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## Overview of topics

#### technical challenges for Cosmic Explorer

Selected

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#### Overview

Signal recycling cavity

Frequency control

Test mass size and figure

References

- Signal recycling cavity design
- Frequency control
- Test mass size and figure
  - Clipping loss
  - Radius of curvature tolerance
  - Scattering from surface figure errors

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Overview

#### Signal recycling cavity

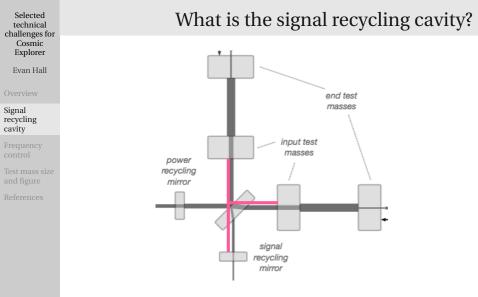
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## Signal recycling cavity design

### How should we choose the transmissivity of the signal recycling mirror and its distance from the arms?



Once the arm bandwidth is fixed (by fixing the ITM transmission), the bandwidth of the detector's GW response is set by choosing the SRM transmission.

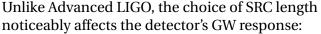
### Effect on GW response

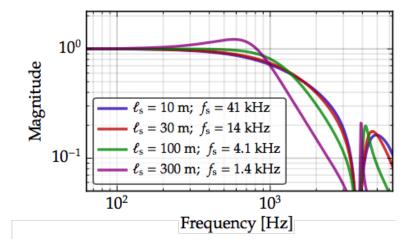
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Signal recycling cavity





### Frequency control

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How should we control the laser so as to keep light resonant in the interferometer and suppress laser frequency fluctuations?

### Frequency control

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Advanced LIGO: laser frequency is locked to the average arm length using Pound–Drever–Hall (PDH) reflection locking.

Servo loop needs a bandwidth > 10 kHz in order to adequately suppress frequency noise below 1 kHz.

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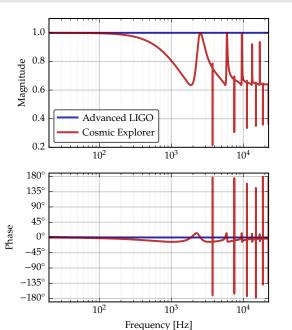
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## Frequency control for CE?

Presence of right-handed zeros above a few kilohertz limits the bandwidth of a PDH frequency-locking loop

Could blend the PDH signal with another signal (perhaps a single non-resonant sideband)

## Test mass size and figure

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How large should the test masses be, and how finely must we control their radius of curvature and surface figure?

### Mirrors for CE

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Compared to Advanced LIGO, Cosmic Explorer will have mirrors that...

- are wider (diameter  $\sim 30 \text{ cm} \rightarrow ??$ )
- are **flatter** (radius of curvature  $\sim 2 \text{ km} \rightarrow \sim 30 \text{ km}$ )

• must accommodate larger laser beams (spot radius  $\sim 5 \text{ cm} \rightarrow \sim 15 \text{ cm}$ )

• may be made of a **different material** (silica  $\rightarrow$  silicon) What challenges will we face as we fabricate these new mirrors?

**Clipping loss** 

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Clipping loss: outermost ring of the laser beam falls off the test mass.

Clipping loss reduces the power stored in the interferometer arms (and must be compensated with higher input laser power).

Simple expression:  $\eta_{clip} = \exp[-2(r_{TM}/w)^2]$ , where  $r_{TM}$  is the radius of the test mass, and *w* is the beam spot radius.

For Advanced LIGO:  $r_{\rm TM} = 17$  cm and  $w_{\rm ETM} = 6.2$  cm, giving  $\eta_{\rm clip} < 1$  ppm.

For CE, 15 ppm might be acceptable (Dwyer et al. [1]).

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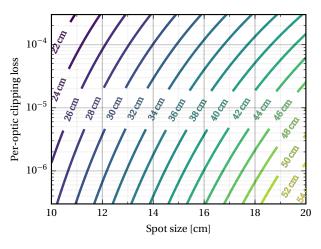
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### For Cosmic Explorer:



Test mass radius > 30 cm required to achieve per-optic clipping loss below 15 ppm...

# Clipping loss

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### Radius of curvature tolerance

RoC mismatch between arms will cause junk light ("contrast defect") at the interferometer's GW readout.

For Advanced LIGO, the manufacturing tolerance on the **sagitta** is  $\Delta s \simeq 5$  nm.

Assuming a similar tolerance for CE, if  $\Re = 30$  km, the resulting RoC error is  $\Delta \Re \simeq 100$  m.

## Radius of curvature tolerance

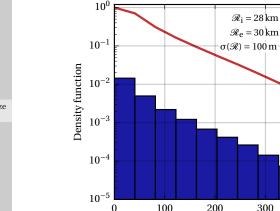
CDF

PDF

400

500

Numerical simulation of mode mismatch:



Contrast defect < 250 ppm for 95 % of realizations — too large by at least a factor of 10. Active modematching will be required.

Mode mismach [ppm]

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## Surface figure

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**Surface figure:** spatial fluctuations across the optic surface *other than* the radius of curvature

The figure of an isotropic surface is characterized by a **spatial power spectrum**  $S_1(k)$ , where k is the spatial wavenumber (units of m<sup>-1</sup>).

Given a particular PSD  $S_1(k)$ , a larger laser beam will experience more scatter because it samples longer wavelengths.

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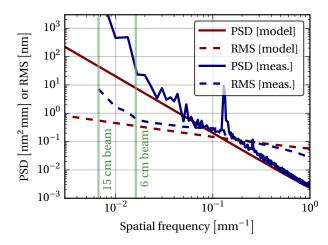
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## Advanced LIGO mirror power spectrum

Advanced LIGO model after Yamamoto [2]:  $S_1(k) = (0.2 \,\mathrm{nm}^2 \,\mathrm{mm}) \times \left(\frac{0.1 \,\mathrm{mm}^{-1}}{k}\right)^2$ 



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Formula for power scattered out of a beam with spot size w [3]:

$$\eta_{\text{scat}}(w) = \left[\frac{4\pi\sigma_{\text{scat}}(w)}{\lambda_0}\right]^2 = \int_0^\infty dk \, S_1(k) \, W_w(k). \tag{1}$$

 $W_w(k)$  suppresses spatial modes larger than the spot size;  $\sigma_{scat}(w)$  is the effective total rms surface roughness for the given spot size.

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Note the scaling: if  $S_1(k) \propto 1/k^2$ , then  $\eta \propto \sigma^2 \propto w$ .

For Advanced LIGO:  $\sigma(6.2 \text{ cm}) = 0.35 \text{ nm}$ , so  $\eta_{scat} = 20 \text{ ppm}$  per optic

For Cosmic Explorer:  $\sigma(15 \text{ cm}) = 0.55 \text{ nm}$ , so  $\eta_{scat} = 48 \text{ ppm}$  per optic

Note the caveats: behavior of  $S_1(k)$  for  $k < 0.01 \text{ mm}^{-1}$  not well known; optical cavity effects not yet included.

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S. Dwyer, D. Sigg, S. W. Ballmer, L. Barsotti, N. Mavalvala, and M. Evans. Gravitational wave detector with cosmological reach. *Phys. Rev. D* **91**(8) (2015), p. 082001.

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W. Winkler, R. Schilling, K. Danzmann, J. Mizuno, A. Rüdiger, and K. A. Strain. Light scattering described in the mode picture. *Appl. Opt.* **33**(31) (1994), pp. 7547–7550.