



# A Sound of Thunder

## Atmospheric Gravitational Noise

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- A. Infrasonic Newtonian Noise
- B. Thermal Newtonian Noise
- C. Transients

## A. Infrasonic Newtonian Noise

$$\bullet \quad \sqrt{S_h} \sim 4\pi G \rho \frac{\sqrt{S_p}}{\gamma p} \frac{1}{\omega^2} \times \begin{cases} v/\omega L & \text{if } \omega \gtrsim v/L \\ Q^2(v/\omega d)^3 & \text{if } \omega \gtrsim Qv/d \end{cases}$$

where  $v \approx 300 \text{ m/s}$ ,  $L = \text{arm length}$ ,  $d = \text{exclusion distance}$ .

- For  $f \sim 10 \text{ Hz}$ ,  $\lambda \sim 30 \text{ m}$ , spatial correlations  $\sim \text{several } \lambda$ .

## A. Infrasonic Newtonian Noise

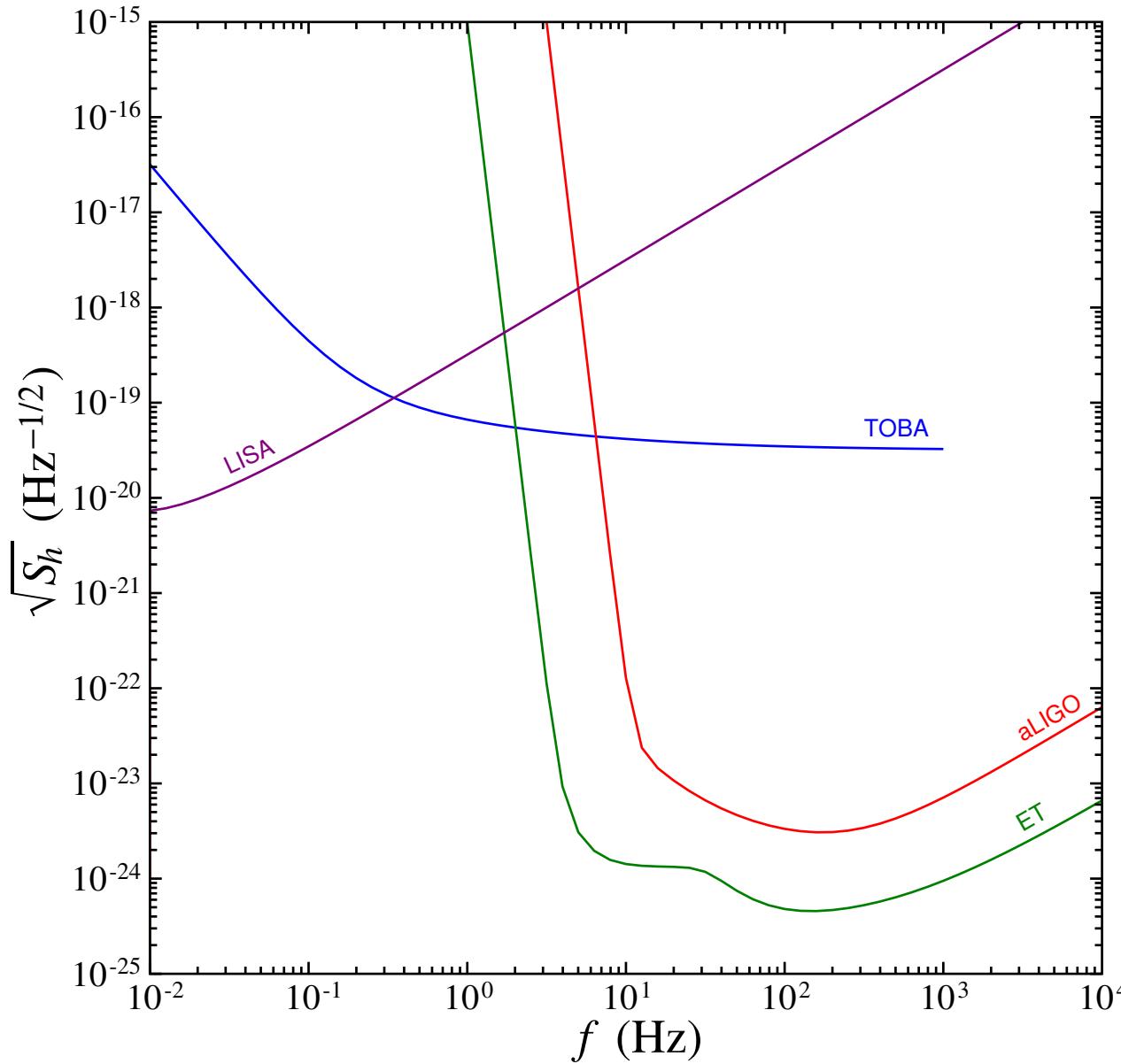
- $\sqrt{S_h} \sim 4\pi G \rho \frac{\sqrt{S_p}}{\gamma p} \frac{1}{\omega^2} \times \begin{cases} v/\omega L & \text{if } \omega \gtrsim v/L \\ Q^2(v/\omega d)^3 & \text{if } \omega \gtrsim Qv/d \end{cases}$

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- Disturbance typically dominated by gravity difference from nearest wave crest and trough.

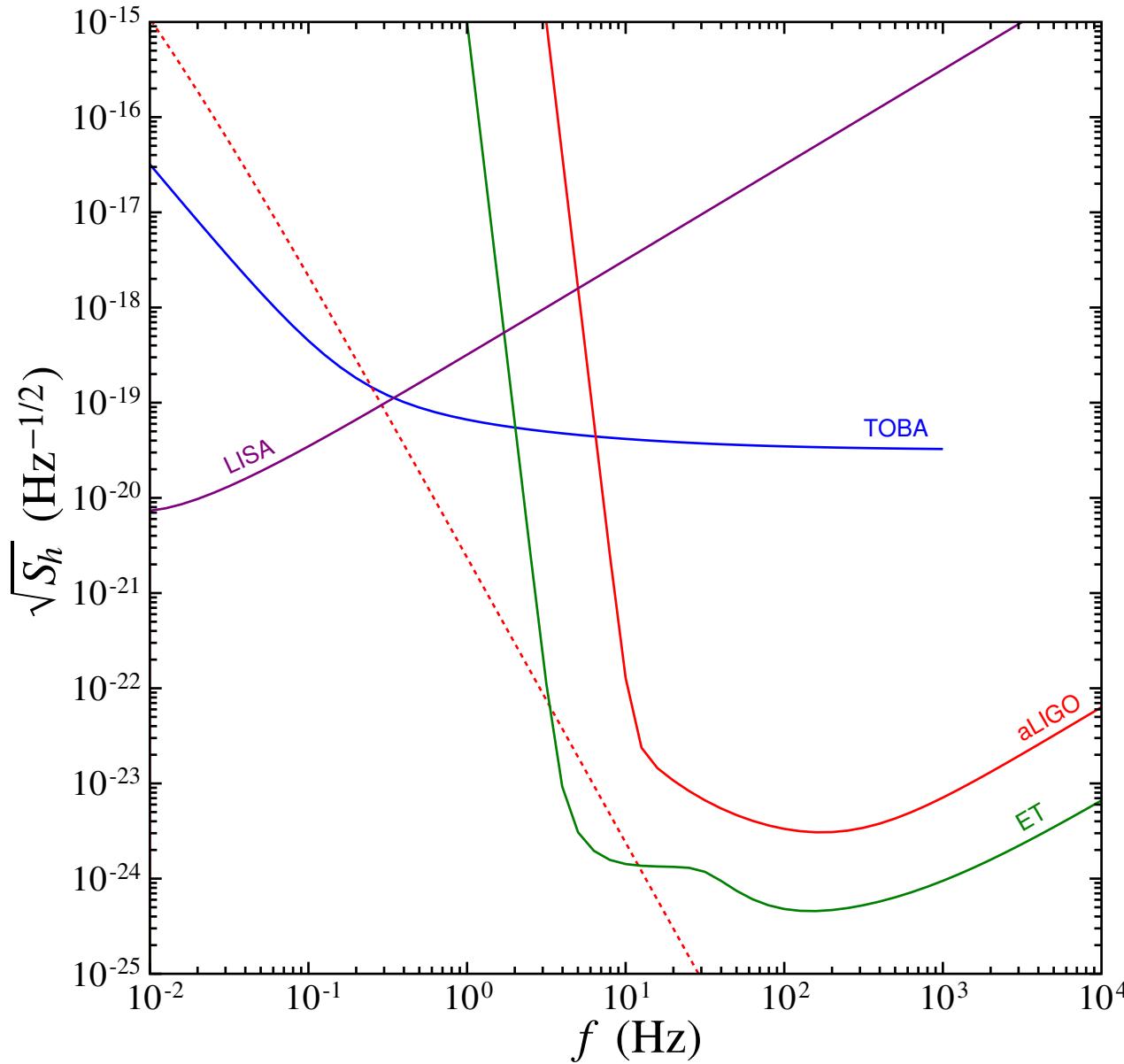
$$\begin{aligned} \sqrt{S_h} &\approx \left(2 \times 10^{-21} \text{ Hz}^{-1/2}\right) \left(\frac{\sqrt{S_p}}{3 \text{ mPa Hz}^{-1/2}}\right) \left(\frac{f}{\text{Hz}}\right)^{-3} && (\text{4km}) \\ &\approx \left(2 \times 10^{-19} \text{ Hz}^{-1/2}\right) \left(\frac{\sqrt{S_p}}{3 \text{ mPa Hz}^{-1/2}}\right) \left(\frac{f}{\text{Hz}}\right)^{-2} && (\text{gradiometer}) \end{aligned}$$

## A. Infrasonic Newtonian Noise

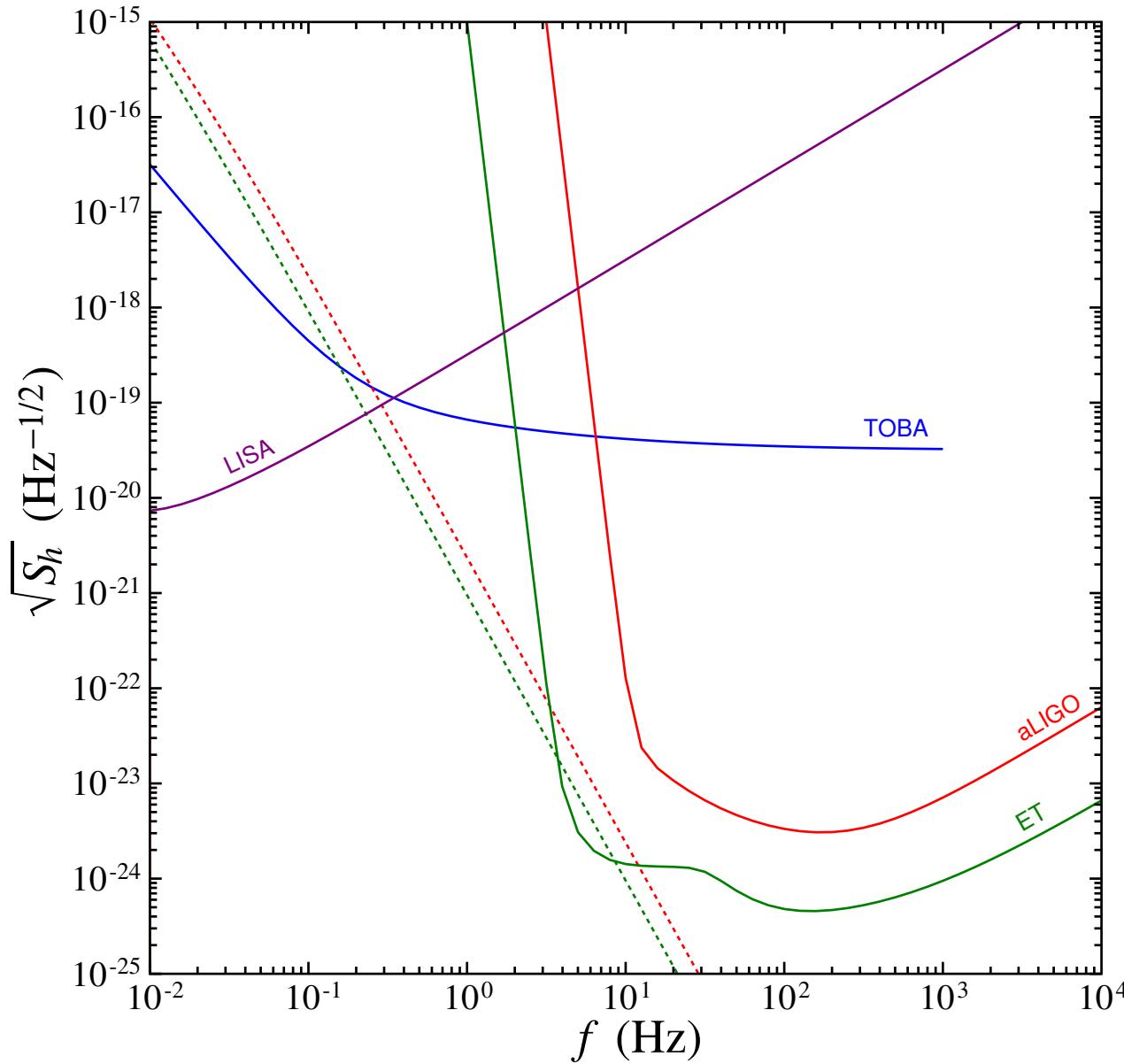


Unshielded NN

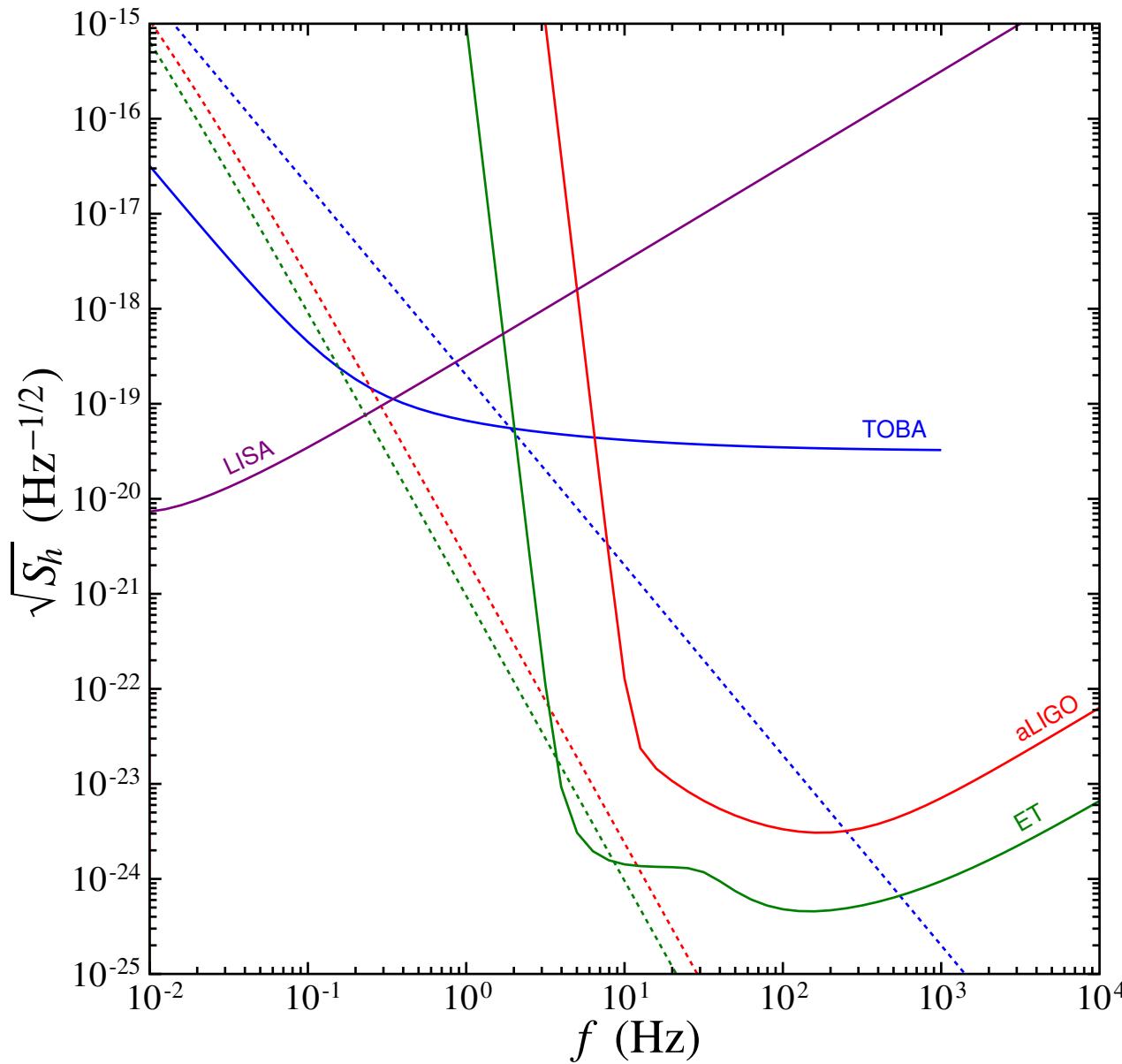
## A. Infrasonic Newtonian Noise



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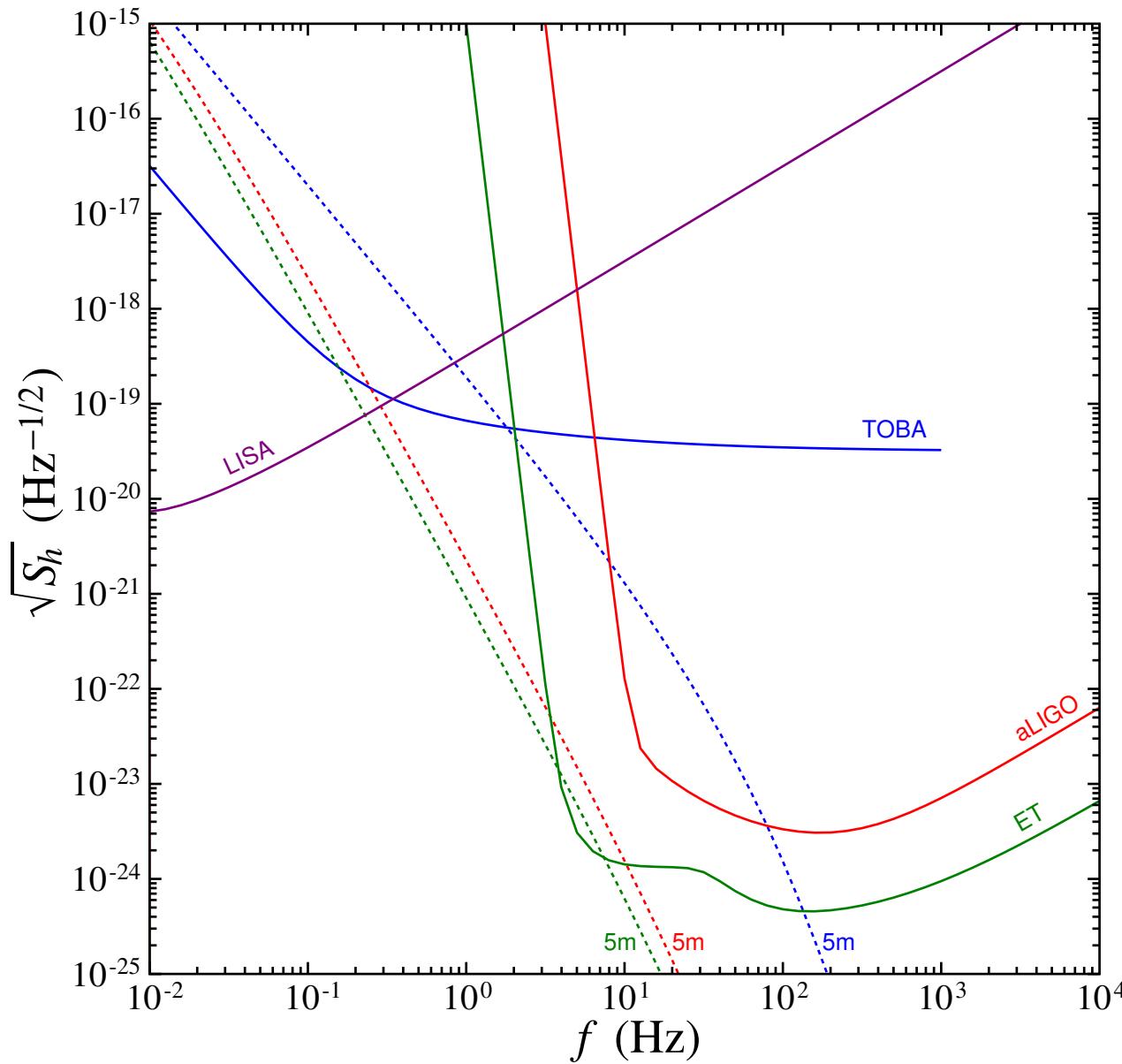


## A. Infrasonic Newtonian Noise



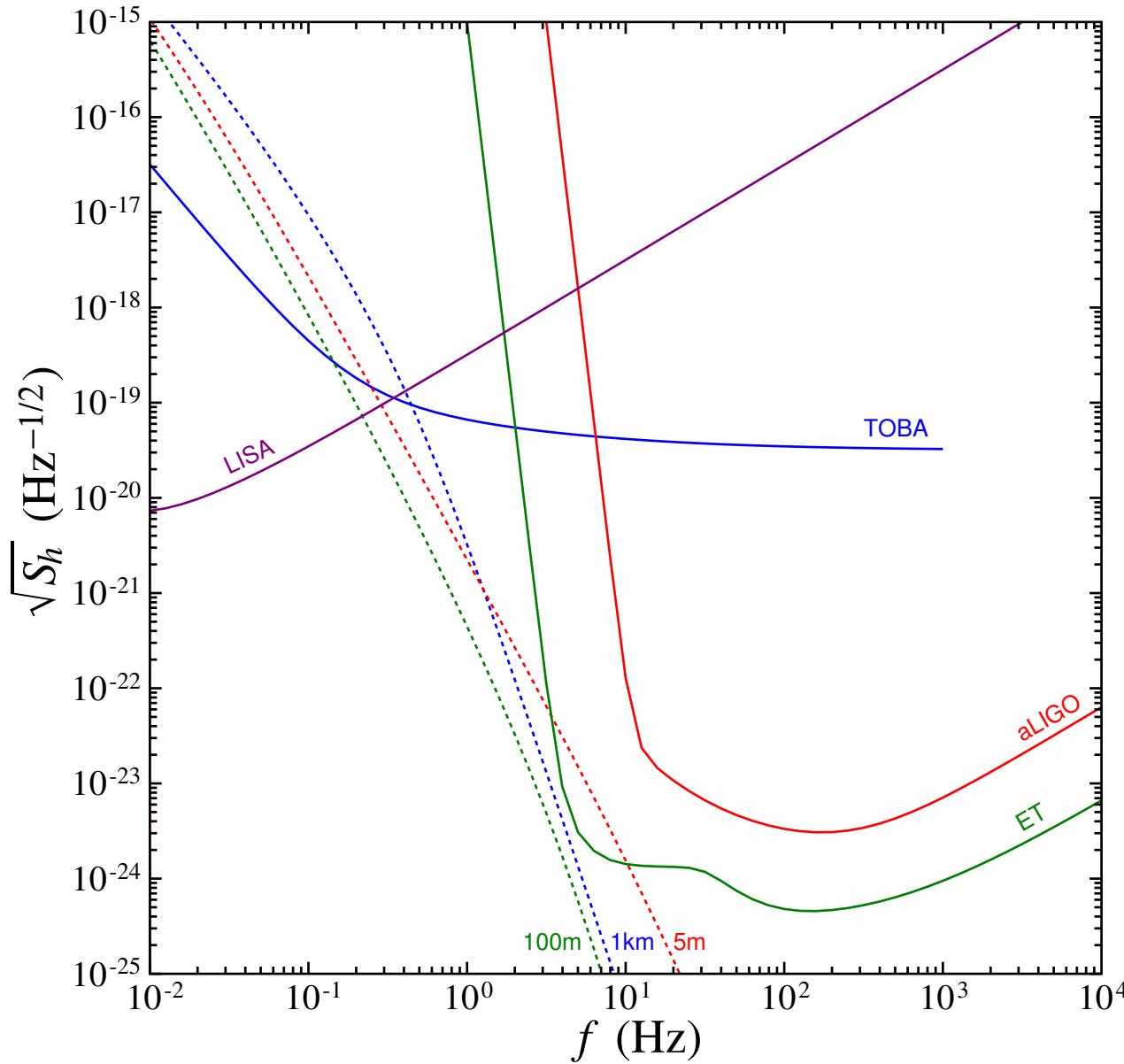
Unshielded NN  
4km arms  
10km arms  
gradiometer

## A. Infrasonic Newtonian Noise



Surface structures  
4km arms  
10km arms  
gradiometer

## A. Infrasonic Newtonian Noise



Underground  
4km arms  
10km arms  
gradiometer

## A. Infrasonic Newtonian Noise

Considerations ( $f \gtrsim$  few Hz):

- Surface: About  $0.1 \times$  amplitude of seismic NN.
- Underground: further damped by  $\lambda/d$ ,  $l_{\text{coh}}^2/d^2$ .
- (If infrasound generated by ground motions, or vice-versa, seismic NN will dominate by factor  $\propto \rho^{3/4} C^{-1/4} \sim 30$  in  $\sqrt{S_h}$ .  
⇒ Only worry about non-seismic sources of infrasound.)
- Since  $\lambda \sim 30$  m, atmosphere is homogeneous on these scales: should be possible to use infrasonic pressure sensors to map exact pressure distribution.

## A. Infrasonic Newtonian Noise

Considerations ( $f \lesssim$  few Hz):

- Underground:  $l_{\text{coh}}/d$  may *not* be small.
- Correlation lengths shorter than seismic, and atmosphere is not stationary or homogeneous on these scales: may be harder to map and subtract.

## B. Thermal Newtonian Noise

- $\sqrt{S_h} \sim 4\pi G \rho \frac{c_T}{T} \left(\frac{v}{\omega}\right)^\alpha \omega^{-5/2} \times \begin{cases} v/\omega L & \text{if } \omega \gtrsim v/L \\ e^{-\omega d/v} & \text{if } \omega \lesssim v/d \end{cases}$

where  $\sqrt{\langle (\Delta T)^2 \rangle} = c_T (\Delta r)^\alpha$ ,  $\alpha \approx 1/3$ ,  $c_T \sim 0.5 \text{ K m}^{-1/3}$ ;  
 and  $v \sim 0 - 30 \text{ m/s}$  (advection).

⇒ Density perturbations *much* larger than infrasound, but  $v$  smaller.

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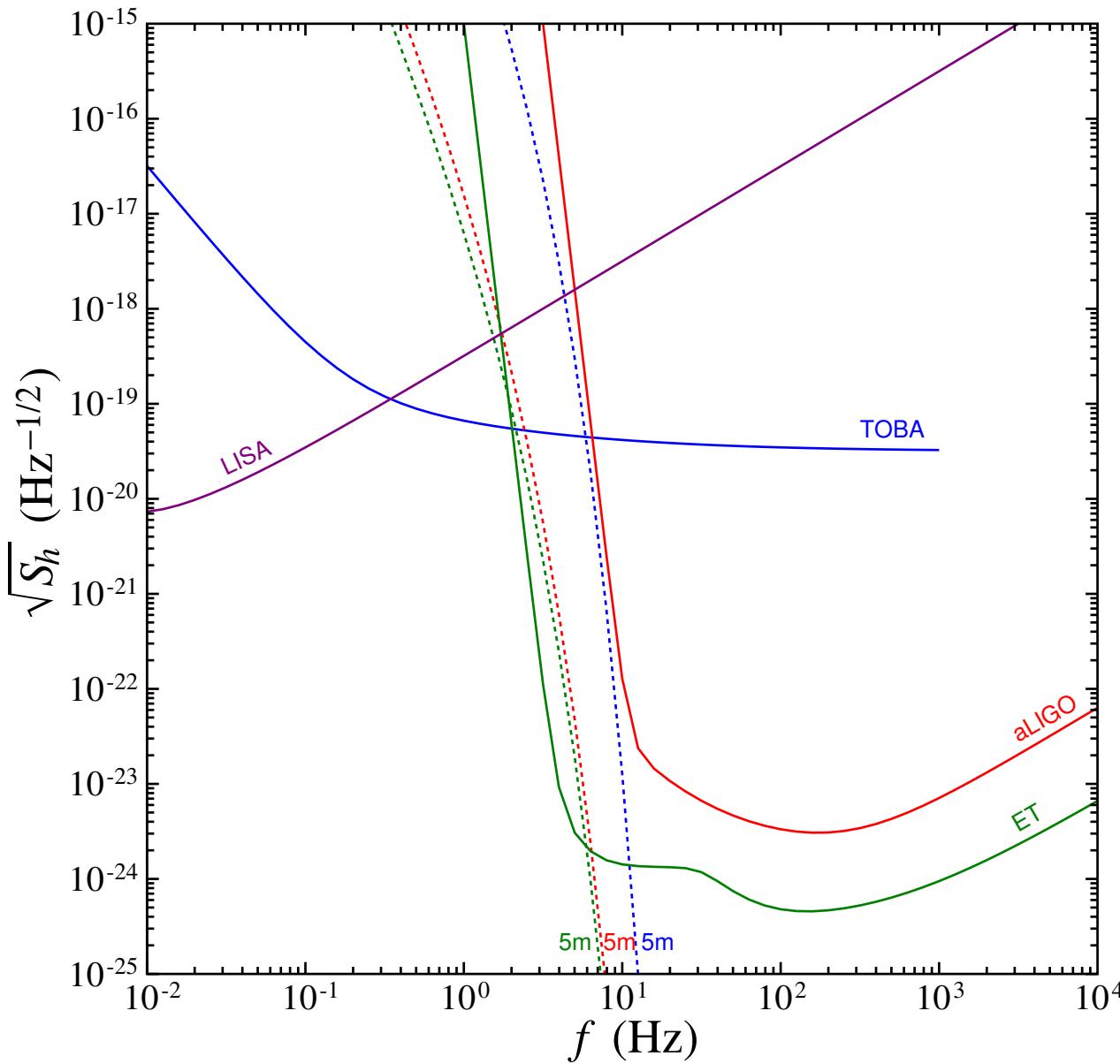
⇒ Density perturbations *much* larger than infrasound, but  $v$  smaller.

$$\sqrt{S_h} \approx \left(10^{-16} \text{ Hz}^{-1/2}\right) \left(\frac{c_T}{0.5 \text{ K m}^{-2/3}}\right) \left(\frac{f}{\text{Hz}}\right)^{-23/6} \quad (4\text{km})$$

$$\sqrt{S_h} \approx \left(10^{-13} \text{ Hz}^{-1/2}\right) \left(\frac{c_T}{0.5 \text{ K m}^{-2/3}}\right) \left(\frac{f}{\text{Hz}}\right)^{-17/6} \quad (\text{gradiometer})$$

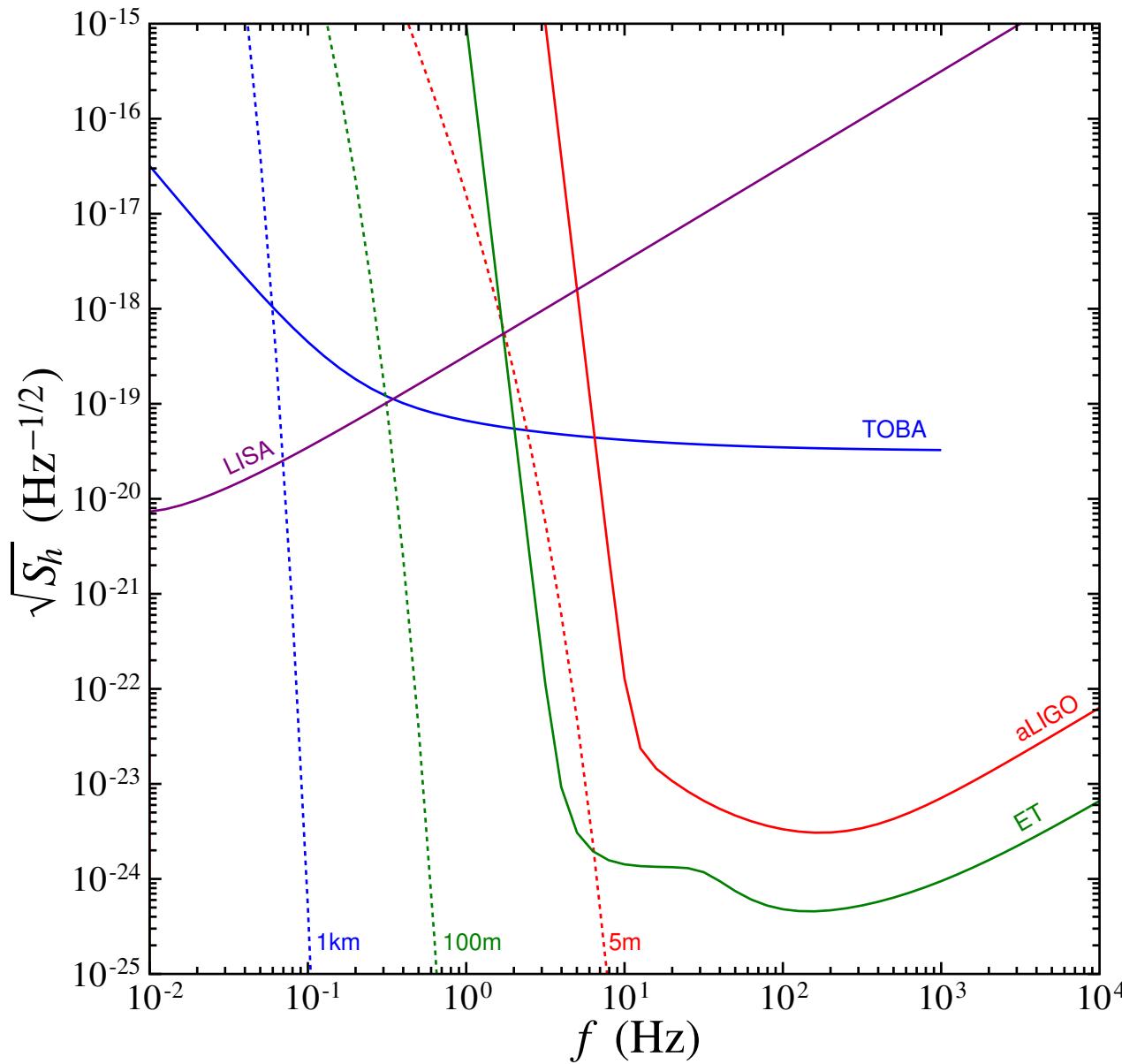
- But above a few Hz, cutoff at  $f \gtrsim v/2\pi d$  usually kills this, *provided* airflow is smooth.

## B. Thermal Newtonian Noise



Smooth airflow  
Surface structures

## B. Thermal Newtonian Noise



Smooth airflow  
Underground

## B. Thermal Newtonian Noise

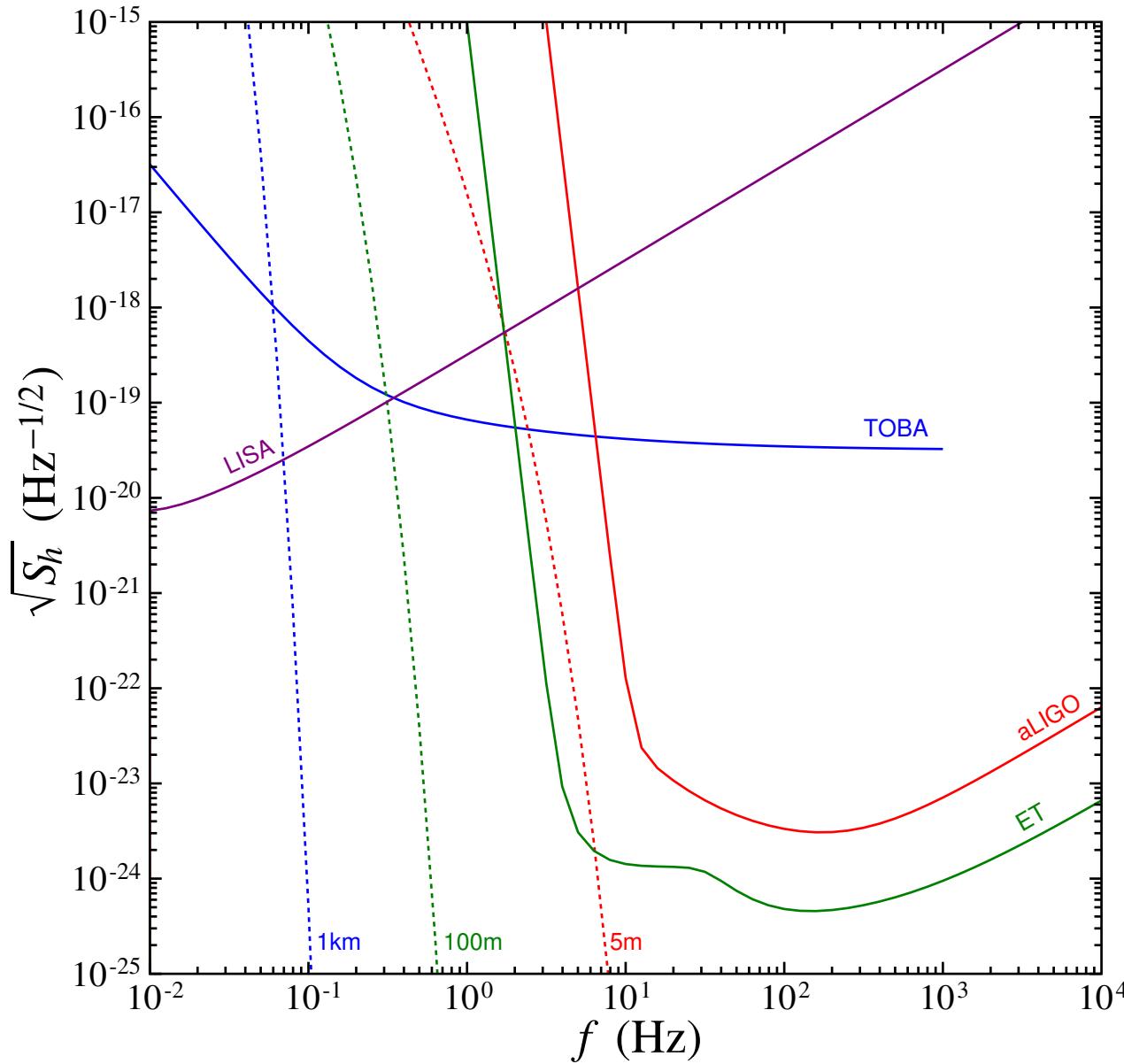
Considerations (any  $f$ ):

- If airflow is not smooth,  $g(t)$  may have discontinuous derivatives  
 $\Rightarrow$  power law tail in frequency.
  - ★ Discontinuous  $\Delta v$  :  $e^{-\omega d/v} \rightarrow v\Delta v/d^2\omega^2$ .
  - ★ Discontinuous  $\Delta a$  :  $e^{-\omega d/v} \rightarrow v\Delta a/d^2\omega^3$ .
  - ★ Vortices :  $e^{-\omega d/v} \rightarrow (v/d\omega)^2$ . (Heuristic!)
- Need hydrodynamic modeling to determine likely streamlines.

$$\sqrt{S_h} \sim \frac{4\pi G\rho}{L} \frac{c_T}{T} \left(\frac{v}{\omega}\right)^{\alpha+2} \omega^{-3/2} \sqrt{\int \left| \text{FT} \left\{ \frac{x(t)}{r(t)^3} \right\} \right|^2 dA}$$

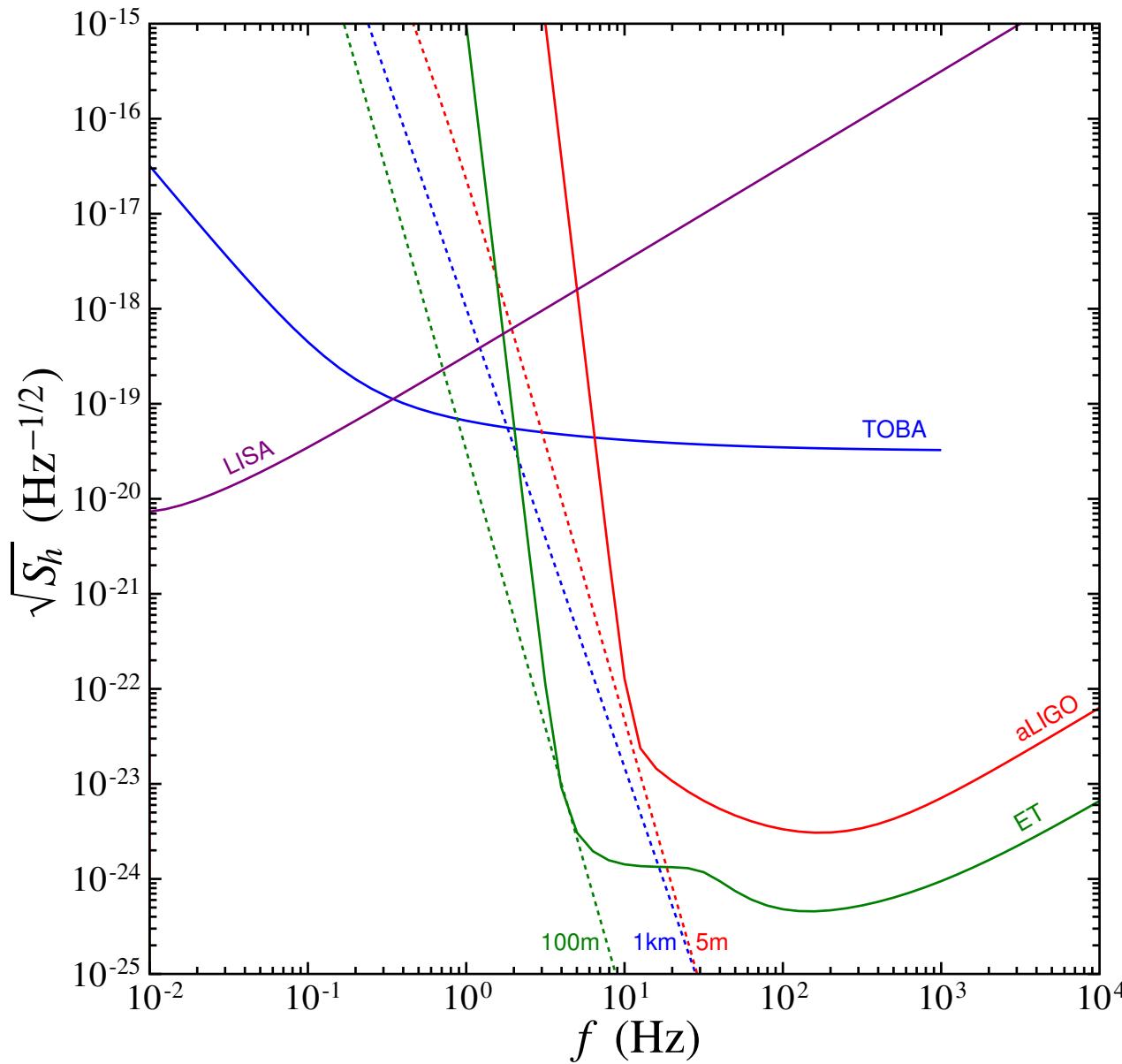
- Likely impossible to map and subtract.

## B. Thermal Newtonian Noise



Smooth airflow  
Underground

## B. Thermal Newtonian Noise



Turbulent airflow  
Underground

CAVEAT:  
approximations  
not based on  
hydrodynamic  
simulations

## C. Transients

- Sharp changes in atmospheric density:

$$\tilde{h} \sim \frac{4\pi G\rho}{L} \frac{v}{\omega} \frac{1}{\omega^3} \times \frac{\Delta p}{\gamma p} \sim 10^{-16} \left( \frac{f}{\text{Hz}} \right)^{-4} \quad (\text{fighter jet})$$

Solid objects passing near test masses:

$$\tilde{h} = \frac{GM}{L\omega^2} \times \text{FT} \left\{ \frac{x(t)}{r(t)^3} \right\} \sim 10^{-18} \left( \frac{f}{\text{Hz}} \right)^{-4} \quad (\text{tumbleweed})$$

- For  $f \gtrsim$  few Hz, assume we can detect & veto individual events; mitigators (exclusion areas, door springs) for anthropogenic sources.
- For  $f \lesssim$  few Hz, may be difficult! (Small- $N$  Poisson noise floor.)