

Suspension Modeling in Mathematica™

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Toolkit Features

 The toolkit is a Mathematica "package", PendUtil.nb, for specifying different configurations (e.g., quad, triple etc) in a (relatively) userfriendly way

• Supported features:

- » 6-DOF rigid bodies for masses (no internal modes)
- » Springs described by an elasticity tensor and a vector of pre-load forces
- » Massless wires (i.e., no violin modes) but detailed elasticity model from beam equation
- » Arbitrary frequency-dependent damping on all sources of elasticity
- » Symbolic up to the point of minimizing the potential to find the equilibrium position
- » Calculates elasticity and mass matrices semi-numerically (symbolic partial derivatives of functions with mostly numeric coefficients)
- » Eigenfrequencies and eigenmodes calculated numerically
- » Reasonable runtime:
 - 2 minutes for quad model with just wire longitudinal elasticity (adequate for most control theory purposes)
 - 2 hours with wire bending elasticity (required for thermal noise estimates)
- » Structured to make version control easy



Normal Mode Calculation (i)

Express the potential energy of the system in terms of the coordinates:

 $E_P = E_P\left(x_1, \dots x_n\right) = E_P\left(\mathbf{x}\right)$

 Express the kinetic energy of the system in terms of the coordinates and coordinate velocities:

 $E_K = E_K \left(x_1, \dots x_n, \dot{x}_1, \dots \dot{x}_n \right)$

 Minimize the potential energy to find the equilibrium values of the coordinates.

$$\mathbf{x}_{eq} = \left(x_{1(eq)}, \dots x_{n(eq)}\right)^{\mathrm{T}}$$



Normal Mode Calculation (ii)

- Create a matrix of second derivatives of the potential energy, a.k.a., the potential energy matrix or the stiffness matrix. $\mathbf{K}:K_{ij} = \frac{\partial E_P}{\partial x_i \partial x_j}\Big|_{\mathbf{x}=\mathbf{x}_{eq}} \qquad E_P = E_P(\mathbf{x}_{eq}) + \frac{1}{2}(\mathbf{x}-\mathbf{x}_{eq})^{\mathrm{T}} \mathbf{K}(\mathbf{x}-\mathbf{x}_{eq})$
- Create a matrix of second derivatives w.r.t. velocity, a.k.a., the kinetic energy matrix or the mass matrix.

$$\mathbf{M}: M_{ij} = \frac{\partial E_K}{\partial \dot{x}_i \partial \dot{x}_j} \bigg|_{\substack{\dot{\mathbf{x}}=0\\\mathbf{x}=\mathbf{x}_{eq}}} E_K = \frac{1}{2} \dot{\mathbf{x}}^{\mathrm{T}} \mathbf{M} \dot{\mathbf{x}}$$



Normal Mode Calculation (iii)

• Do a simultaneous diagonalization of the stiffness and mass matrices to obtain the eigenfrequencies and eigenmodes:

$$\mathbf{K}\mathbf{e}_{i} = \boldsymbol{\omega}_{i}^{2}\mathbf{M}\mathbf{e}_{i} \qquad \mathbf{x}_{i}(t) = \mathbf{x}_{eq} + \mathbf{e}_{i}e^{\boldsymbol{\omega}_{i}t} \qquad f_{i} = \boldsymbol{\omega}_{i}/2\pi$$

 For a practical calculation potential matrix step needs to be considerably elaborated, partly for efficiency and partly to support additional calculations such as transfer functions and thermal noise estimates.



Other coordinates

 Also need to consider coordinates of structure - constant during normal mode motion but movable when injecting displacement inputs:

$$\mathbf{s} = (s_1, \dots s_l)^{\mathrm{T}} \qquad \mathbf{s}_{nom} = (s_{1(nom)}, \dots s_{l(nom)})^{\mathrm{T}}$$

 And "floats", coordinates of things such as junctions between elastic elements in series - not independent of normal mode coordinates:

$$\mathbf{q} = (q_1, \dots, q_m)^{\mathrm{T}}$$

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Master Potential Matrix

 To work with all types of coordinates efficiently, define master potential matrix:

$$\mathbf{P}: E_P = E_P \left(\mathbf{x}_{eq}, \mathbf{q}_{eq}, \mathbf{s}_{nom} \right) + \frac{1}{2} \left(\mathbf{x}^{\mathrm{T}} - \mathbf{x}_{eq}^{\mathrm{T}} - \mathbf{q}_{eq}^{\mathrm{T}} - \mathbf{q}_{eq}^{\mathrm{T}} - \mathbf{s}_{nom}^{\mathrm{T}} \right) \mathbf{P} \begin{pmatrix} \mathbf{x} - \mathbf{x}_{eq} \\ \mathbf{q} - \mathbf{q}_{eq} \\ \mathbf{s} - \mathbf{s}_{nom} \end{pmatrix}$$

• It has a block structure with many useful submatrices:

$$\mathbf{P} = \begin{pmatrix} \mathbf{K} & \mathbf{C}_{XQ} & \mathbf{C}_{XS} \\ \mathbf{C}_{QX} & \mathbf{Q} & \mathbf{C}_{QS} \\ \mathbf{C}_{SX} & \mathbf{C}_{SQ} & \mathbf{S} \end{pmatrix} \qquad \qquad \mathbf{C}_{XQ} = \mathbf{C}_{QX}^{\mathbf{T}}$$

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LIGO Effective Potential and Coupling Matrices

 If there are any float coordinates, K submatrix of P is not appropriate to use in the normal mode analysis, since it assumes q=const whereas actually:

$$\mathbf{q} = \mathbf{q}_{eq} - \mathbf{Q}^{-1} \mathbf{C}_{QX} \left(\mathbf{x} - \mathbf{x}_{eq} \right)$$

• The effective potential matrix is

$$\mathbf{K}_{eff} = \mathbf{K} - \mathbf{C}_{XQ} \mathbf{Q}^{-1} \mathbf{C}_{QX}$$

 Similarly the effective coupling matrix converting displacement inputs of the structure to forces on the normal mode coordinates is

$$f_{xs} = \mathbf{C}_{XS(eff)} \left(\mathbf{s} - \mathbf{s}_{nom} \right) = \left(C_{XS} - \mathbf{C}_{XQ} \mathbf{Q}^{-1} \mathbf{C}_{QS} \right) \left(\mathbf{s} - \mathbf{s}_{nom} \right)$$





Dissipation Dilution and Pendulums (i)

- Two independent reasons why pendulums have low loss:
 - » Restoring force is gravitational
 - » Restoring force is sideways component of a tension
- Reason #2 would still apply if the tension were supplied by a second spring:
- Why? Because when a spring is used to create a restoring force by first generating a static force and then coupling that to the load by a variable mechanical advantage, the length change is only second order in amplitude.





Dissipation Dilution and Pendulums (ii)

 Why is it important to get this right? Because the normal mode formalism mixes up the two cases depending on the coordinates used and the stiffness of the wire:



- No stretch of spring for pendulum mode in polar coordinates vs. second order stretch in Cartesian coordinates.
- Solution: recompute potential matrix with tension zeroed out, then: $\mathbf{P} = \sum \left(\mathbf{P}_i \Big|_{tension_off} \left(\boldsymbol{\varepsilon}'_i(f) + i \boldsymbol{\varepsilon}''_i(f) \right) \right) + \sum \left(\mathbf{P}_i \Big|_{tension_off} - \mathbf{P}_i \Big|_{tension_off} \right)$



Equations of Motion

• The net equation of motion is then,

$$\mathbf{K}_{eff}\left(\mathbf{x} - \mathbf{x}_{eq}\right) + \mathbf{M}\dot{\mathbf{x}} = \mathbf{f}_{x} + \mathbf{C}_{XS(eff)}\left(\mathbf{s} - \mathbf{s}_{nom}\right)$$

• Or in the frequency domain:

$$\mathbf{K}_{eff} \left(\mathbf{x} - \mathbf{x}_{eq} \right) - \left(2\pi f \right)^2 \mathbf{M} \left(\mathbf{x} - \mathbf{x}_{eq} \right) = \mathbf{f}_x + \mathbf{C}_{XS(eff)} \left(\mathbf{s} - \mathbf{s}_{nom} \right)$$

- This can be solved for x for a sequence of different values of f to give force-input or displacement-input transfer functions as a function of frequency.
- Thermal noise is calculated in usual way from complex admittance.



Models

- Two major families of models have been defined:
 - » The triple models reflect a generic GEO-style pendulum with 3 masses, 6 blade springs and 10 wires.
 - » The quad models reflect a standard AdvLIGO quad pendulum, with 4 masses, 6 blade springs and 14 wires.
- Within each family there are three variants
 - » The "full" version, where the tips of the blade springs are modeled as 6-DOF rigid bodies attached to their bases by 6-DOF springs
 - » The "lite" version, where the tips of the blade springs are connected to their bases by geometrical constraints in 5 DOFs and elastically in 1 DOF.
 - » The xtra-lite version, where the tips of the blade springs are massless.
- The "xtra-lite" models are preferred for time-domain simulation because they have the smallest matrices and no high-frequency eigenmodes due to the blades.
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Triple Pendulum Model

- 2 blade springs
- 2 wires
- "upper" mass
- 4 blade springs
- 4 wires
- "intermediate" mass
- 4 fibres

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• optic







Defining a Model (i)

- Define the "variables" (cf. x in the theory example from the xtra-lite triple):
- allvars = {
 - » x1,y1,z1,yaw1,pitch1,roll1,
 - » x2,y2,z2,yaw2,pitch2,roll2,
 - » x3,y3,z3,yaw3,pitch3,roll3

```
};
```

• Define the "floats" (cf. **q** in the theory):

```
»allfloats = {
    -qul,qur,qlf,qlb,qrf,qrb
»};
```

• Define the "parameters" (cf. **s** in the theory):

```
allparams = {
    x00, y00, z00, yaw00, pitch00, roll00
    };
```



Defining a Model (ii)

- Define coordinate lists for rigid bodies of interest:
- optic = {x3, y3, z3, yaw3, pitch3, roll3};
- support = {x00, y00, z00, yaw00, pitch00, roll00};

• Define coordinate lists for points on rigid bodies

massUl={0,-n1,d0}; (* left wire attachment point on upper mass *)

• Define list of gravitational potential terms:

- gravlist = {}; (* initialize list *)
- AppendTo[gravlist, m3 g z3]; (* typical item *)



Defining a Model (iii)

• Define list of wires, each with the following format

- }
 - » coordinate list defining first mass,
 - » attachment point for first mass (local coordinates),
 - » attachment vector for first mass,
 - » coordinate list defining second mass,
 - » attachment point for second mass (local coordinates),
 - » attachment vector for second mass,
 - » Young's modulus,
 - » unstretched length,
 - » longitudinal elasticity,
 - » vector defining principal axis 1,
 - » moment of area along principal axis 1,
 - » moment of area along principal axis 2,
 - » linear elasticity type,
 - » angular elasticity type,
 - » torsional elasticity type,
 - » shear modulus,
 - » cross sectional area for torsional calculations,
 - » torsional stiffness geometric factor

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Defining a Model (iv)

• Define list of springs, each with following format:

• {

- » coordinate list defining first mass,
- » attachment point for first mass (local coordinates),
- » attachment angles for first mass (yaw, pitch, roll),
- » coordinate list defining second mass,
- » attachment point for second mass (local coordinates),
- » attachment angles for second mass (yaw, pitch, roll),
- » damping type,
- » 6x6 elasticity matrix,
- » 1*6 pre-load force/torque vector
- }

Define kinetic energy

```
    IM3 = {{I3x, 0, 0}, {0, I3y, 0}, {0, 0, I3z}}; (* typical MOI tensor)
    kinetic = (
    ...
    +(1/2) m3 Plus@@(Dt[b2s[optic,COM],t]^2)
    +(1/2) omegaB[yaw3, pitch3, roll3].IM3.omegaB[yaw3, pitch3, roll3]
```

-);

```
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```

•••



Defining a Model (v)

- Define default values of constants
- defaultvalues = {
 - » g -> 9.81, (* value given numerically *)
 - » ...
 - » m3 -> Pi*r3^2*t3, (* value given in terms of other constants *)
 - » ····
 - » x00 -> 0, (* value for nominal position of structure *)
 - » y00 -> 0,
 - » z00 -> 0,

••••

- » ...
- » damping[imag,dampingtype] -> (phi&) (* value for frequency dependence of damping *)
- **
- };

Define starting point for finding equilibrium position:



Defining a Model (vi)

- Define model-specific utilities:
 - » A function to list eigenmodes in a table
 - » pretty[eigenvector]
 - » A function to plot eigenmode shapes
 - » eigenplot[eigenvector, amplitude, {viewpoint}]
 - » Vectors representing force and displacement inputs and displacement outputs of interest

```
» structurerollinput = makeinputvector[roll00];
```

```
» opticxinput = makefinputvector[x3];
```

```
» opticx = makeoutputvector[x3];
```

» Rotation matrices to put angle variables in a more easily interpretable basis:

```
» e2ni;
```



Sample Output (i)

 Transfer function from x displacement of support to x motion of optic (quad model, reference parameters of 20031114):



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Sample Output (ii)

 Thermal noise in x motion of optic (quad model, reference parameters of 20031114):

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Export to Matlab/Simulink

Export to E2E

Status

- Toolkit and quad and triple models defined and published:
 - » http://www.ligo.caltech.edu/~mbarton/SUSmodels/indexMB.html
 - » T020205-01
- Export of state-space matrices to E2E programmed.
- Simple E2E model using triple data is working.
- Need to do similar test model for quad data.
- Need to add local control systems.
- Need to define E2E blocks for all pendulum designs.
- Need to add global control (when finalized ages off yet).
- Need to start on SEI.