

Extraction of Gravitational Wave Signatures from Highly Magnetized Core-Collapse Supernovae

Adam Lloyd Bruce

University of California, Berkeley

Caltech LIGO Summer Undergraduate Research Fellowship (SURF)
Presentations



Purpose

- Core-Collapse supernovae (CCSNe) are expected to be good transient sources for GWs.



Purpose

- **Core-Collapse supernovae** (CCSNe) are expected to be good **transient sources** for GWs.
- We can use physics we already know (General Relativity + Nuclear Physics + Magnetohydrodynamics) to **simulate** CCSNe and **isolate** the expected GW signatures.



Purpose

- **Core-Collapse supernovae** (CCSNe) are expected to be good **transient sources** for GWs.
- We can use physics we already know (General Relativity + Nuclear Physics + Magnetohydrodynamics) to **simulate** CCSNe and **isolate** the expected GW signatures.
- This gives us an idea of where CCSNe signatures may appear in the **LIGO frequency band**.



Purpose

- **Core-Collapse supernovae** (CCSNe) are expected to be good **transient sources** for GWs.
- We can use physics we already know (General Relativity + Nuclear Physics + Magnetohydrodynamics) to **simulate** CCSNe and **isolate** the expected GW signatures.
- This gives us an idea of where CCSNe signatures may appear in the **LIGO frequency band**.
- We run these simulations (with full GR!) using Caltech's **Zelmani** suite of **Numerical Relativity** tools.



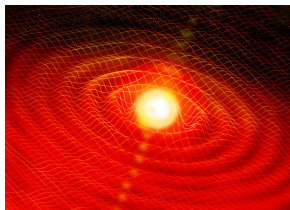
Purpose

- **Core-Collapse supernovae** (CCSNe) are expected to be good **transient sources** for GWs.
- We can use physics we already know (General Relativity + Nuclear Physics + Magnetohydrodynamics) to **simulate** CCSNe and **isolate** the expected GW signatures.
- This gives us an idea of where CCSNe signatures may appear in the **LIGO frequency band**.
- We run these simulations (with full GR!) using Caltech's **Zelmani** suite of **Numerical Relativity** tools.

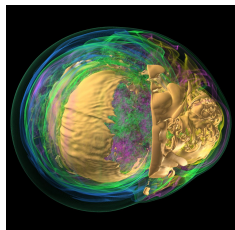


Purpose

- **Core-Collapse supernovae** (CCSNe) are expected to be good **transient sources** for GWs.
- We can use physics we already know (General Relativity + Nuclear Physics + Magnetohydrodynamics) to **simulate** CCSNe and **isolate** the expected GW signatures.
- This gives us an idea of where CCSNe signatures may appear in the **LIGO frequency band**.
- We run these simulations (with full GR!) using Caltech's **Zelmani** suite of **Numerical Relativity** tools.



(k) GW from Sne (Artist's Rendering)

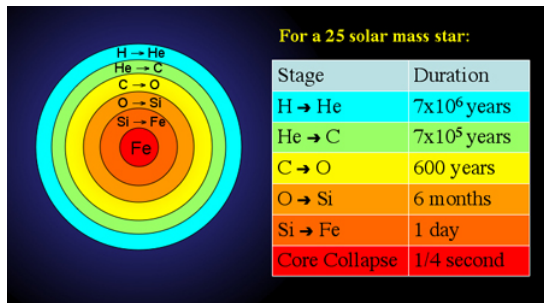


(l) Argonne Nat Lab (3D Simulation)



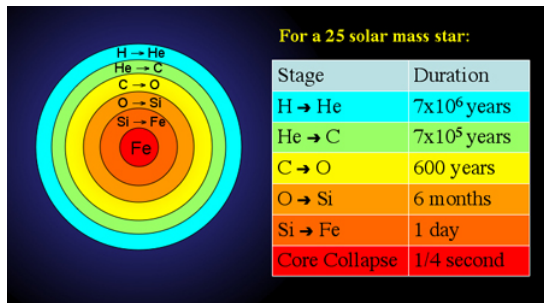
Core-Collapse Basics: Why do We Care?

- CCSNe result from **gravitational collapse** of the core of a **massive star** ($M > \sim 8-10M_{sun}$); observed as Type II, Ib, Ic SNe.



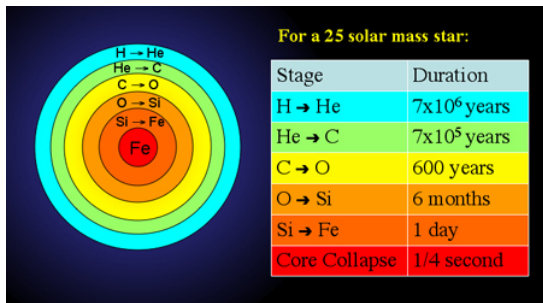
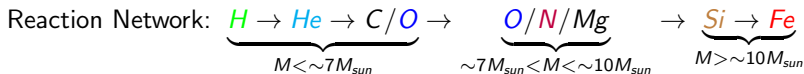
Core-Collapse Basics: Why do We Care?

- CCSNe result from **gravitational collapse** of the core of a **massive star** ($M > \sim 8-10M_{sun}$); observed as Type II, Ib, Ic SNe.
- Birth sites of **neutron stars** and **black holes**.



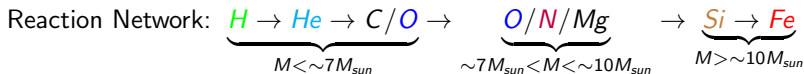
Core-Collapse Basics: Why do We Care?

- CCSNe result from **gravitational collapse** of the core of a **massive star** ($M > \sim 8-10M_{sun}$); observed as Type II, Ib, Ic SNe.
- Birth sites of **neutron stars** and **black holes**.
- **Nucleosynthesizes Fe-group elements** \Rightarrow "Iron Core."

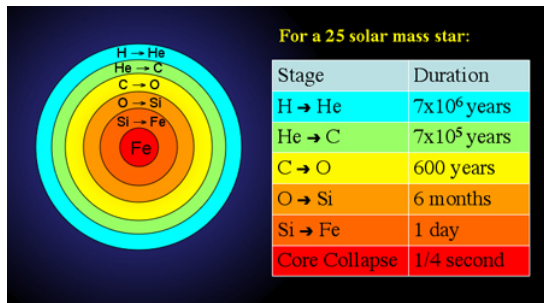


Core-Collapse Basics: Why do We Care?

- CCSNe result from **gravitational collapse** of the core of a **massive star** ($M > \sim 8-10M_{sun}$); observed as Type II, Ib, Ic SNe.
- Birth sites of **neutron stars** and **black holes**.
- **Nucleosynthesizes Fe-group elements** \Rightarrow "Iron Core."



- Rate of 1 to ~ 4 SN per 90 yrs in the **Local Group** (most events within ~ 100 kpc of Earth.)



Eta Carinae (NASA WFPC2): likely CCSNe.



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**
 - affected by many intricate processes (photodissociation, electron capture, neutrino trapping, etc.)



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**
 - affected by many intricate processes (photodissociation, electron capture, neutrino trapping, etc.)
- At **nuclear statistical equilibrium** ($\bar{\rho} \sim 10^{14} \text{ g/cm}^3$) the inner core **“bounces”** into the **supersonically infalling** outer core.



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**
 - affected by many intricate processes (photodissociation, electron capture, neutrino trapping, etc.)
- At **nuclear statistical equilibrium** ($\bar{\rho} \sim 10^{14} \text{ g/cm}^3$) the inner core “**bounces**” into the **supersonically infalling** outer core.
- High-energy **shockwave** is created as the proto-**neutron star** at the CCSNe center and inner core **supersonically collide**.



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**
 - affected by many intricate processes (photodissociation, electron capture, neutrino trapping, etc.)
- At **nuclear statistical equilibrium** ($\bar{\rho} \sim 10^{14} \text{ g/cm}^3$) the inner core “**bounces**” into the **supersonically infalling** outer core.
- High-energy **shockwave** is created as the proto-**neutron star** at the CCSNe center and inner core **supersonically collide**.
 - evolves into a **standing accretion shock**, immediately losing most of its kinetic energy via **photodissociation**.



Core-Collapse Mechanism

- **Radial instability** occurs as $M \geq M_{\text{Ch}} \approx 5.8(Y_e)^2 M_{\text{sun}}$ and **causes collapse**
 - affected by many intricate processes (photodissociation, electron capture, neutrino trapping, etc.)
- At **nuclear statistical equilibrium** ($\bar{\rho} \sim 10^{14} \text{ g/cm}^3$) the inner core “**bounces**” into the **supersonically infalling** outer core.
- High-energy **shockwave** is created as the proto-**neutron star** at the CCSNe center and inner core **supersonically collide**.
 - evolves into a **standing accretion shock**, immediately losing most of its kinetic energy via **photodissociation**.
 - Shock can be revived via **neutrino heating**, **SASI**, and **magnetorotational effects**, causing a **massive detonation** which releases $\sim 10^{51}$ ergs of energy!



2D Simulation Example

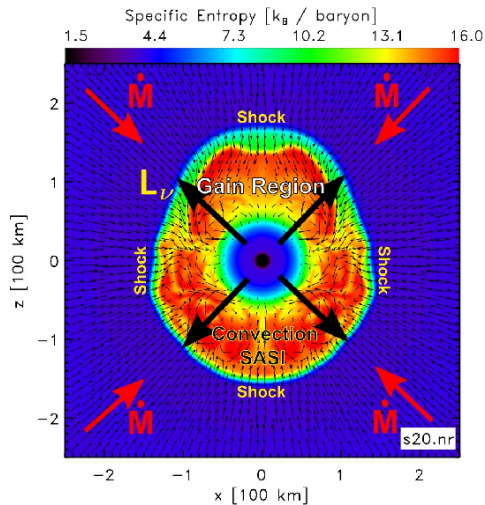


Figure: Shortly after first bounce (~ 20 -50 ms evolution).



- Typically use the **quadrupole formula**:

$$h_{ij}^{TT}(t, \mathbf{x}) = \left[\frac{2}{c^4} \frac{G}{|\mathbf{x}|} \ddot{I}_{ij} \left(t - \frac{|\mathbf{x}|}{c} \right) \right]^{TT} \quad (1)$$



- Typically use the **quadrupole formula**:

$$h_{ij}^{TT}(t, \mathbf{x}) = \left[\frac{2}{c^4} \frac{G}{|\mathbf{X}|} \ddot{I}_{ij} \left(t - \frac{|\mathbf{X}|}{c} \right) \right]^{TT} \quad (1)$$

- First proposed 1st-order approximation; recently shown to recover the **analytic quadrupole** for CCSNe **almost exactly!** [C. Reisswig (2011)]



- Typically use the **quadrupole formula**:

$$h_{ij}^{TT}(t, \mathbf{x}) = \left[\frac{2}{c^4} \frac{G}{|\mathbf{X}|} \ddot{I}_{ij} \left(t - \frac{|\mathbf{X}|}{c} \right) \right]^{TT} \quad (1)$$

- First proposed 1st-order approximation; recently shown to recover the **analytic quadrupole** for CCSNe **almost exactly!** [C. Reisswig (2011)]
- implementation boils down to finding the second-time derivative of the **quadrupole tensor**, \ddot{I}_{ij} .



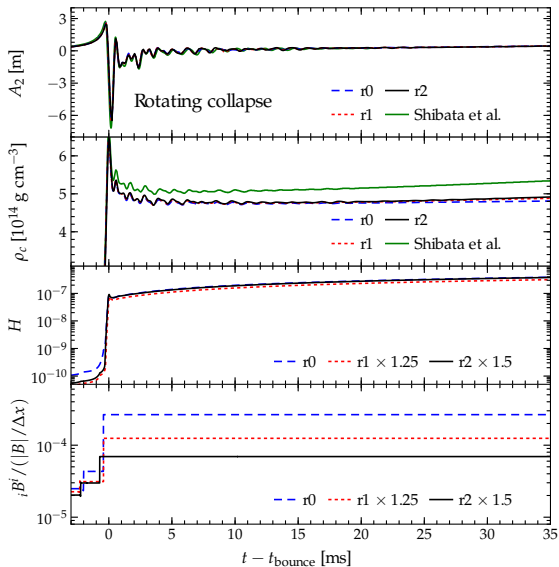


Figure: Mösta et al. (2013) [MHD in simulation, not extraction.]



Duelling Models

- Derived two possible models to include MHD in the GW extraction:



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" T^{ij} model")



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" T^{ij} model")
 - (ii) $I^{ij} = \int T^{00} x^i x^j d^3x \sim \int (\rho + b^2) x^i x^j d^3x$ (" b^2 model")



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" T^{ij} model")
 - (ii) $I^{ij} = \int T^{00} x^i x^j d^3x \sim \int (\rho + b^2) x^i x^j d^3x$ (" b^2 model")
 - (ρ = density, b^2 = Norm-Squared of comoving B -Field.)



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" T^{ij} model")
 - (ii) $I^{ij} = \int T^{00} x^i x^j d^3x \sim \int (\rho + b^2) x^i x^j d^3x$ (" b^2 model")
 - (ρ = density, b^2 = Norm-Squared of comoving B -Field.)



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" T^{ij} model")
 - (ii) $I^{ij} = \int T^{00} x^i x^j d^3x \sim \int (\rho + b^2) x^i x^j d^3x$ (" b^2 model")
 - (ρ = density, b^2 = Norm-Squared of comoving B -Field.)

Early Result I:

T^{ij} model discredited in early simulations! Incorrect morphology now explained analytically via the absence of a Green's function (violation of causality).



Duelling Models

- Derived two possible models to include MHD in the GW extraction:
 - (i) $\ddot{I}^{ij} \sim 2 \int T_{\text{eff}}^{ij} d^3x$ (" **T^{ij}** model")
 - (ii) $I^{ij} = \int T^{00} x^i x^j d^3x \sim \int (\rho + b^2) x^i x^j d^3x$ (" **b^2** model")
 - (ρ = density, b^2 = Norm-Squared of comoving B -Field.)

Early Result I:

T^{ij} model discredited in early simulations! Incorrect morphology now explained analytically via the absence of a Green's function (violation of causality).

Early Result II:

b^2 contributions **don't affect signature** in early simulation \Rightarrow initially $\mathcal{O}(\text{rotational effects}) \ll \mathcal{O}(\text{hydrodynamical effects})$.



Simulation Progress

- Simulation: currently on the Zwicky supercomputing cluster, will run for $\sim 120 - 160$ hours in 12 hour iterations before the full simulation is complete



Simulation Progress

- Simulation: currently on the Zwicky supercomputing cluster, will run for $\sim 120 - 160$ hours in 12 hour iterations before the full simulation is complete
- Currently on the 4th iteration: just observed successful first bounce!



Simulation Progress

- Simulation: currently on the Zwicky supercomputing cluster, will run for $\sim 120 - 160$ hours in 12 hour iterations before the full simulation is complete
- Currently on the 4th iteration: just observed successful first bounce!
- effects appear to be at 5-10% level, as expected.



Future Directions

- Extraction Specific: Implement higher-order magnetic effects in other hydrodynamics, off-diagonal terms ($b_i b_j$).

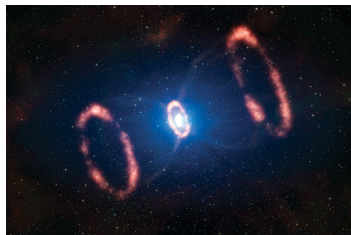


Figure: CCSNe 87A (European Southern Observatory)



Future Directions

- Extraction Specific: Implement higher-order magnetic effects in other hydrodynamics, off-diagonal terms ($b_i b_j$).
- General: Implement high fidelity microphysics into the Zelmani code.

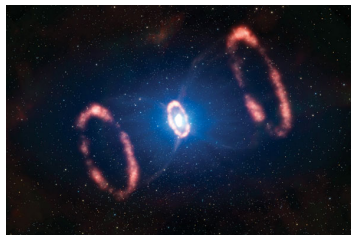


Figure: CCSNe 87A (European Southern Observatory)



Future Directions

- Extraction Specific: Implement higher-order magnetic effects in other hydrodynamics, off-diagonal terms ($b_i b_j$).
- General: Implement high fidelity microphysics into the Zelmani code.
 - mostly done by postdoc Philipp Mösta

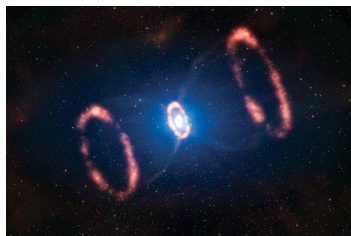


Figure: CCSNe 87A (European Southern Observatory)



Future Directions

- Extraction Specific: Implement higher-order magnetic effects in other hydrodynamics, off-diagonal terms ($b_i b_j$).
- General: Implement high fidelity microphysics into the Zelmani code.
 - mostly done by postdoc Philipp Mösta
 - author will implement tracer particles to track nucleosynthesis (started).

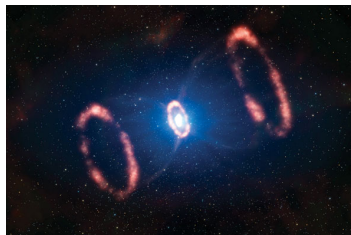


Figure: CCSNe 87A (European Southern Observatory)



Questions?



© Stephan Pastis/Dist. by LFS, Inc.

