

Higher order laser modes in gravitational wave detectors

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Higher order laser modes

- **Longitudinal sensing and control**
 - Plane wave calculation was sufficient
- **Alignment, mode matching, mode selection**
 - higher order modes need to be taken into account

Eigenmodes of the lasers

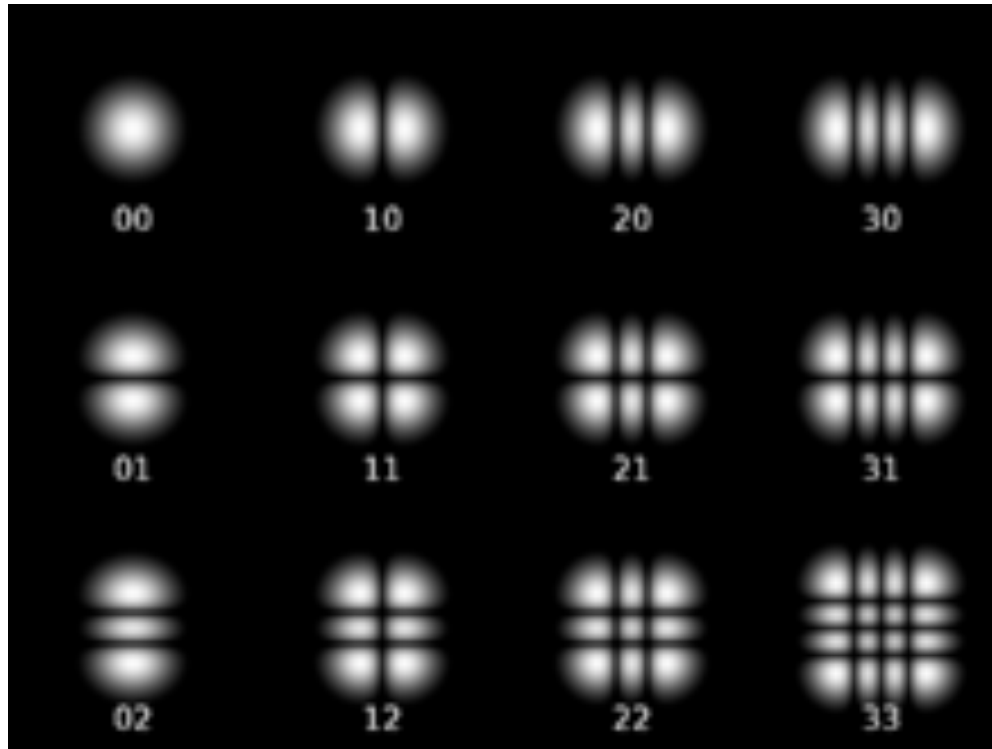
- Solution of Maxwell's equation for propagating electromagnetic wave under the paraxial approximation
- => Laser beams change their **intensity distributions** and **wavefront shapes** as they are propagated
- => **Any laser beam** can be decomposed and expressed as a **unique linear combination of eigenmodes**

In this sense, a (given) set of eigenmodes are ortho-normal basis

Hermite Gaussian modes

- HG modes (TEM mods) : one example of the eigenmodes

$$E_{mn}(x, y, z) = E_0 \frac{w_0}{w} H_m \left(\frac{\sqrt{2}x}{w} \right) H_n \left(\frac{\sqrt{2}y}{w} \right) \exp \left[-(x^2 + y^2) \left(\frac{1}{w^2} + \frac{jk}{2R} \right) - jkz - j(m + n + 1)\zeta(z) \right]$$

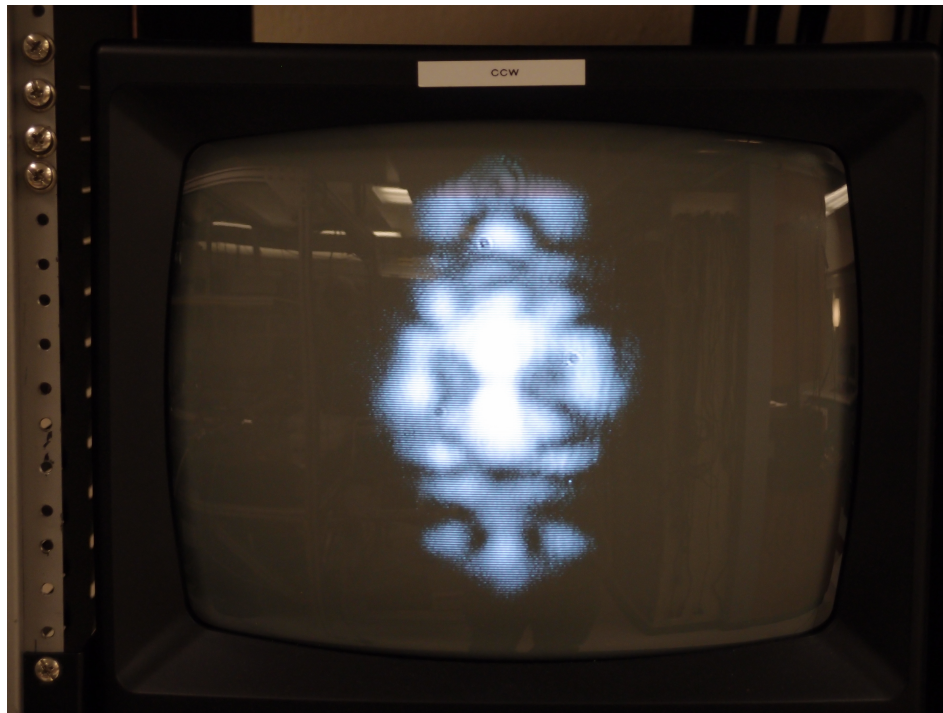


A. E. Siegman, *Lasers*, University Science Books, Mill Valley, CA (1986)

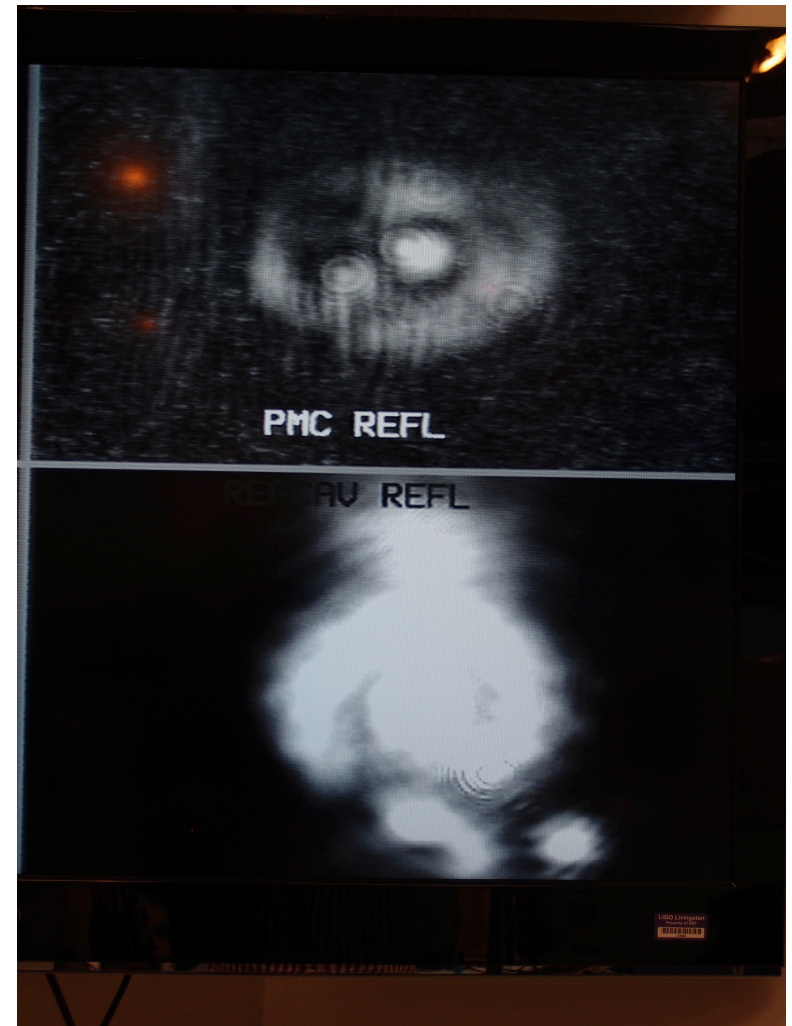
H. Kogelnik and T. Li, *Appl. Opt.* 5 (1966) 1550-1567

Wikipedia http://en.wikipedia.org/wiki/Transverse_mode

Any beam can be decomposed...



- Astronaut



- World cup football

Hermite Gaussian modes (TEM modes)

Trivia

- **There are infinite sets of HG modes**
 - A TEM₀₀ mode for an HG modes can be decomposed into infinite modes for other HG modes
- **The complex coefficients of the mode decomposition is invariant along the propagation axis**
 - Where ever the decomposition is calculated, the coefficients are unique.
- **No matter how a beam is decomposed, the laser frequency stays unchanged!**

(sounds trivial but frequently misunderstood)

Useful to note

- Beam size at z

$$w(z) = w_0 \sqrt{1 + \left(\frac{z}{z_R}\right)^2},$$

- Wavefront curvature at z

$$R(z) = z + \frac{z_R^2}{z},$$

- Gouy phase

$$\eta(z) = \tan^{-1} \left(\frac{z}{z_R} \right)$$

- Rayleigh range

$$j \frac{\pi w_0^2}{\lambda} = j z_R.$$

cf. Huygens' principle

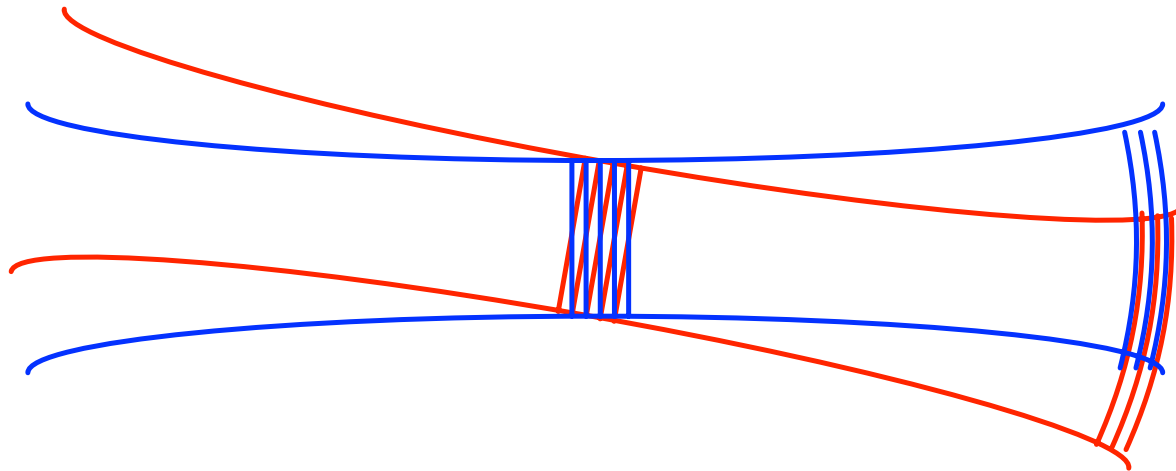
Gouy phase shift

$$E_{mn}(x, y, z) = E_0 \frac{w_0}{w} H_m \left(\frac{\sqrt{2}x}{w} \right) H_n \left(\frac{\sqrt{2}y}{w} \right) \exp \left[-(x^2 + y^2) \left(\frac{1}{w^2} + \frac{jk}{2R} \right) - jkz - j(m + n + 1)\zeta(z) \right]$$

- **Gouy phase shift:**

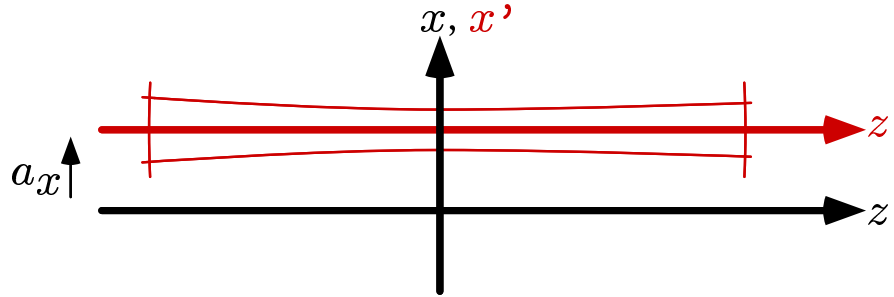
- Relative Phase shift between the transverse modes**

- Different optical phase of the modes for the same distance
=> Different resonant freq in a cavity (will see later)
- “Near field” and “Far field”



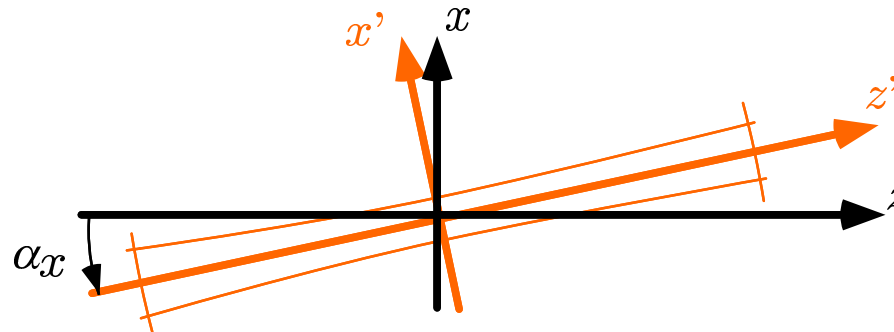
Decomposition of misaligned modes

■ Lateral shift:



$$P_x(a_x) * U_{00+}(x, y, z) \simeq \left[1 - \frac{1}{2} \left(\frac{a_x}{w_0} \right)^2 \right] U_{00+} + \frac{a_x}{w_0} U_{10+}$$

■ Rotational shift:



$$R_x(\alpha_x) * U_{00+}(x, y, z) \simeq \left[1 - \frac{1}{2} \left(\frac{\alpha_x}{\alpha_0} \right)^2 \right] U_{00+} - i \frac{\alpha_x}{\alpha_0} U_{10+}$$

"Cavity" eigenmodes

- TEM modes with matched wavefront RoC

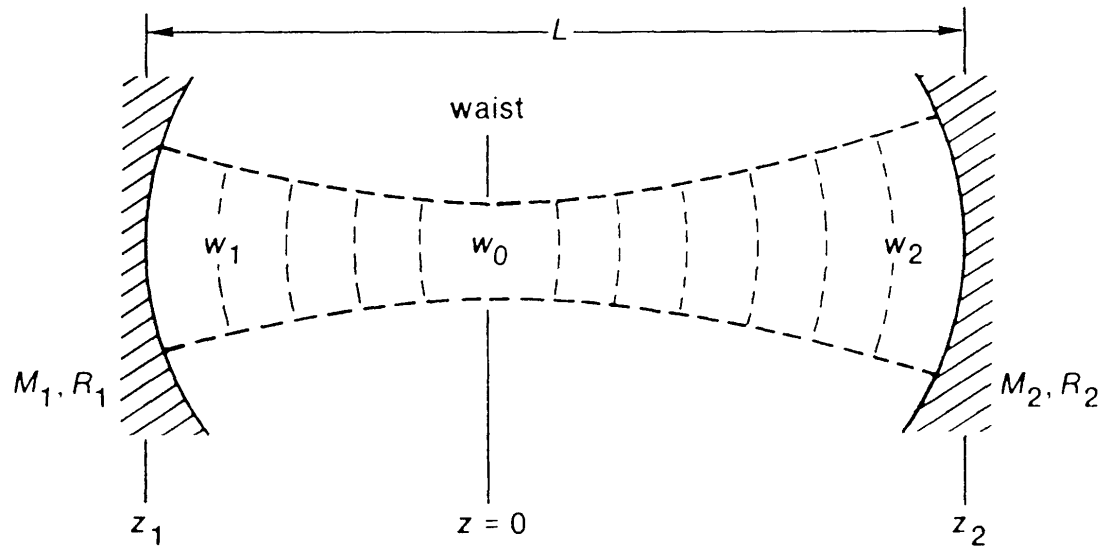
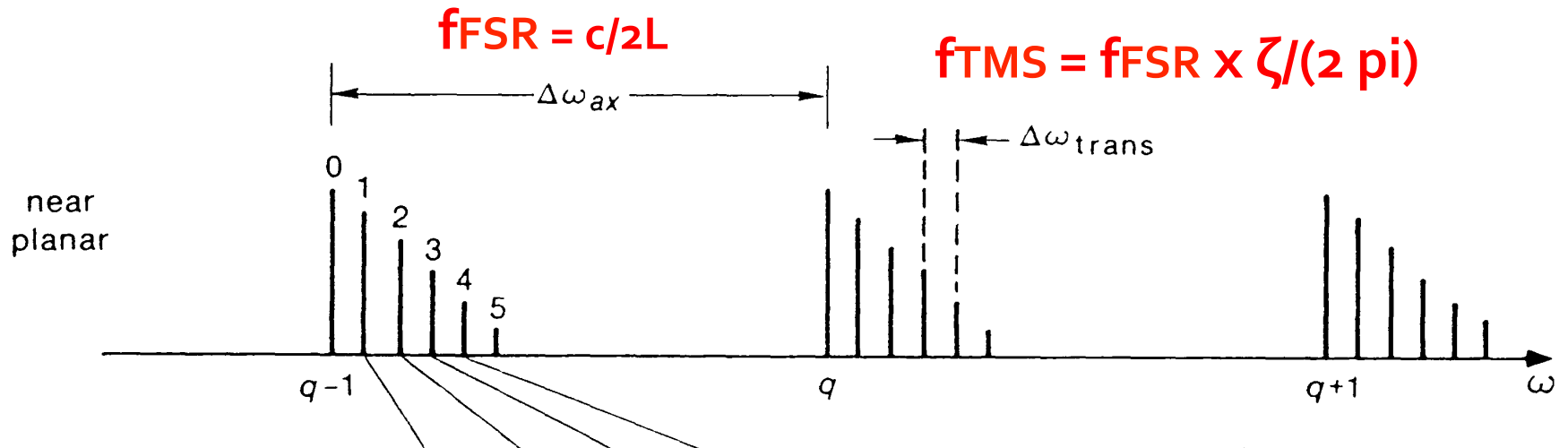


FIGURE 19.2

Notation and analytical model for analyzing a simple stable two-mirror cavity.

"Cavity" eigenmodes

- Due to different Gouy phase shifts between TEM modes, their resonant frequencies are different



ζ : cavity round trip Gouy phase shift

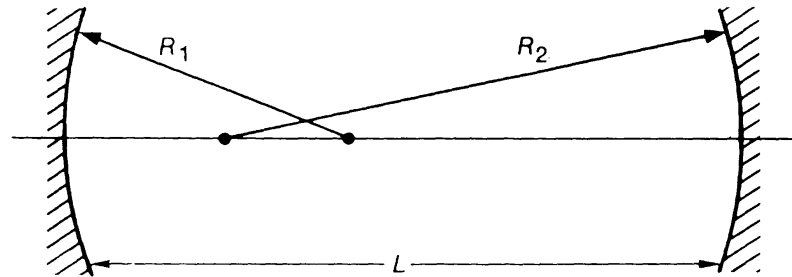
Optical resonator stability

■ g-factors

$$g_1 \equiv 1 - L/R_1$$

$$g_2 \equiv 1 - L/R_2$$

FIGURE 19.3
The resonator g parameters.



■ Stability criteria

$$0 \leq g_1 g_2 \leq 1.$$

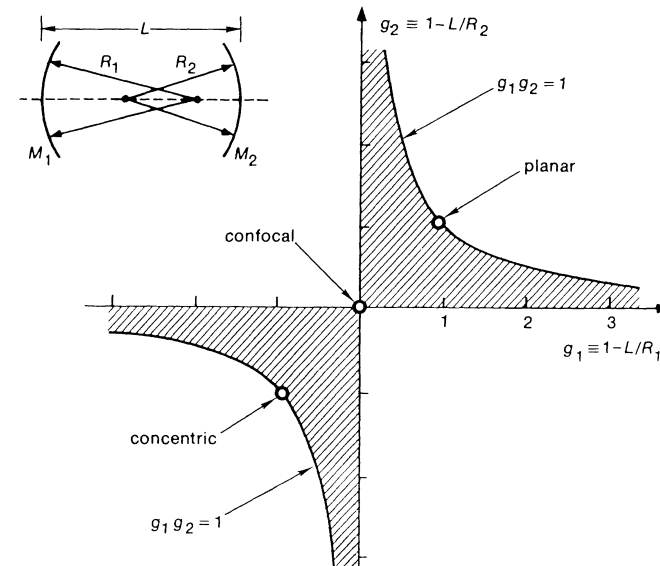


FIGURE 19.4
The stability diagram for a two-mirror optical resonator.

A. E. Siegman, *Lasers*,
University Science Books,
Mill Valley, CA (1986)

Optical resonator stability

- **General case**

derived that the accumulated round-trip Gouy phase shift can be computed only from the round-trip ABCD matrix of the cavity as:

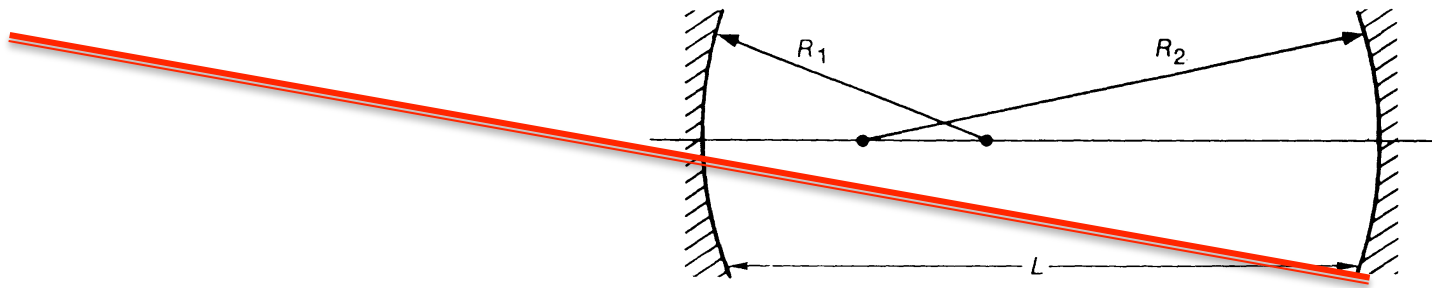
$$\zeta = \text{sgn}B \cdot \cos^{-1} \left(\frac{A + D}{2} \right) , \quad (12)$$

- **The cavity is stable when this quantity ζ exists**

$$-1 \leq \frac{A + D}{2} \leq 1.$$

Cavity alignment

- To match the input beam axis and the cavity axis



- Corresponds to the suppression of $TEM_{01/10}$ mode in the beam with regard to the cavity mode

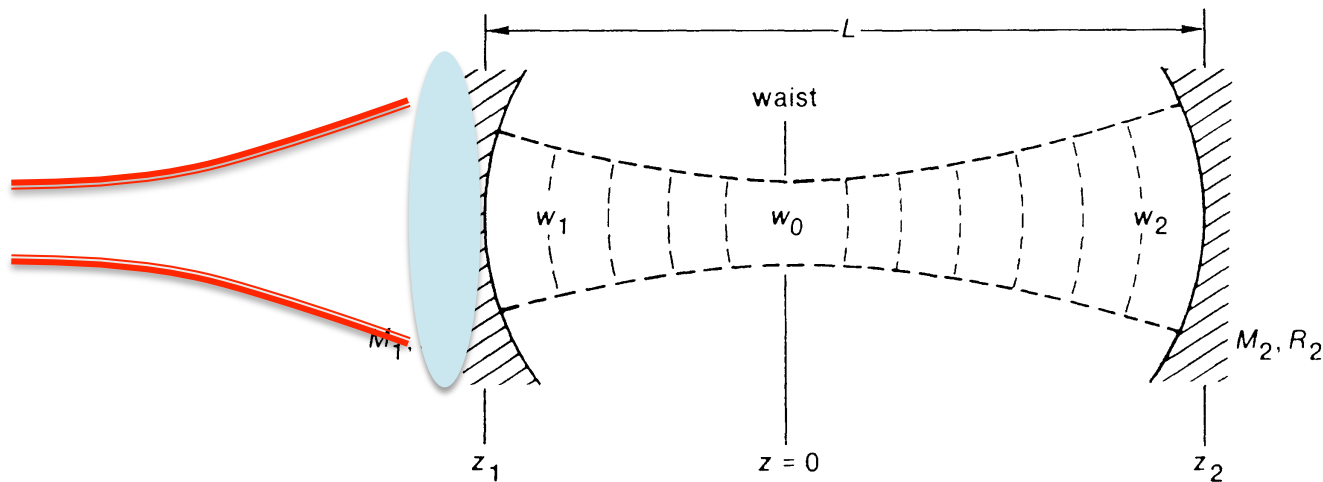
- 4 d.o.f.: (Horizontal, Vertical) x (translation, rotation)

Note: it is most intuitive to define the trans/rot at the waist

- To move the mirrors or to move the beam?

Cavity mode matching

- To match the waist size and position of the input beam to these of the cavity



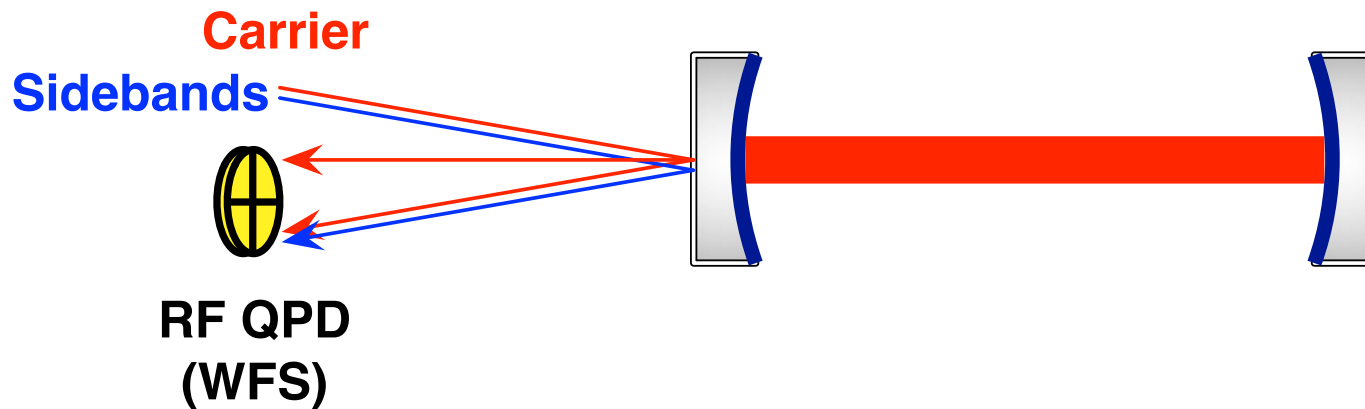
- Corresponds to the suppression of TEM_{02/11/20} mode in the beam with regard to the cavity mode

Angular global control

■ Wave Front Sensing

- Misalignment between the incident beam and the cavity axis
- **The carrier is resonant in the cavity**
- The reflection port has
 - Prompt reflection of the modulation sidebands
 - Prompt reflection of the carrier
 - Leakage field from the cavity internal mode

no signal
spatially distributed amplitude modulation

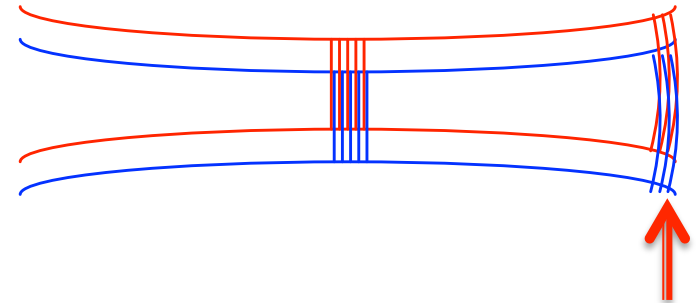


Angular global control

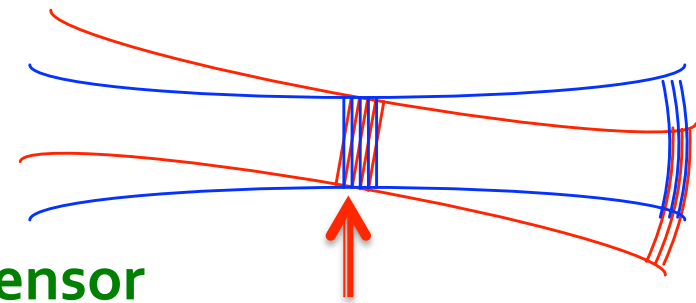
■ Wave Front Sensing

- WFS becomes sensitive when there is an angle between the wave fronts of the CA and SB
- Can detect rotation and translation of the beam separately, depending on the “location” of the sensor
- Use lens systems to adjust the “location” of the sensors.
i.e. Gouy phase telescope

Sensitive at the far field



Sensitive at the near field

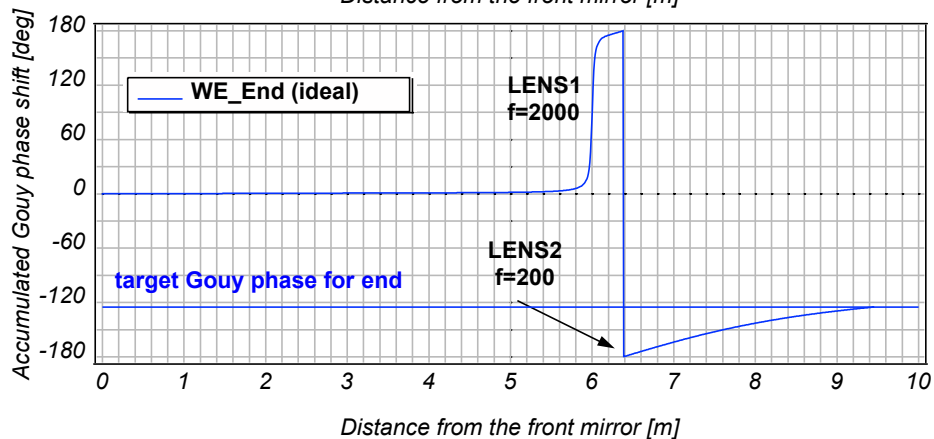
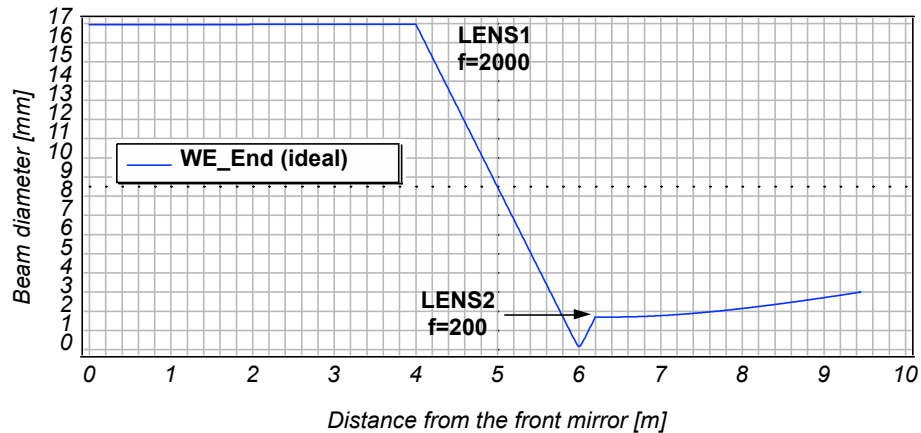


Angular global control

- Frequent mistake:**

What we want to adjust is the accumulated Gouy phase shift!

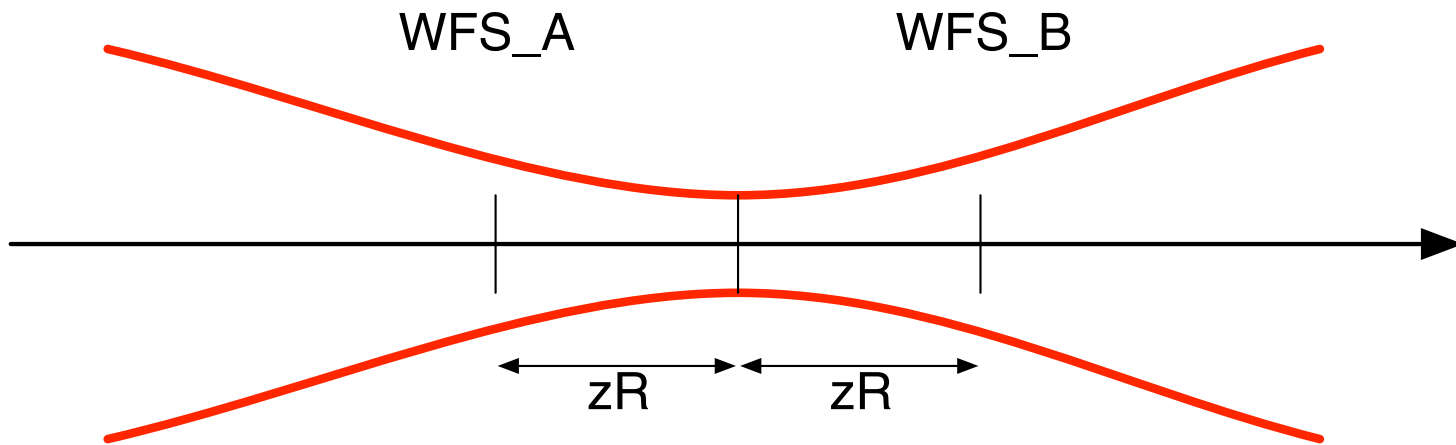
Not the one for the final mode!



Angular global control

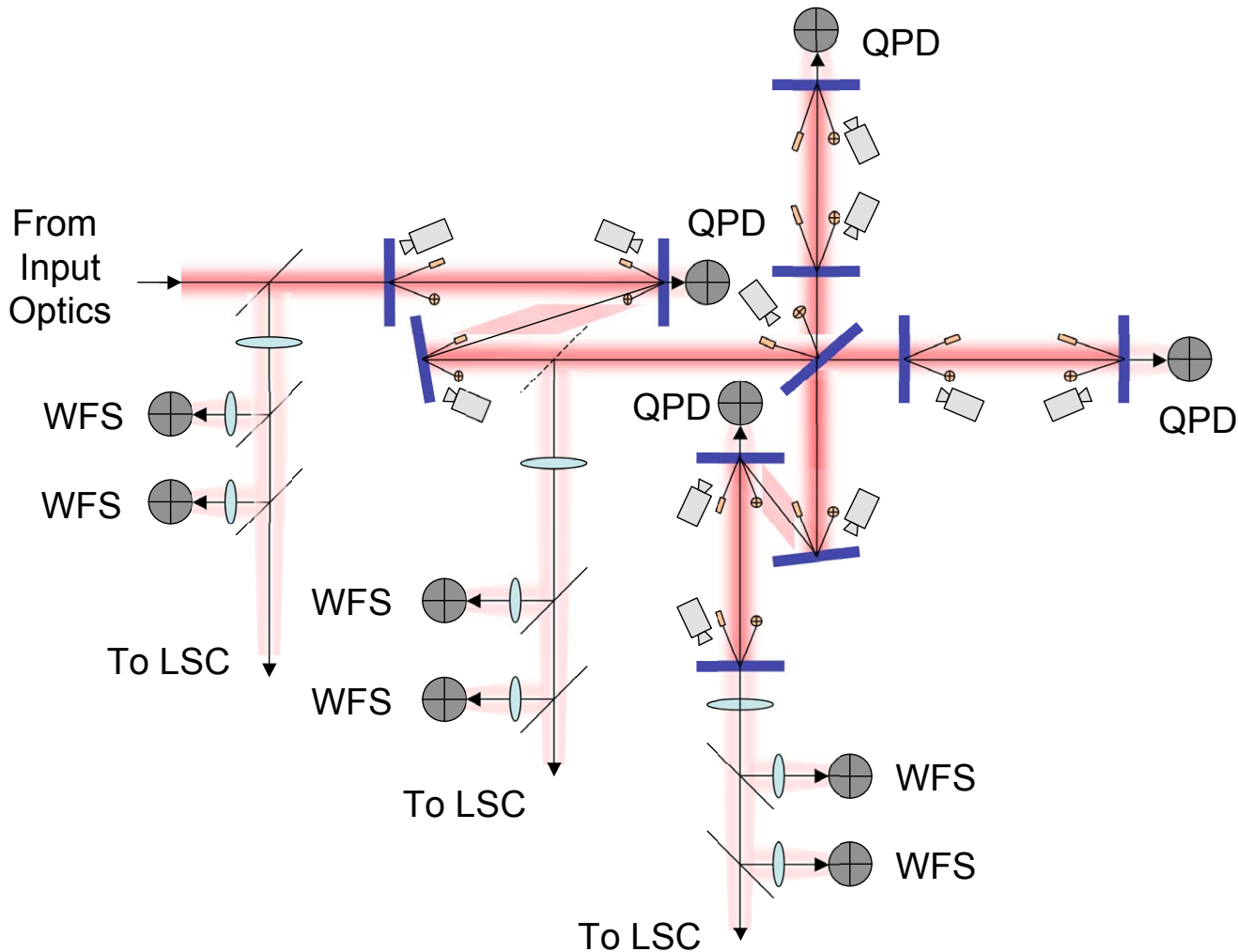
- **aLIGO implementation:**

Separate Gouy phase of a set of two WFSs with 90 deg



aLIGO angular control

- Combine WFS, DC QPD, digital CCD cameras



Other topics

- **PRC/SRC Degeneracy**
- **Sigg-Sidles instability & alignment modes**
 - G0900594
- **Impact on the noise**
 - G0900278 / P0900258
- **Parametric Instability**
 - HOM in the arm cavity
 - >Rad Press.
 - >Mirror acoustic mode
 - >Scattering of TEM₀₀->HOM