas, L. Balzarini, J. Degallaix, J. Dolique, R. Flaminio, D. Forest, M. Granata, B. Lagrange, N. Straniero, J. Teillon, and G. Cagnoli, “The Mirrors Used in the LIGO Interferometers for the First-time Detection of Gravitational Waves,” , in *Optical Interference Coatings 2016*, MB.3, Optical Society of America, 2016. <http://www.osapublishing.org/abstract.cfm?URI=OIC-2016-MB.3>.
- [141] V. B. Braginsky, M. L. Gorodetsky, and S. P. Vyatchanin, “Thermodynamical fluctuations and photo-thermal shot noise in gravitational wave antennae,” *Phys. Lett. A*, **264**, 1, 1999.
- [142] V. B. Braginsky, M. L. Gorodetsky, and S. P. Vyatchanin, “Thermo-refractive noise in gravitational wave antennae,” *Phys. Lett. A*, **271**, 303, 2000.
- [143] Y. Levin, “Fluctuation-dissipation theorem for thermo-refractive noise,” *Physics Letters A*, **372**, no. 12, pp. 1941 – 1944, 2008. <http://www.sciencedirect.com/science/article/pii/S0375960107015964>.
- [144] G. H. Ogin, *Measurement of thermo-optic properties of thin film dielectric coatings*. California Institute of Technology, 2013.
- [145] T. Chalermsongsak, E. D. Hall, G. D. Cole, D. Follman, F. Seifert, K. Arai, E. K. Gustafson, J. R. Smith, M. Aspelmeyer, and R. X. Adhikari, “Coherent cancellation of photothermal noise in GaAs/Al<sub>0.92</sub>Ga<sub>0.08</sub>As Bragg mirrors,” *Metrologia*, **53**, no. 2, 860, 2016.
- [146] V. Braginsky and S. Vyatchanin, “Thermodynamical fluctuations in optical mirror coatings,” *Physics Letters A*, **312**, no. 3-4, pp. 244–255, 2003.

- [147] M. N. Inci, “Simultaneous measurements of the thermal optical and linear thermal expansion coefficients of a thin film etalon from the reflection spectra of a superluminescent diode,” *Journal of Physics D: Applied Physics*, **37**, no. 22, 3151, 2004.
- [148] D. Crooks, G. Cagnoli, M. Fejer, A. Gretarsson, G. Harry, J. Hough, N. Nakagawa, S. Penn, R. Route, S. Rowan, *et al.*, “Experimental measurements of coating mechanical loss factors,” *Classical and Quantum Gravity*, **21**, no. 5, S1059, 2004.
- [149] A. Gretarsson, “Thermo-optic noise from doped tantala/silica coatings,” *LIGO DCC*, pages G080151–00, 2008.
- [150] E. Çetinörgü, B. Baloukas, O. Zabeida, J. E. Klemberg-Sapieha, and L. Martinu, “Mechanical and thermoelastic characteristics of optical thin films deposited by dual ion beam sputtering,” *Applied optics*, **48**, no. 23, pp. 4536–4544, 2009.
- [151] H.-W. Pan, S.-J. Wang, L.-C. Kuo, S. Chao, M. Principe, I. M. Pinto, and R. DeSalvo, “Thickness-dependent crystallization on thermal anneal for titania/silica nm-layer composites deposited by ion beam sputter method,” *Opt. Express*, **22**, pp. 29847–29854, Dec 2014. <http://www.opticsexpress.org/abstract.cfm?URI=oe-22-24-29847>.
- [152] R. Bassiri, K. B. Borisenko, D. J. H. Cockayne, J. Hough, I. MacLaren, and S. Rowan, “Probing the atomic structure of amorphous Ta<sub>2</sub>O<sub>5</sub> coatings,” *Applied Physics Letters*, **98**, no. 3, 031904, 2011.
- [153] R. Bassiri, K. Evans, K. Borisenko, M. Fejer, J. Hough, I. MacLaren, I. Martin, R. Route, and S. Rowan, “Correlations between the mechanical loss and atomic structure of amorphous TiO<sub>2</sub>-doped Ta<sub>2</sub>O<sub>5</sub> coatings,” *Acta Materialia*, **61**, no. 4, pp. 1070–1077, 2013.
- [154] R. Bassiri, F. Liou, M. R. Abernathy, A. C. Lin, N. Kim, A. Mehta, B. Shyam, R. L. Byer, E. K. Gustafson, M. Hart, *et al.*, “Order within disorder: The atomic structure of ion-beam sputtered amorphous tantala (a-Ta<sub>2</sub>O<sub>5</sub>),” *APL materials*, **3**, no. 3, 036103, 2015.
- [155] R. Bassiri, M. R. Abernathy, F. Liou, A. Mehta, E. K. Gustafson, M. J. Hart, H. N. Isa, N. Kim, A. C. Lin, I. MacLaren, *et al.*, “Order, disorder and mixing: The atomic structure of amorphous mixtures of titania and tantala,” *Journal of Non-Crystalline Solids*, **438**, pp. 59–66, 2016.
- [156] B. Shyam, K. H. Stone, R. Bassiri, M. M. Fejer, M. F. Toney, and A. Mehta, “Measurement and modeling of short and medium range order in amorphous Ta<sub>2</sub>O<sub>5</sub> thin films,” *Scientific reports*, **6**, 32170, 2016.
- [157] M. J. Hart, R. Bassiri, K. B. Borisenko, M. Véron, E. F. Rauch, I. W. Martin, S. Rowan, M. M. Fejer, and I. MacLaren, “Medium range structural order in amorphous tantala spatially resolved with changes to atomic structure by thermal annealing,” *Journal of Non-Crystalline Solids*, **438**, pp. 10–17, 2016.
- [158] J. P. Trinastic, R. Hamdan, Y. Wu, L. Zhang, and H.-P. Cheng, “Unified interatomic potential and energy barrier distributions for amorphous oxides,” *The Journal of Chemical Physics*, **139**, no. 15, 154506, 2013.

- [159] J. P. Trinastic, R. Hamdan, C. Billman, and H.-P. Cheng, “Molecular dynamics modeling of mechanical loss in amorphous tantala and titania-doped tantala,” *Physical Review B*, **93**, no. 1, 014105, 2016.
- [160] S. Singh, M. D. Ediger, and J. J. De Pablo, “Ultrastable glasses from in silico vapour deposition,” *Nature materials*, **12**, no. 2, 139, 2013.
- [161] I. Lyubimov, M. D. Ediger, and J. J. de Pablo, “Model vapor-deposited glasses: Growth front and composition effects,” *The Journal of chemical physics*, **139**, no. 14, 144505, 2013.
- [162] S. Gras, H. Yu, W. Yam, D. Martynov, and M. Evans, “Audio-band coating thermal noise measurement for Advanced LIGO with a multimode optical resonator,” *Phys. Rev. D*, **95**, 022001, Jan 2017. <https://link.aps.org/doi/10.1103/PhysRevD.95.022001>.
- [163] M. Principe *et al.*, “Chapter 12, Reflectivity and thickness optimization,” , in *Optical Coatings and Thermal Noise in Precision Measurements* (G. Harry *et al.*, eds.), Cambridge University Press, 2012.
- [164] A. Villar *et al.*, “Measurement of Thermal Noise in Multilayer Coatings with Optimized Layer Thicknesses,” *Phys. Rev. D*, **81**, 122001, 2010.
- [165] W. Yam, S. Gras, and M. Evans, “Multimaterial coatings with reduced thermal noise,” *Phys. Rev. D*, **91**, 042002, Feb 2015. <https://link.aps.org/doi/10.1103/PhysRevD.91.042002>.
- [166] J. Steinlechner, I. W. Martin, J. Hough, C. Krüger, S. Rowan, and R. Schnabel, “Thermal noise reduction and absorption optimization via multimaterial coatings,” *Phys. Rev. D*, **91**, 042001, Feb 2015. <https://link.aps.org/doi/10.1103/PhysRevD.91.042001>.
- [167] M. Evans, L. Barsotti, and P. Fritschel, “A general approach to optomechanical parametric instabilities,” *Phys. Lett. A*, **374**, no. 4, pp. 665–671, 2010. <http://dx.doi.org/10.1016/j.physleta.2009.11.023>.
- [168] C. Zhao *et al.* *Phys. Rev. Lett.*, **4**, 121102, 2005.
- [169] S. Gras *et al.* *Class. Quantum Grav.*, **27**, 205019, 2010.
- [170] S. Strigin *et al.* *Phys. Lett. A*, **372**, 5727, 2008.
- [171] B. Lantz, “Simple Calculations for Charge Noise for Advanced LIGO,” *LIGO Document: LIGO-T080214*, 2010. <https://dcc.ligo.org/LIGO-T080214>.
- [172] R. Weiss, “Note on Electrostatics in the LIGO Suspensions,” *LIGO Document: LIGO-T960137*, 1996. <https://dcc.ligo.org/LIGO-T960137>.
- [173] D. Ugolini and C. Fitzgerald, “Developing a Charging Testbed in an Advanced LIGO Geometry,” *LIGO Document: LIGO-G1200853*, 2012. <https://dcc.ligo.org/DocDB/0096/G1200853/001/G1200853.pdf>.

- [174] V. Mitrofanov and L. Prokhorov, “Space charge polarization in fused silica test masses of gravitational wave detector associated with the electrostatic drive,” *LIGO Document: LIGO-P1000077*, 2010. <https://dcc.ligo.org/LIGO-P1000077>.
- [175] R. Weiss, “ionic\_neutralization,” *LIGO Document: LIGO-G1000383*, 2010. <https://dcc.ligo.org/LIGO-G1000383>.
- [176] D. Ugolini, “Discharging Optics with Ionized Gases,” *LIGO Document: LIGO-T1000135*, 2010. <https://dcc.ligo.org/LIGO-T1000135>.
- [177] R. Weiss, “Surface charge control of the Advanced LIGO mirrors using externally introduced ions,” *LIGO Document: LIGO-T1100332*, 2011. <https://dcc.ligo.org/LIGO-T1100332>.
- [178] G. D. Cole, “Outlook on Size-Scaling of Substrate-Transferred Crystalline Coatings,” 2018. <https://dcc.ligo.org/LIGO-G1801024>.
- [179] I. W. Martin, E. Chalkley, R. Nawrodt, H. Armandula, R. Bassiri, C. Comtet, M. Fejer, A. Gretarsson, G. Harry, D. Heinert, *et al.*, “Comparison of the temperature dependence of the mechanical dissipation in thin films of Ta<sub>2</sub>O<sub>5</sub> and Ta<sub>2</sub>O<sub>5</sub> doped with TiO<sub>2</sub>,” *Classical and Quantum Gravity*, **26**, no. 15, 155012, 2009.
- [180] I. W. Martin, R. Bassiri, R. Nawrodt, M. Fejer, A. Gretarsson, E. Gustafson, G. Harry, J. Hough, I. MacLaren, S. Penn, *et al.*, “Effect of heat treatment on mechanical dissipation in Ta<sub>2</sub>O<sub>5</sub> coatings,” *Classical and Quantum Gravity*, **27**, no. 22, 225020, 2010.
- [181] G. D. Cole, W. Zhang, M. J. Martin, J. Ye, and M. Aspelmeyer, “Tenfold reduction of Brownian noise in high-reflectivity optical coatings,” *Nature Photonics*, **7**, no. 8, 644, 2013.
- [182] A. C. Lin, R. Bassiri, S. Omar, A. S. Markosyan, B. Lantz, R. Route, R. L. Byer, J. S. Harris, and M. M. Fejer, “Epitaxial growth of GaP/AlGaP mirrors on Si for low thermal noise optical coatings,” *Optical Materials Express*, **5**, no. 8, pp. 1890–1897, 2015.
- [183] A. Cumming, K. Craig, I. Martin, R. Bassiri, L. Cunningham, M. Fejer, J. Harris, K. Haughian, D. Heinert, B. Lantz, *et al.*, “Measurement of the mechanical loss of prototype GaP/AlGaP crystalline coatings for future gravitational wave detectors,” *Classical and Quantum Gravity*, **32**, no. 3, 035002, 2015.
- [184] I. W. Martin, R. Nawrodt, K. Craig, C. Schwarz, R. Bassiri, G. Harry, J. Hough, S. Penn, S. Reid, R. Robie, *et al.*, “Low temperature mechanical dissipation of an ion-beam sputtered silica film,” *Classical and Quantum Gravity*, **31**, no. 3, 035019, 2014.
- [185] P. G. Murray, I. W. Martin, K. Craig, J. Hough, R. Robie, S. Rowan, M. R. Abernathy, T. Pershing, and S. Penn, “Ion-beam sputtered amorphous silicon films for cryogenic precision measurement systems,” *Phys. Rev. D*, **92**, 062001, Sep 2015. <https://link.aps.org/doi/10.1103/PhysRevD.92.062001>.

- [186] J. Steinlechner, I. Martin, J. Hough, S. Rowan, R. Schnabel, M. Fletcher, R. Robie, P. Murray, and A. Bell, “Silicon-based optical mirror coatings for ultra-high precision metrology and sensing,” *LIGO Document: LIGO-P1800004*, 2018. <https://dcc.ligo.org/LIGO-P1800004>.
- [187] I.Martin *et al.*, “Mechanical loss of crystalline and amorphous coatings.” GWADW 2014. [http://www.gravity.ircs.titech.ac.jp/GWADW2014/slide/Iain\\_Martin.pdf](http://www.gravity.ircs.titech.ac.jp/GWADW2014/slide/Iain_Martin.pdf).
- [188] S. Chao, H.-W. Pan, and L.-C. Kuo, “Mechanical loss of silicon nitride films grown by plasma-enhanced chemical vapor deposition (PECVD) method,” *LIGO Document: LIGO-G1300171*, 2013. <https://dcc.ligo.org/LIGO-G1300171>.
- [189] S. Chao, H.-W. Pan, Y.-H. Juang, and S.-Y. Huang, “Mechanical loss of silicon cantilever coated with a high-stress SiNx film,” *LIGO Document: LIGO-G1400851*, 2014. <https://dcc.ligo.org/LIGO-G1400851>.
- [190] I. Pinto, “nm-layered coatings: status and perspectives.” ELiTES: Second general meeting 2013. <https://events.ego-gw.it/indico/conferenceOtherViews.py?view=standard&confId=7>.
- [191] J. Agresti, G. Castaldi, R. DeSalvo, V. Galdi, V. Pierro, and I. M. Pinto, “Optimized multilayer dielectric mirror coatings for gravitational wave interferometers,” , in *Advances in Thin-Film Coatings for Optical Applications III*, vol. 6286, 628608, International Society for Optics and Photonics, 2006.
- [192] A. E. Villar, E. D. Black, R. DeSalvo, K. G. Libbrecht, C. Michel, N. Morgado, L. Pinard, I. M. Pinto, V. Pierro, V. Galdi, *et al.*, “Measurement of thermal noise in multilayer coatings with optimized layer thickness,” *Physical Review D*, **81**, no. 12, 122001, 2010.
- [193] B. Lantz, R. Schofield, B. O’Reilly, D. E. Clark, and D. DeBra, “Review: Requirements for a Ground Rotation Sensor to Improve Advanced LIGO,” *Bulletin of the Seismological Society of America*, **99**, pp. 980–989, 2009. <http://bssa.geoscienceworld.org/cgi/content/abstract/99/2B/980>.
- [194] V. Dergachev and R. DeSalvo, “A high precision mechanical ground rotation sensor ,” *LIGO Document: LIGO-G1200187*, 2012. <https://dcc.ligo.org/G1200187>.
- [195] K. Venkateswara, “TiltWash: Update on tiltmeter work at UW,” *LIGO Document: LIGO-G1200225*, 2012. <https://dcc.ligo.org/G1200225>.
- [196] W. Korth, A. Heptonstall, R. Adhikari, E. Gustafson, and B. Lantz, “Status of the Caltech Laser Gyro,” *LIGO Document: LIGO-G1200261*, 2012. <https://dcc.ligo.org/G1200261>.
- [197] M. Evans and F. Matchard, “Tilt Free Inertial Sensing,” *LIGO Document: LIGO-T0900628*, 2009. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=7974>.



- [198] K. Venkateswara, “Improving active seismic isolation in aLIGO using a ground rotation sensor,” *LIGO Document: LIGO-G1600083*, 2016. <https://dcc.ligo.org/LIGO-G1600083>.
- [199] G. Rajalakshmi, P. K. Madhu, and C. Unnikrishnan, “Atomic Spin Gyroscopes for Rotation Sensing,” *LIGO Document: LIGO-G1601427*, 2016. <https://dcc.ligo.org/LIGO-G1601427>.
- [200] T. W. Kornack, R. K. Ghosh, and M. V. Romalis, “Nuclear Spin Gyroscope Based on an Atomic Comagnetometer,” *Phys. Rev. Lett.*, **95**, 230801, Nov 2005. <http://link.aps.org/doi/10.1103/PhysRevLett.95.230801>.
- [201] B. Lantz, “Thoughts on Hanford Wind,” *LIGO Document: LIGO-G1501371*, 2015. <https://dcc.ligo.org/LIGO-G1501371>.
- [202] D. Tshilumba, L. Nuttall, T. Mac Donald, R. Mittleman, B. Lantz, F. Matichard, and C. Collette, “Vibration analysis and control of the LIGO observatories large chambers and support piers,” , in *PROCEEDINGS OF ISMA2014 INCLUDING USD2014*, pp. 187–200, 2014. <https://dcc.ligo.org/LIGO-p1400109>.
- [203] T. MacDonald, “BSC HEPI Motion Measurements and Modeling,” *LIGO Document: LIGO-G1401167*, 2014. <https://dcc.ligo.org/G1401167>.
- [204] J. Driggers, V. Frolov, D. Atkinson, H. Miao, M. Landry, R. Adhikari, and R. DeRosa, “Global feed-forward vibration isolation in a km scale interferometer,” *LIGO-Documents: LIGO-P1000088*, 2010. <https://dcc.ligo.org/P1000088>.
- [205] B. N. Shapiro, *Adaptive Modal Damping for Advance LIGO Suspensions*. PhD thesis, Massachusetts Institute of Technology, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=91198>.
- [206] R. Kurdyumov, C. Kucharczyk, and B. Lantz, “Blend Switching User Guide,” *LIGO-Documents: LIGO-T1200126*, 2012. <https://dcc.ligo.org/T1200126>.
- [207] M. Coughlin *et al.*, “Earthquake monitoring for aLIGO,” *LIGO Document: LIGO-G1400811*, 2014. <https://dcc.ligo.org/G1400811>.
- [208] N. Mukund, “update on lockloss monitoring & prediction,” 2018. LIGO-G1800569, <https://dcc.ligo.org/LIGO-G1800569>.
- [209] B. Lantz, S. Biscans, J. Warner, A. Pele, H. Radkins, N. Mukund, and M. Coughlin, “Seismic Control during Earthquakes: Review of new technique,” 2018. LIGO-G1800399, <https://dcc.ligo.org/LIGO-G1800399>.
- [210] C. Collette and F. Matichard, “Sensor Fusion Methods for High Performance Active Vibration Isolation Systems,” *Journal of Sound and Vibration*, no. P1400022, 2014. <https://dcc.ligo.org/LIGO-P1400022>.
- [211] C. Collette and F. Matichard, “Vibration control of flexible structures using fusion of inertial sensors and hyper-stable actuator/sensor pairs,” *LIGO Document: LIGO-P1400099*, 2014. <https://dcc.ligo.org/LIGO-P1400099>.

- [212] K.-F. J. Tseng, “System Failure Diagnosis for the Advanced LIGO HAM Chamber Seismic Isolation System,” *LIGO Document: LIGO-P1000086*, 2010. <https://dcc.ligo.org/P1000086>.
- [213] B. N. Shapiro, R. Adhikari, J. Driggers, J. Kissel, B. Lantz, J. Rollins, and K. Youcef-Toumi, “Noise and control decoupling of Advanced LIGO suspensions,” *Classical and Quantum Gravity*, **32**, no. 1, 015004, 2015. <http://stacks.iop.org/0264-9381/32/i=1/a=015004>.
- [214] E. Bonilla, “Improving DARM with ISI - $\zeta$  SUS Feedforward,” 2018. LIGO-G18003467, <https://dcc.ligo.org/LIGO-G1800467>.
- [215] P. Fritschel, “Potential Advanced LIGO post-Project upgrades,” Tech. Rep. T1300176, LIGO Laboratory, Mar 2013. <https://dcc.ligo.org/LIGO-T1300176>.
- [216] N. Robertson, “Beamsplitter Actuation: A Note on Sweetspot Positioning and Magnet Size Considerations,” *LIGO Document: LIGO-T1600328*, 2016. <https://dcc.ligo.org/LIGO-T1600328>.
- [217] N. A. Robertson, H. Miller, and C. Torrie, “Risk Reduction Work on Advanced LIGO Suspensions,” *LIGO Document: LIGO-G1400804*, 2014. <https://dcc.ligo.org/LIGO-G1400804>.
- [218] N. Robertson and M. Barton, “Design of a Larger Beamsplitter Suspension,” *LIGO Document: LIGO-T1400296*, 2014. <https://dcc.ligo.org/LIGO-T1400296>.
- [219] N. Robertson and H. Miller, “Revised Design of HSTS with Improved Vertical Isolation,” *LIGO Document: LIGO-T1400290*, 2014. <https://dcc.ligo.org/LIGO-T1400290>.
- [220] N. Robertson, “Relook at Signal Recycling Cavity Displacement Noise,” *LIGO Document: LIGO-T1500323*, 2015. <https://dcc.ligo.org/LIGO-T1500323>.
- [221] N. Robertson, C. Torrie, E. Sanchez, and R. Abbott, “BRDs for HSTS: Design, Results and Conclusions,” *LIGO Document: LIGO-T1700111*, 2017. <https://dcc.ligo.org/LIGO-T1700111>.
- [222] N. Robertson, C. Torrie, and E. Sanchez, “BRDs for HLTS: Design, Results and Conclusions,” *LIGO Document: LIGO-T1700155*, 2017. <https://dcc.ligo.org/LIGO-T1700155>.
- [223] N. A. Robertson and J. Hough, “Gas Damping in Advanced LIGO Suspensions,” *LIGO Document: LIGO-T0900416-v2*, 2009. <https://dcc.ligo.org/cgi-bin/DocDB/ShowDocument?docid=5085>.
- [224] A. Cavalleri, G. Ciani, R. Dolesi, A. Heptonstall, M. Hueller, D. Nicolodi, S. Rowan, D. Tombolato, S. Vitale, P. J. Wass, and W. J. Weber, “Increased Brownian Force Noise from Molecular Impacts in a Constrained Volume,” *Phys. Rev. Lett.*, **103**, 140601, 2009. <http://link.aps.org/doi/10.1103/PhysRevLett.103.140601>.

- [225] R. Dolesi, M. Hueller, D. Nicolodi, D. Tombolato, S. Vitale, P. J. Wass, W. J. Weber, M. Evans, P. Fritschel, R. Weiss, J. H. Gundlach, C. A. Hagedorn, S. Schlamminger, G. Ciani, and A. Cavalleri, “Brownian force noise from molecular collisions and the sensitivity of advanced gravitational wave observatories,” *Phys. Rev. D*, **84**, 063007, 2011. <http://link.aps.org/doi/10.1103/PhysRevD.84.063007>.
- [226] G. Billingsley, N. Robertson, and B. Shapiro, “Annular End Reaction Mass Conceptual/Final Design Document,” *LIGO Document: LIGO-E1500264*, 2015. <https://dcc.ligo.org/LIGO-E1500264>.
- [227] N. Robertson and C. Torrie, “Recent Results and Conclusions from Tests of the UIM Blade Non-Magnetic Damper,” *LIGO Document: LIGO-T1600046*, 2016. <https://dcc.ligo.org/LIGO-T1600046>.
- [228] N. Robertson and C. Torrie, “A Tale of Two Dampers: Bounce and Roll Mode Damping for the Quads,” *LIGO Document: LIGO-G1600371*, 2016. <https://dcc.ligo.org/LIGO-G1600371>.
- [229] N. Robertson, P. Fritschel, B. Shapiro, C. Torrie, and M. Evans, “Design of a Tuned Mass Damper for High Quality Factor Suspension Modes in Advanced LIGO,” *Review of Scientific Instruments*, **88**, no. 3, 035117, 2017.
- [230] N. Robertson, C. Torrie, and E. Sanchez, “BRDs for the Beamsplitter: Design, Results and Conclusions,” *LIGO Document: LIGO-T1700176*, 2017. <https://dcc.ligo.org/LIGO-T1700176>.
- [231] R. Mittleman et al., “Violin Mode Damper,” *LIGO document G1701702*, 2017. <https://dcc.ligo.org/LIGO-G1701702>.
- [232] Norna A Robertson, “Simple Model for Analysing Passive Violin Mode Damping,” *LIGO document T1800161*, 2018. <https://dcc.ligo.org/LIGO-T1800161>.
- [233] M. Zucker et al., “The A+ Upgrade: Expanding the Advanced LIGO Horizon,” *LIGO Document: LIGO-M1800040*, 2018. <https://dcc.ligo.org/LIGO-M1800040>.
- [234] M. Zucker et al., “A+ Overview,” *LIGO Document: LIGO-G1800514*, 2018. <https://dcc.ligo.org/LIGO-G1800514>.
- [235] G. Hammond *et al.*, “Progress on silica and silicon suspensions,” *LIGO Document: LIGO-G1400849*, 2014. <https://dcc.ligo.org/G1400849>.
- [236] D. Madden-Fong, “LIGO III Quad Pendulum Conceptual Design Optimization,” *LIGO Document: LIGO-T1300786*, 2013. <https://dcc.ligo.org/T1300786>.
- [237] B. Shapiro, “Alternative design approach for caging, sensing and control next generation suspensions,” *LIGO Document: LIGO-G1601426*, 2016. <https://dcc.ligo.org/LIGO-G1601426>.
- [238] N. Lockerbie, K. Tokmakov, and S. Jawahar, “New OSEM results using a displacement-doubling prism-based flag,” *LIGO Document: LIGO-G1601136*, 2016. <https://dcc.ligo.org/LIGO-G1601136>.

- [239] C. C. Speake and S. M. Aston, “An interferometric sensor for satellite drag-free control,” *Class. Quantum Grav.*, **22**, S269, 2005. <http://dx.doi.org/10.1088/0264-9381/22/10/019>.
- [240] S. J. Cooper, C. J. Collins, A. C. Green, D. Hoyland, C. C. Speake, A. Freise and C. M. Mow-Lowry, “A compact, large-range interferometer for precision measurement and inertial sensing,” *Classical and Quantum Gravity*, **35**, 095007, 2018. <http://iopscience.iop.org/article/10.1088/1361-6382/aab2e9/meta>.
- [241] Y. Aso, M. Ando, K. Kawabe, S. Otsuka, and K. Tsubono, “Stabilization of a Fabry–Perot interferometer using a suspension-point interferometer,” *Phys. Lett. A*, **327**, no. 1, pp. 1–8, 2004. <http://dx.doi.org/10.1016/j.physleta.2004.04.066>.
- [242] D. Clark, B. Lantz, and D. DeBra, “Seismic Platform Interferometer - Progress at Stanford,” *LIGO Document: LIGO-G1200178*, 2012. <https://dcc.ligo.org/G1200178>.
- [243] D. Clark, “Control of Differential Motion Between Adjacent Advanced LIGO Seismic Isolation Platforms,” Tech. Rep. P1300043, LIGO Laboratory, Mar 2013. <https://dcc.ligo.org/LIGO-P1300043>.
- [244] K. Dahl, A. Bertolini, M. Born, Y. Chen, D. Gering, S. Gößler, C. Gräf, G. Heinzl, S. Hild, F. Kawazoe, O. Kranz, G. Kühn, H. Lück, K. Mossavi, R. Schnabel, K. Somiya, K. A. Strain, J. R. Taylor, A. Wanner, T. Westphal, B. Willke, and K. Danzmann, “Towards a Suspension Platform Interferometer for the AEI 10 m Prototype Interferometer,” *J. Phys.: Conf. Ser.*, **228**, 012027, 2010. <http://dx.doi.org/10.1088/1742-6596/228/1/012027>.
- [245] C. Collette, F. Nassif, J. Amar, C. Depouhon, and S.-P. Gorza, “Prototype of interferometric absolute motion sensor,” *Sensors and Actuators A: Physical*, 2015.
- [246] J. vanHeijningen and J. van den Brand, “An interferometric readout for a monolithic accelerometer, towards a fm/rtHz resolution,” *LIGO Document: LIGO-G1300971*, 2013. <https://dcc.ligo.org/LIGO-G1300971>.
- [247] D. Martynov et al., “Sensitivity of the Advanced LIGO detectors at the beginning of gravitational wave astronomy,” *Physical Review D*, **93**, 112004, 2016. <https://journals.aps.org/prd/abstract/10.1103/PhysRevD.93.112004>.
- [248] Y. Hang et al., “Prospects for Detecting Gravitational Waves at 5 Hz with Ground-Based Detectors,” *Physical Review Letters*, **120**, 141102, 2018. <https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.120.141102>.
- [249] C. M. Mow-Lowry and D. Martynov, “A 6D interferometric inertial isolation system,” *arXiv:1801.01468 [astro-ph.IM]*, 2018. <https://arxiv.org/abs/1801.01468>.
- [250] A. Cumming, R. Jones, M. Barton, G. Cagnoli, C. A. Cantley, D. R. M. Crooks, G. D. Hammond, A. Heptonstall, J. Hough, S. Rowan, and K. A. Strain, “Apparatus for dimensional characterization of fused silica fibers for the suspensions of advanced gravitational wave detectors,” *Rev. Sci. Instrum.*, **82**, 044502, 2011. <http://link.aip.org/link/doi/10.1063/1.3581228>.

- [251] A. Heptonstall, M. Barton, C. Cantley, A. Cumming, G. Cagnoli, J. Hough, R. Jones, R. Kumar, I. Martin, S. Rowan, C. Torrie, and S. Zech, “Investigation of mechanical dissipation in CO2 laser-drawn fused silica fibres and welds,” *Class. Quantum Grav.*, **27**, 035013, 2010. <http://dx.doi.org/10.1088/0264-9381/27/3/035013>.
- [252] A. Heptonstall, M. A. Barton, A. Bell, G. Cagnoli, C. A. Cantley, and et al., “Invited Article: CO2 laser production of fused silica fibers for use in interferometric gravitational wave detector mirror suspensions,” *Rev. Sci. Instrum.*, **82**, 011301, 2011. <http://link.aip.org/link/doi/10.1063/1.3532770>.
- [253] L. Cunningham, P. Murray, A. Cumming, E. Elliffe, G. Hammond, K. Haughian, J. Hough, M. Hendry, R. Jones, I. Martin, S. Reid, S. Rowan, J. Scott, K. Strain, K. Tokmakov, C. Torrie, and A. van Veggel, “Re-evaluation of the mechanical loss factor of hydroxide-catalysis bonds and its significance for the next generation of gravitational wave detectors,” *Phys. Lett. A*, **374**, no. 39, pp. 3993 – 3998, 2010. <http://dx.doi.org/10.1016/j.physleta.2010.07.049>.
- [254] G. D. Hammond et al., “Suspension Activities in Glasgow,” *LIGO G1800423*, 2018. <https://dcc.ligo.org/LIGO-G1800423>.
- [255] B. Sorazu et al., “Characterisation of the aLIGO monolithic suspensions,” *LIGO G1601163*, 2016. <https://dcc.ligo.org/LIGO-G1601163>.
- [256] B. Sorazu, K. A. Strain, I. S. Heng, , and R. Kumar, “Violin mode amplitude glitch monitor for the presence of excess noise on the monolithic silica suspensions of GEO 600,” *Class. Quantum Grav.*, **27**, 155017, 2010. <http://dx.doi.org/10.1088/0264-9381/27/15/155017>.
- [257] X. Ni, E. Quintero, and G. Vajente, “Proposal for an Upgrade of the Crackle Experiment,” *LIGO Document: LIGO-T1400407*, 2014. <https://dcc.ligo.org/LIGO-T1400407>.
- [258] G. Vajente, E. A. Quintero, X. Ni, K. Arai, E. K. Gustafson, N. A. Robertson, E. J. Sanchez, J. R. Greer, and R. X. Adhikari, “An instrument to measure mechanical up-conversion phenomena in metals in the elastic regime,” *Review of Scientific Instruments*, **87**, no. 6, 2016. <http://scitation.aip.org/content/aip/journal/rsi/87/6/10.1063/1.4953114>.
- [259] G Vajente, “Crackling noise in advanced gravitational wave detectors: A model of the steel cantilevers used in the test mass suspensions,” *Phys. Rev D*, **96**, 022003, 2017. <https://journals.aps.org/prd/abstract/10.1103/PhysRevD.96.022003>.
- [260] G Vajente, “Crackling noise scaling: from the Crackle2 experiment to the QUAD suspension,” *LIGO document G1700347*, 2017. <https://dcc.ligo.org/LIGO-G1700347>.
- [261] N. Robertson and C. Torrie, “Redesign of Blades for Improved Vertical Isolation for Triple Suspensions in Advanced LIGO,” *LIGO Document: LIGO-T080267*, 2008. <https://dcc.ligo.org/LIGO-T080267>.

- [262] M. Barsanti, M. Beghini, F. Frasconi, R. Ishak, B. D. Monelli, R. Valentini, “Experimental study of hydrogen embrittlement in Maraging steels,” *Procedia Structural Integrity*, 2018. <https://www.sciencedirect.com/science/article/pii/S2452321617305425>.
- [263] Y. Ikushima, R. Li, T. Tomaru, N. Sato, T. Suzuki, T. Haruyama, T. Shintomi, and A. Yamamoto, “Ultra-low-vibration pulse-tube cryocooler system - cooling capacity and vibration,” *Cryogenics*, **48**, no. 9-10, pp. 406 – 412, 2008. <http://www.sciencedirect.com/science/article/pii/S0011227508000490>.
- [264] B. Shapiro, “Cryogenic Test Mass Work at Stanford,” *LIGO Document: LIGO-G1400926*, 2014. <https://dcc.ligo.org/G1400926>.
- [265] O. Aguiar, “Multi-Nested Pendula System: measured, calculated and simulated cooling performances,” *LIGO Document: LIGO-G1400845*, 2014. <https://dcc.ligo.org/G1400845>.
- [266] O. Aguiar, M. Constancio, E. Ferreira, A. Silva, and M. Okada, “Multi-Nested Pendula System: Mechanics and Cryogenics Update,” *LIGO Document: LIGO-G1500172*, 2015. <https://dcc.ligo.org/G1500172>.
- [267] B. Shapiro, “A Cryogenic LIGO Mirror for 3rd Generation Observatories,” *LIGO Document: LIGO-G1600766*, 2016. <https://dcc.ligo.org/G1600766>.
- [268] B. Shapiro, “LIGO Voyager Cryogenics at Stanford,” 2017. LIGO-G1700404, <https://dcc.ligo.org/LIGO-G1700404/public>.
- [269] D. I. G. Jones, *Handbook of Viscoelastic Vibration Damping*. Chichester: J. Wiley, 2001. pp. 171-174.
- [270] Y. Sakakibara et al., “A study of cooling time reduction of interferometric cryogenic gravitational wave detectors using a high-emissivity coating,” *AIP Conf. Proc.*, **1573**, 176, 2013. <https://arxiv.org/abs/1309.4836>.
- [271] Acktar Ltd., “Acktar black coating services,” <http://acktar.com/category/BlackOpticalCoating>.
- [272] J. Lehman, C. Yung, N. Tomlin, D. Conklin, and M. Stephens, “Carbon nanotube-based black coatings,” *Applied physics reviews*, **5**, 011103, 2018. <https://aip.scitation.org/doi/pdf/10.1063/1.5009190>.
- [273] M R Abernathy, N Smith, W Z Korth, R X Adhikari, L G Prokhorov, D V Koptsov and V P Mitrofanov, “Measurement of mechanical loss in the Acktar Black coating of silicon wafers,” *Classical and Quantum Gravity*, **33**, 185002. <http://iopscience.iop.org/article/10.1088/0264-9381/33/18/185002/pdf>.
- [274] L.G. Prokhorov, V.P. Mitrofanov, B. Kamai, A. Markowitz, Xiaoyue Ni, R. X. Adhikari, “Measurement of mechanical losses in the carbon nanotube NASA black coating of silicon wafers,” *LIGO-G1800447*, 2018. <https://dcc.ligo.org/LIGO-G1800447>.

- [275] G. Hammond, “Suspension Thermal Noise,” *LIGO Document: LIGO-G1200579*, 2012. <https://dcc.ligo.org/G1200579>.
- [276] R. Birney, A. V. Cumming, P. Campsie, D. Gibson, G. D. Hammond, J. Hough, I. W. Martin, S. Reid, S. Rowan, S. Song, “Coatings and surface treatments for enhanced performance suspensions for future gravitational wave detectors,” *Classical and Quantum Gravity*, **34**, 235012, 2017. <http://iopscience.iop.org/article/10.1088/1361-6382/aa9354/meta>.
- [277] L.G. Prokhorov, D.V. Koptsov, M.S. Matiushechkina, V.P. Mitrofanov, K. Haughian, J. Hough, S. Rowan, A.A. van Veggel, P.G. Murray, G.D. Hammond, K. Tokmakov, “Cryogenic measurement of hydroxide catalysis bond loss using a silicon tuning fork oscillator,” *LIGO-G1700340*, 2017. <https://dcc.ligo.org/LIGO-G1700340>.
- [278] B. Sathyaprakash et al., “Scientific objectives of Einstein Telescope,” *Classical and Quantum Gravity*, **29**, 124013, 2012. <http://iopscience.iop.org/article/10.1088/0264-9381/29/12/124013/meta>.
- [279] B.P. Abbott et al., “Exploring the sensitivity of next generation gravitational wave detectors,” *Classical and Quantum Gravity*, **34**, 044001, 2017. <http://iopscience.iop.org/article/10.1088/1361-6382/aa51f4/pdf>.
- [280] V. Guerrero, “Monitoring Void Fraction and Bubble Formation in Flowing Liquid Nitrogen,” *LIGO G1701534*, 2017. <https://dcc.ligo.org/LIGO-G1701534>.
- [281] O. Aguiar, “Cryogenic Updates for Voyager,” *LIGO G1800357*, 2018. <https://dcc.ligo.org/LIGO-G1800357>.
- [282] L. Winkelmann, O. Puncken, R. Kluzik, C. Veltkamp, P. Kwee, J. Poeld, C. Bogan, B. Willke, M. Frede, J. Neumann, P. Wessels, and D. Kracht, “Injection-locked single-frequency laser with an output power of 220 W,” *Applied Physics B: Lasers and Optics*, **102**, pp. 529–538, 2011. <http://www.springerlink.com/content/t2r71860p9631681/>.
- [283] M. Frede, B. Schulz, R. Wilhelm, P. Kwee, F. Seifert, B. Willke, and D. Kracht, “Fundamental mode, single-frequency laser amplifier for gravitational wave detectors,” *Opt. Express*, **15**, pp. 459–465, 2007. <http://www.opticsinfobase.org/oe/abstract.cfm?URI=oe-15-2-459>.
- [284] M. Frede, R. Wilhelm, D. Kracht, and C. Fallnich, “Nd:YAG ring laser with 213 W linearly polarized fundamental mode output power,” *Opt. Express*, **13**, pp. 7516–7519, 2005. <http://www.opticsinfobase.org/oe/abstract.cfm?URI=oe-13-19-7516>.
- [285] P. Kwee, B. Willke, and K. Danzmann, “Shot-noise-limited laser power stabilization with a high-power photodiode array,” *Opt. Lett.*, **34**, pp. 2912–2914, 2009. <http://www.opticsinfobase.org/abstract.cfm?URI=ol-34-19-2912>.
- [286] D. A. Bryan, R. Gerson, and H. E. Tomaschke, “Increase optical damage resistance in lithium niobate,” *Appl. Phys. Lett.*, **44**, 847, 1984. <http://dx.doi.org/10.1063/1.94946>.

- [287] M. Ganija, D. Ottaway, P. Veitch, and J. Munch, “Cryogenic, high power, near diffraction limited, Yb:YAG slab laser,” *Opt. Express*, **21**, pp. 6973–6978, Mar 2013. <http://www.opticsexpress.org/abstract.cfm?URI=oe-21-6-6973>.
- [288] M. Evans and P. Fritschel, “TCS and the Golden Shield,” *LIGO Document: LIGO-T0900359-v2*, 2009. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=3967>.
- [289] G. Mueller, M. Arain, P. Sainathan, and G. Ciani, “aLIGO TCS Ring Heater development at UF - Krakow 2010 talk,” *LIGO Document: LIGO-G1000945-v1*, 2010. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=21294>.
- [290] R. Abbott, R. Adhikari, S. Ballmer, L. Barsotti, M. Evans, P. Fritschel, V. Frolov, G. Mueller, B. Slagmolen, and S. Waldman, “AdvLIGO Interferometer Sensing and Control Conceptual Design,” *LIGO Document: LIGO-T070247-01*, 2007. <https://dcc.ligo.org/cgi-bin/DocDB/ShowDocument?docid=5306>.
- [291] K. Arai, M. Ando, S. Moriwaki, K. Kawabe, and K. Tsubono, “New signal extraction scheme with harmonic demodulation for power-recycled Fabry-Perot-Michelson interferometers,” *Phys. Lett. A*, **273**, pp. 15–24, 2000. [http://dx.doi.org/10.1016/S0375-9601\(00\)00467-9](http://dx.doi.org/10.1016/S0375-9601(00)00467-9).
- [292] T. Fricke, N. Smith-Lefebvre, R. Abbott, R. Adhikari, K. L. Dooley, M. Evans, P. Fritschel, V. Frolov, K. Kawabe, J. S. Kissel, and S. Waldman, “DC readout experiment in Enhanced LIGO,” *Class. Quantum Grav.*, 2012. <http://dx.doi.org/10.1088/0264-9381/29/6/065005>.
- [293] S. Hild, H. Grote, J. Degallaix, S. Chelkowski, K. Danzmann, A. Freise, M. Hewitson, J. Hough, H. Lück, M. Prijatelj, K. A. Strain, J. R. Smith, and B. Willke, “DC-readout of a signal-recycled gravitational wave detector,” *Class. Quantum Grav.*, **26**, no. 5, 055012, 2009. <http://dx.doi.org/10.1088/0264-9381/26/5/055012>.
- [294] K. Goda, O. Miyakawa, E. E. Mikhailov, S. Saraf, R. Adhikari, K. McKenzie, R. Ward, S. Vass, A. J. Weinstein, and N. Mavalvala, “A quantum-enhanced prototype gravitational-wave detector,” *Nature Physics*, **4**, pp. 472–476, 2008. <http://dx.doi.org/10.1038/nphys920>.
- [295] O. Miyakawa, R. Ward, R. Adhikari, M. Evans, B. Abbott, R. Bork, D. Busby, J. Heefner, A. Ivanov, M. Smith, R. Taylor, S. Vass, A. Weinstein, M. Varvella, S. Kawamura, F. Kawazoe, S. Sakata, and C. Mow-Lowry, “Measurement of optical response of a detuned resonant sideband extraction gravitational wave detector,” *Phys. Rev. D*, **74**, 022001, 2006. <http://link.aps.org/doi/10.1103/PhysRevD.74.022001>.
- [296] R. L. Ward, R. Adhikari, B. Abbott, R. Abbott, D. Barron, R. Bork, T. Fricke, V. Frolov, J. Heefner, A. Ivanov, O. Miyakawa, K. McKenzie, B. Slagmolen, M. Smith, R. Taylor, S. Vass, S. Waldman, and A. Weinstein, “DC readout experiment at the Caltech 40m prototype interferometer,” *Class. Quantum Grav.*, **25**, no. 11, 114030, 2008. <http://dx.doi.org/10.1088/0264-9381/25/11/114030>.



- [297] M. Evans, P. Fritschel, D. McClelland, J. Miller, A. Mullavey, D. Shaddock, B. Slagmolen, S. Waldman, and et al., “Adv. LIGO Arm Length Stabilization Design,” *LIGO Document: LIGO-T0900144*, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=1625>.
- [298] A. J. Mullavey, B. J. J. Slagmolen, J. Miller, M. Evans, P. Fritschel, D. Sigg, S. J. Waldman, D. A. Shaddock, and D. E. McClelland, “Arm-length stabilisation for interferometric gravitational-wave detectors using frequency-doubled auxiliary lasers,” *Opt. Express*, **20**, pp. 81–89, 2012. <http://dx.doi.org/10.1364/OE.20.000081>.
- [299] K. Izumi, K. Arai, B. Barr, J. Betzwieser, A. Brooks, K. Dahl, S. Doravari, J. C. Driggers, W. Z. Korth, H. Miao, J. G. Rollins, S. Vass, D. Yeaton-Massey, and R. Adhikari, “Multi-color Cavity Metrology,” *LIGO Document: LIGO-P1200019*, 2012. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=87878>.
- [300] J. A. Sidles and D. Sigg, “Optical torques in suspended Fabry–Perot interferometers,” *Phys. Lett. A*, **354**, pp. 167–172, 2006. <http://dx.doi.org/10.1016/j.physleta.2006.01.051>.
- [301] N. Smith-Lefebvre, S. Ballmer, M. Evans, S. Waldman, K. Kawabe, V. Frolov, and N. Mavalvala, “Optimal alignment sensing of a readout mode cleaner cavity,” *Opt. Lett.*, **36**, no. 22, pp. 4365–4367, 2011. <http://dx.doi.org/10.1364/OL.36.004365>.
- [302] J. Lissel, “Absolute Calibration of a Galactic Gravitational Wave Detector Network,” 2018. LIGO-G1800953, <https://dcc.ligo.org/LIGO-G1800953>.
- [303] S. Karki, D. Tuyenbayev, S. Kandhasamy, B. Abbott, T. Abbott, E. Anders, J. Berliner, J. Betzwieser, C. Cahillane, L. Canete, C. Conley, H. Daveloza, N. De Lillo, J. Gleason, E. Goetz, K. Izumi, J. Kissel, G. Mendell, V. Quetschke, M. Rodruck, S. Sachdev, T. Sadecki, P. Schwinberg, A. Sottile, M. Wade, A. Weinstein, M. West, and R. Savage, “The Advanced LIGO photon calibrators,” *Review of Scientific Instruments*, **87**, 114503, 2016. (LIGO-P1500249), <https://aip.scitation.org/doi/10.1063/1.4967303>.
- [304] C. Cahillane<sup>1</sup>, J. Betzwieser, D. Brown, E. Goetz, E. Hall, K. Izumia, S. Kandhasamy, S. Karki, J. Kissel, G. Mendell, R. Savage, D. Tuyenbayev, A. Urban, A. Viets, M. Wade, and A. Weinstein, “Calibration Uncertainty for Advanced LIGOs First and Second Observing Runs,” *Physical Review D*, **96**, 102001, 2017. (LIGO-P1600139), <https://journals.aps.org/prd/abstract/10.1103/PhysRevD.96.102001>.
- [305] D. Tuyenbayev, T. Sadecki, and R. Savage, “Photon Calibrator Procedure for Measuring Response Ratios of Power Calibration Standards,” 2018. (LIGO-T1400442), <https://dcc.ligo.org/LIGO-T1400442>.
- [306] Y. Inoue, S. Haino, N. Kanda, Y. Ogawa, T. Suzuki, T. Tomaru, T. Yamamoto, and T. Yokozawa, “Improving the absolute accuracy of the gravitational wave detectors by combining the photon pressure and gravity field calibrators,” *arXiv preprint arXiv:1804.08249*, 2018. <https://arxiv.org/abs/1804.08249>.

- [307] B. Mours, “First tests of a Newtonian Calibrator (NCal),” *LIGO Document Control Center*, 2018. LIGO-G1800442, <https://dcc.ligo.org/LIGO-G1800442>.
- [308] M. Rakhmanova, R. Savage, D. Reitze, and D. Tanner, “Dynamic resonance of light in FabryPerot cavities,” *Physics Letters A*, **305**, pp. 239–244, 2002. LIGO-P010013, [https://doi.org/10.1016/S0375-9601\(02\)01469-X](https://doi.org/10.1016/S0375-9601(02)01469-X).
- [309] J. C. Driggers, M. Evans, K. Pepper, and R. Adhikari, “Active noise cancellation in a suspended interferometer,” *Rev. Sci. Instrum.*, **83**, 2012. <http://dx.doi.org/10.1063/1.3675891>.
- [310] G. Vajente, “Machine Learning and Controls in GW detectors,” 2018. LIGO-G1800874, <https://dcc.ligo.org/LIGO-G1800874>.
- [311] K. Arai, J. Betzwieser, P. Kalmus, and R. Adhikari, “Simulated Plant Approach,” *LIGO Document: LIGO-G1000546*, 2010. <https://dcc.ligo.org/cgi-bin/private/DocDB/ShowDocument?docid=11804>.
- [312] N. Mukund, “update on lockloss monitoring & prediction,” 2018. LIGO-G1800569, <https://dcc.ligo.org/LIGO-G1800569>.
- [313] L. Price, “Optimal Transfer Functions,” 2014. LIGO-G1400084, <https://dcc.ligo.org/LIGO-G1400084>.
- [314] R. Pintelon and J. Schoukens, *System Identification: A Frequency Domain Approach*. John Wiley & Sons, 2012.
- [315] M. Coughlin, P. Earle, J. Harms, S. Biscansa, C. Buchanan, E. Coughlin, F. Donovan, J. Fee, H. Gabbard, M. Guy, and *et al.*, “Limiting the effects of earthquakes on gravitational-wave interferometers,” *Classical and Quantum Gravity*, 044004, 2017. LIGO-P1600321, <https://arxiv.org/abs/1611.09812v2>.
- [316] N. Mukund, M. Coughlin, J. Harms, N. Arnaud, D. Barker, S. Biscans, E. Coughlin, F. Donovan, P. Earle, J. Fee, I. Fiori, H. Gabbard, M. Guy, B. Lantz, R. Mittleman, A. Pele, H. Radkins, B. Swinkels, K. Thorne, and J. Warner, “Status of Seismon,” 2018. LIGO-G1800118, <https://dcc.ligo.org/LIGO-G1800118>.
- [317] S. Biscans, J. Warner, R. Mittleman, C. Buchanan, M. Coughlin, M. Evans, H. Gabbard, J. Harms, B. Lantz, N. Mukund, A. Pele, C. Pezerat, P. Picart, H. Radkins, and T. Shaffer, “Control strategy to limit duty cycle impact of earthquakes on the LIGO gravitational-wave detectors,” *Classical and Quantum Gravity*, **35**, no. 5, 055004, 2018. LIGO-P1700163, <http://stacks.iop.org/0264-9381/35/i=5/a=055004>.
- [318] A. Pele, M. Kasprzack, R. Adhikari, H. Yu, E. Quintero, and J. Driggers, “L2A Decoupling webpage,” *Classical and Quantum Gravity*. [https://wiki.ligo.org/CSWG/L2A\\_Decoupling/WebHome](https://wiki.ligo.org/CSWG/L2A_Decoupling/WebHome).
- [319] B. Lantz, S. Biscans, J. Warner, A. Pele, H. Radkins, N. Mukund, and M. Coughlin, “Seismic Control during Earthquakes: Review of new technique,” LIGO-G1800399, <https://dcc.ligo.org/LIGO-G1800399>.

- [320] M. Kasprzack, “Angular control issues in Advanced LIGO,” 2018. LIGO-G1800989, <https://dcc.ligo.org/LIGO-G1800989>.
- [321] H. Yu and D. Martynov, “Optimal and robust control of the arm alignment with mu-synthesis,” 2018. LIGO-T1800077, <https://dcc.ligo.org/LIGO-T1800077>.
- [322] H. Yu and D. Sigg, “Alignment sensing in the signal recycling cavity using a new 118.3 MHz sideband scheme,” 2017. LIGO-T1700215, <https://dcc.ligo.org/LIGO-T1700215>.