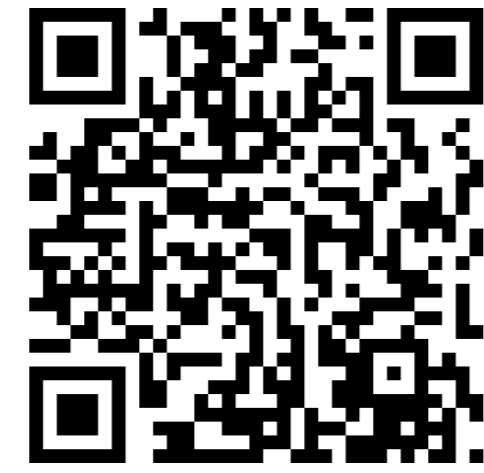




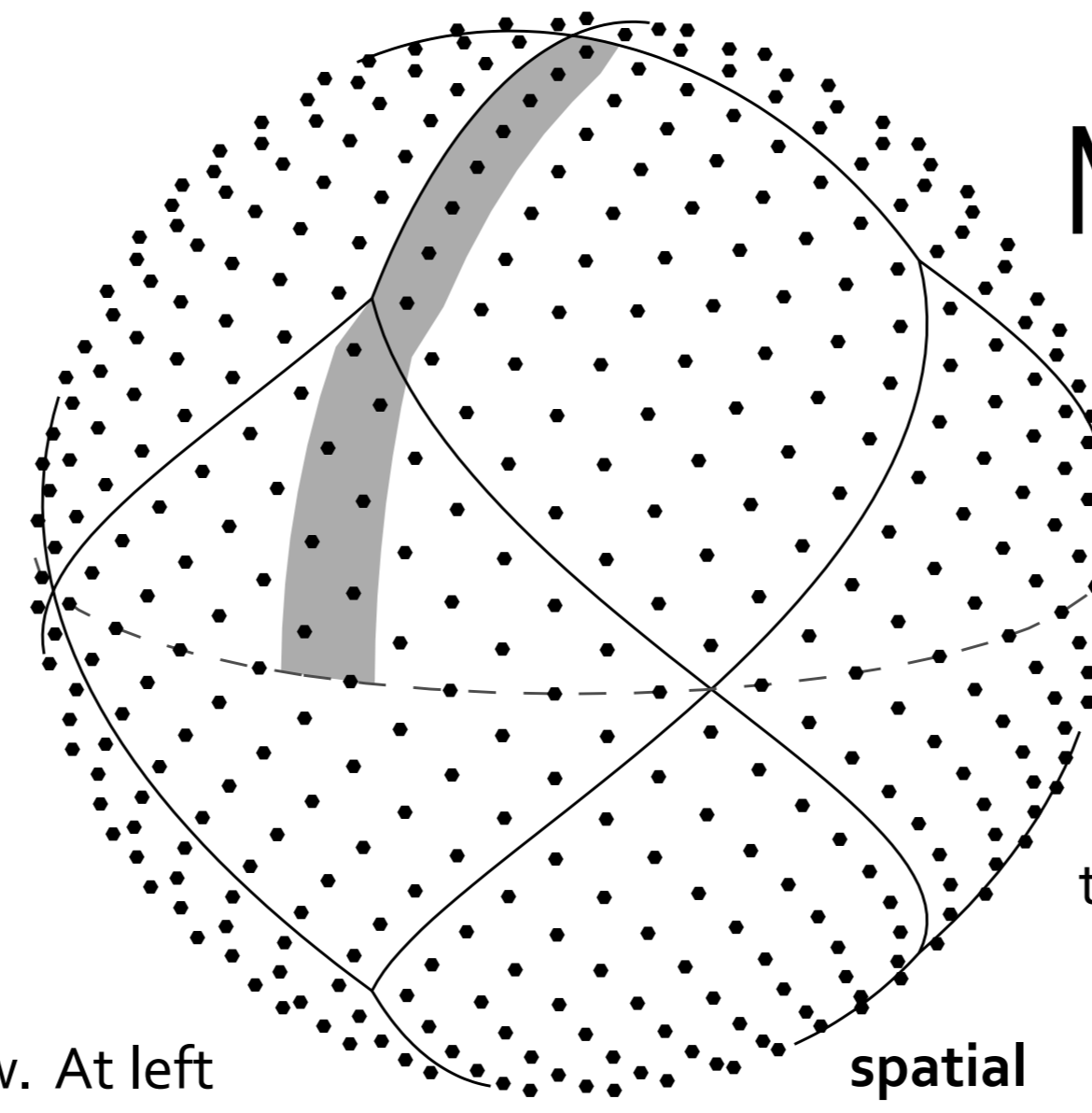
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LSC-VIRGO FALL 2011

figure 6d of Górski et al. (2005)
Healpix, 768 pixels



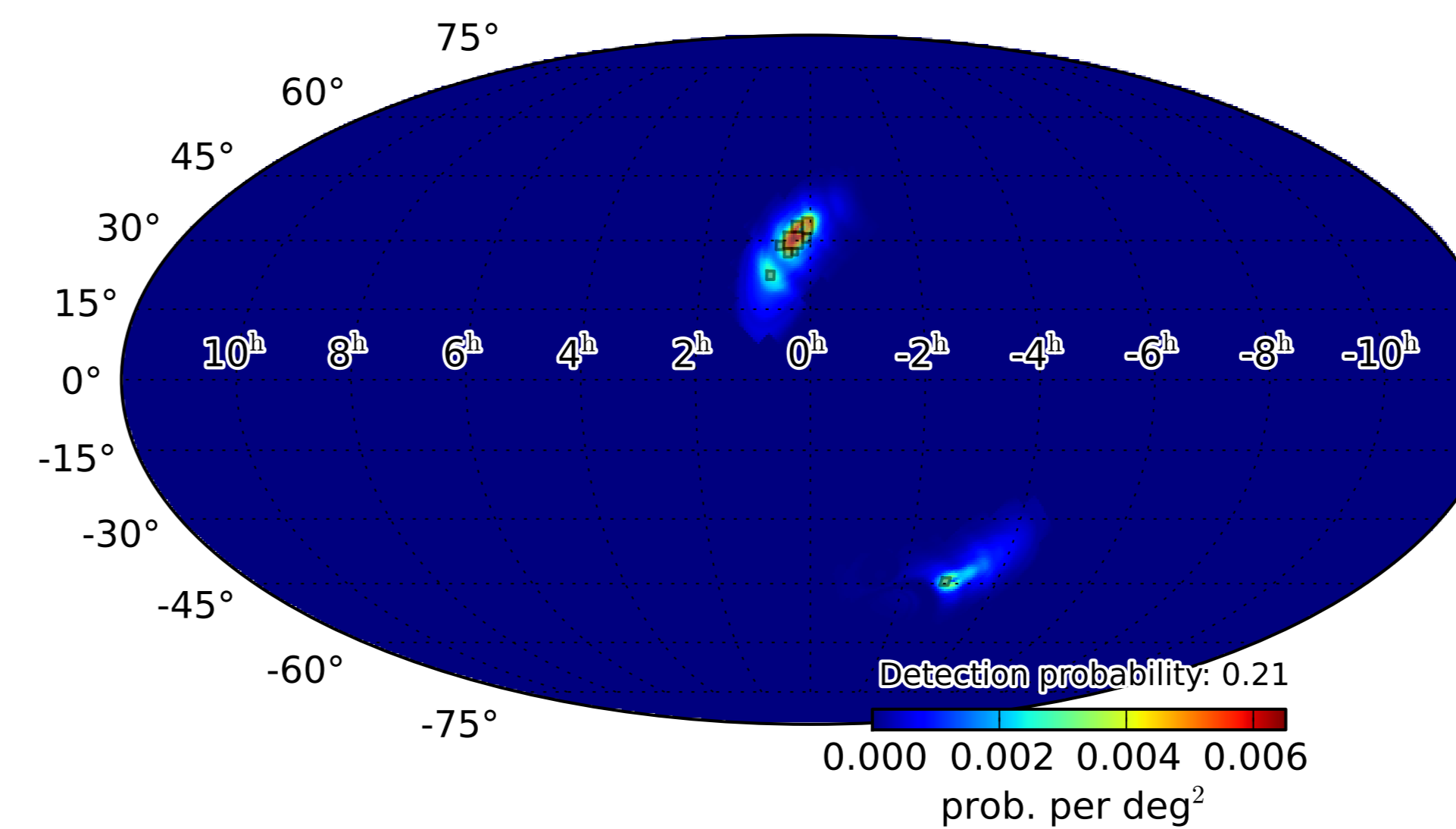
Numerical implementation

Equation (1) resembles a cross-correlation integral. For scalar functions, the convolution theorem and the fast Fourier transform (FFT) make cross-correlation in the frequency domain very efficient. Driscoll & Healy (1994), Wandelt & Górski (2001), and others have described analogous fast convolution algorithms for functions defined on the unit sphere. We used *HEALPix* to write two convolution algorithms:

spatial uses nearest neighbor interpolation of the rotated kernel to approximate the integral in spherical polar coordinates.

multipole involves a spherical harmonic transform of both the masked sky map $v_i(\omega)s(\omega)$ and the FOV $b_i(\omega)$, a weighted inner product of the spherical harmonic coefficients, and a 2D inverse FFT to return to polar coordinates.

We checked convergence and run time of both algorithms. At a resolution of $\approx 0.05 \text{ deg}^2$, **spatial** takes $\approx 1000 \text{ s}$ while **multipole** takes $\approx 25 \text{ s}$ of CPU time. Both versions are *OpenMP* accelerated to exploit multiple cores. On the LIGO-Caltech cluster head node, the **multipole** algorithm has achieved run times as short as 5 s, though further speedup is possible.



Multiple telescope case

For multiple telescopes, the figure of merit is the probability of imaging the source with at least one telescope:

$$p_{\geq 1} \equiv p(\text{EM}_1 \cup \text{EM}_2 \cup \dots \cup \text{EM}_N | \gamma_1, \gamma_2, \dots, \gamma_N, \text{GW}) = 1 - \int [1 - b_1(\gamma_1^{-1}\omega) v_1(\omega)] [1 - b_2(\gamma_2^{-1}\omega) v_2(\omega)] \dots [1 - b_N(\gamma_N^{-1}\omega) v_N(\omega)] s(\omega) d\Omega. \quad (2)$$

The integral expression above can be evaluated iteratively one telescope at a time, suggesting a greedy algorithm depicted in the flow chart at right.

The most sophisticated planning algorithm that we tried used simulated annealing to simultaneously vary the pointings of all of the telescopes. We used the Python module `scipy.optimize.anneal` and a modified version of the "very fast" cooling schedule described by Ingber (1989).

Case study

We compared the detection efficiency of the noncoordinated planner and our two coordinated planners using a set of 2126 GW skymaps from low-mass inspiral signals injected into 24 hours of simulated initial LIGO noise.

We computed the environmental masks $v_i(\omega)$ using the Python edition of *NOVAS*, the US Naval Observatory's positional astronomy library.

Surprisingly, we found that (a) *greedy* performed almost as well as *anneal* in terms of detection efficiency, and (b) both coordinated planners had roughly *double* the detection efficiency of the *noncoordinated* planner.

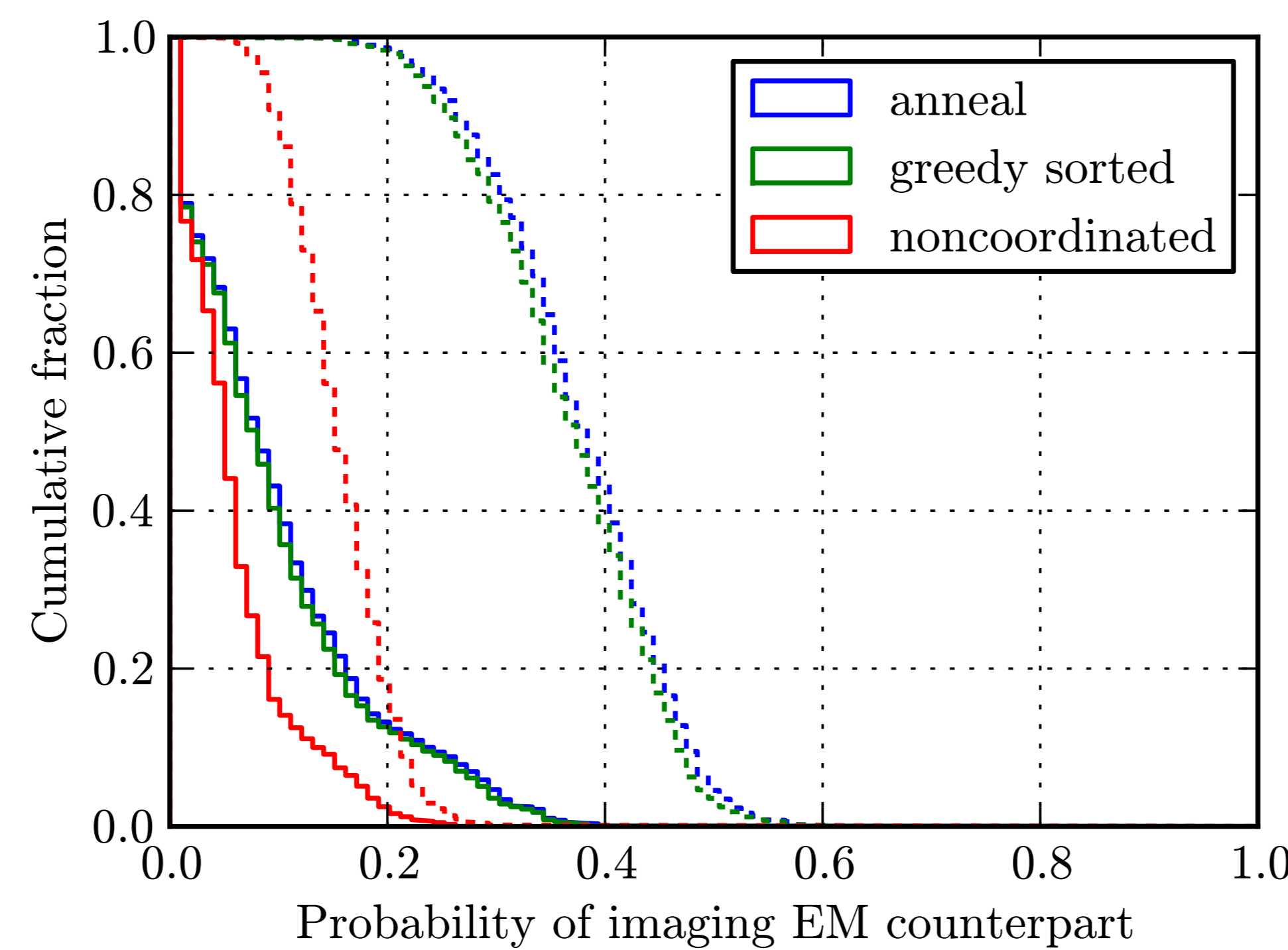
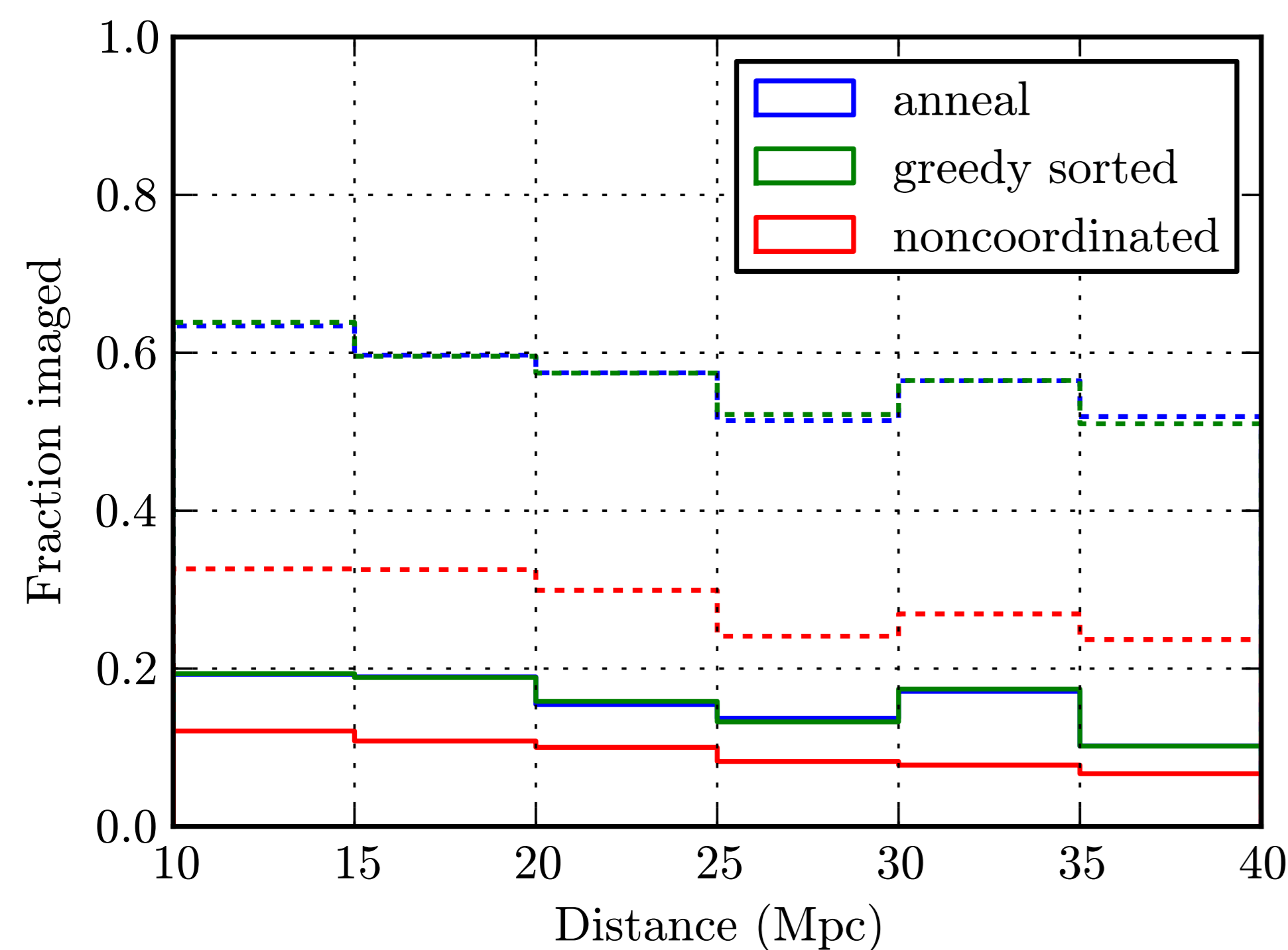
Our key results are the two figures below. At left is the fraction of injected signals that we would have imaged with one pointing of each telescope at the time of the trigger as a function of luminosity distance. Solid lines represent observing plans that account for interference from the Sun and Earth. Dashed lines represent observing plans in which these considerations are neglected.

We tested three different planning algorithms:

noncooperative Each telescope is independently pointed where it is most likely to observe an EM counterpart.

greedy sorted Suppose that we have chosen pointings for telescopes 1, 2, ..., i . The pointing of telescope $i + 1$ is chosen to maximize the probability of detection, subject to telescopes 1, 2, ..., i , remaining fixed.

anneal Uses simulated annealing to the probability of imaging the source by varying the configurations of all of the telescopes simultaneously.



Single telescope case

The posterior distribution of source location ω given all GW observations is commonly called the GW sky map, denoted $p(\omega | \text{GW})$.

Let EM_i denote the event of observing an optical transient with telescope i . Let γ_i represent the pointing of telescope i . The probability of observing an EM counterpart in telescope i given its pointing γ_i is

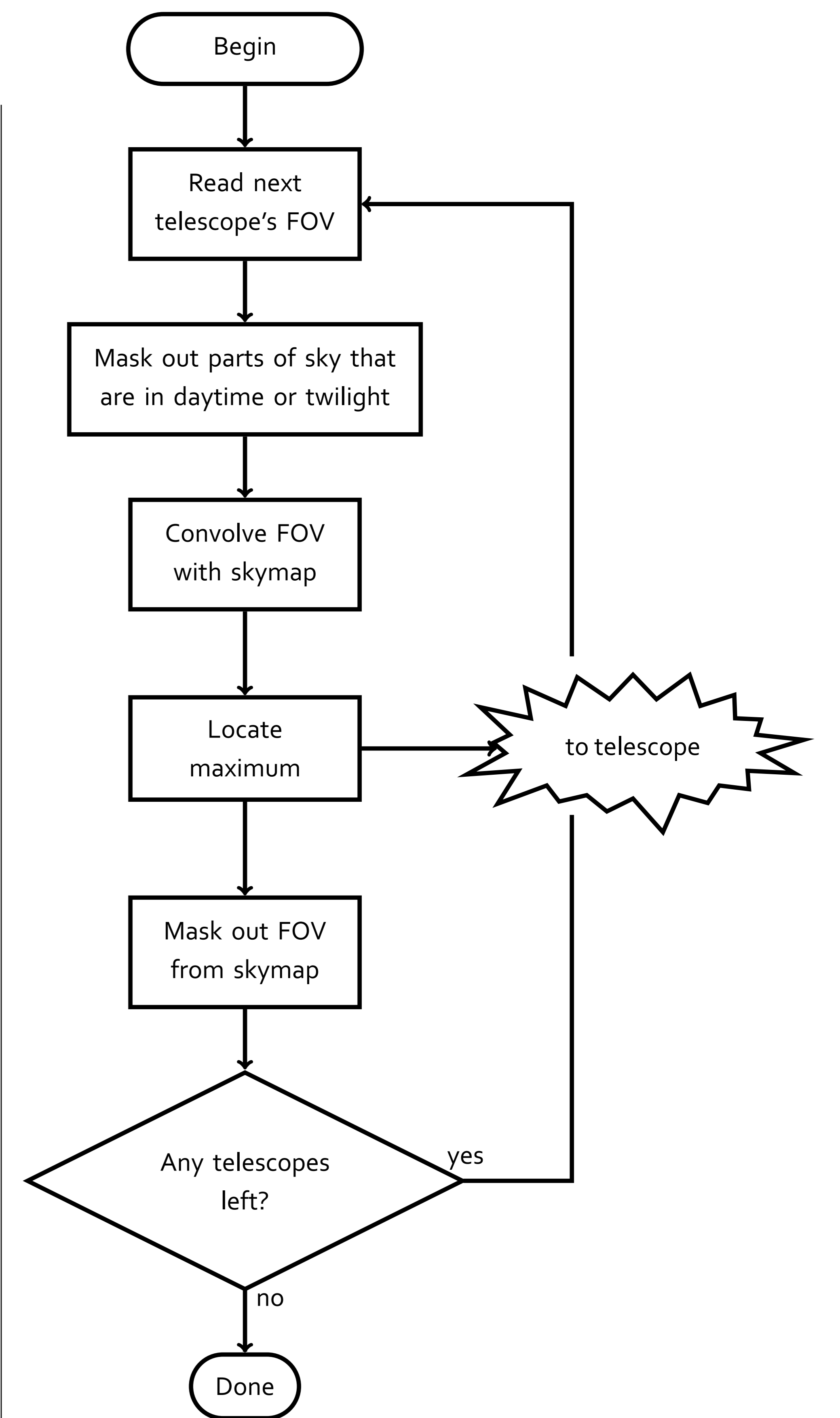
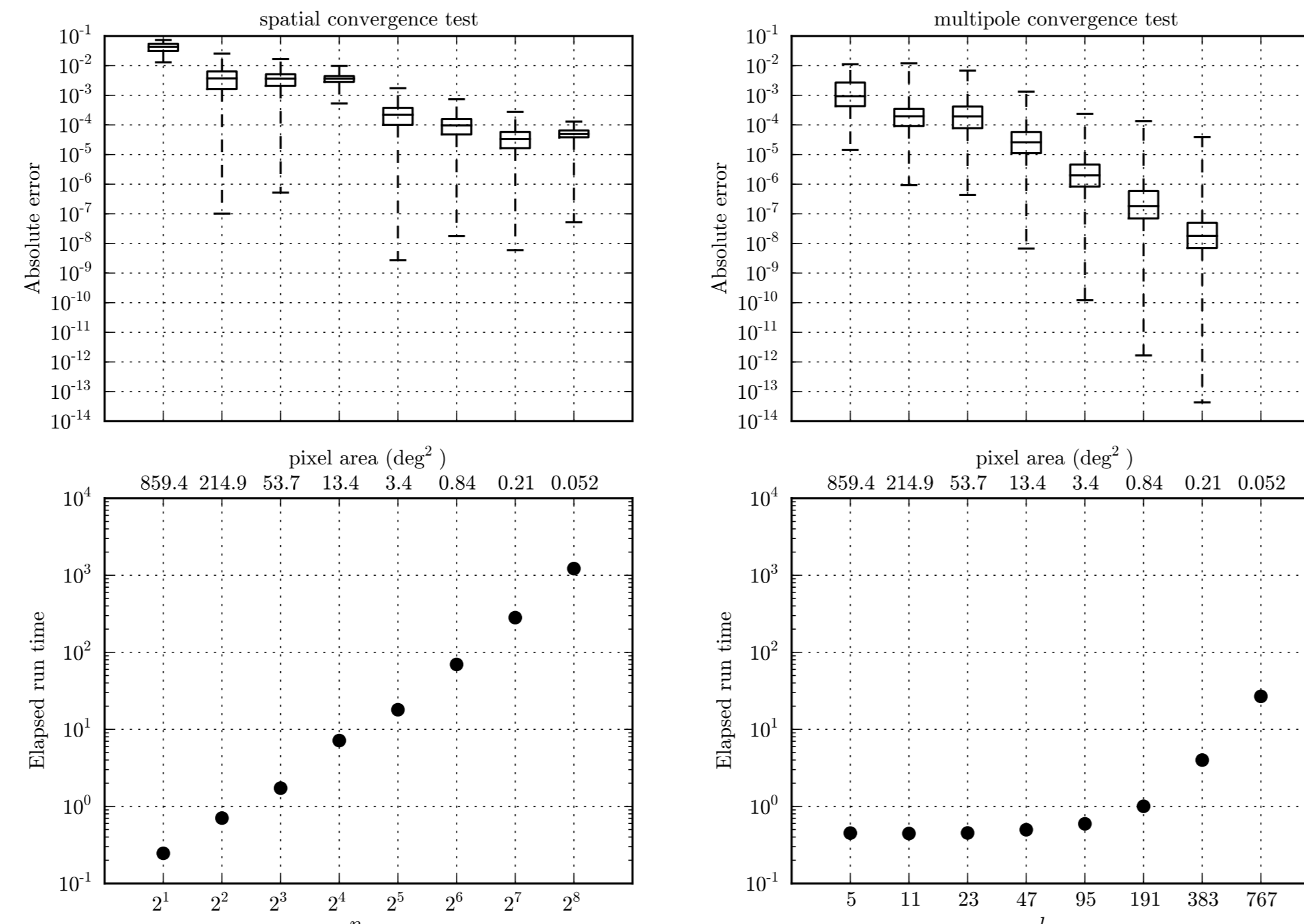
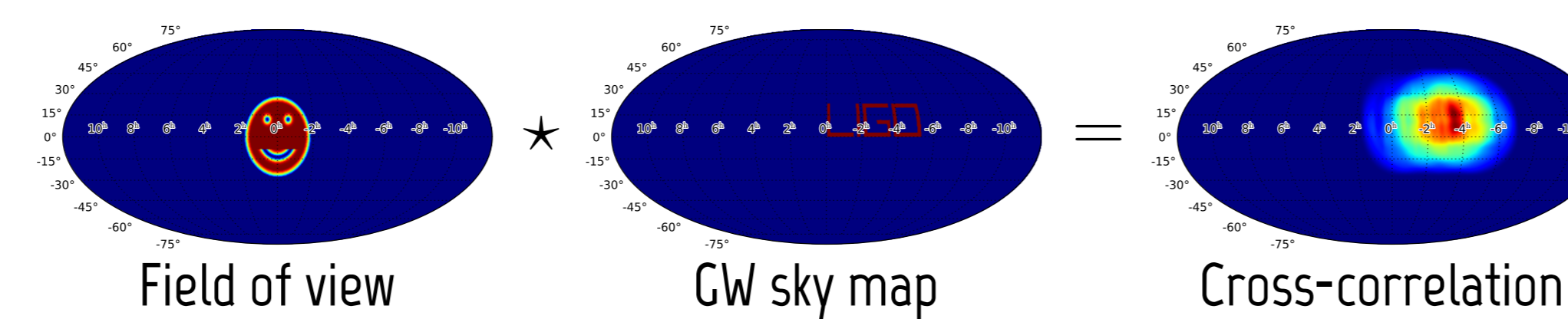
$$p(\text{EM}_i | \gamma_i, \omega) \equiv b_i(\gamma_i^{-1}\omega) v_i(\omega).$$

Here, the function $b_i(\gamma_i^{-1}\omega)$ describes the telescope's (rotated) FOV and v_i describes environmental features such as the twilight/nighttime terminator, the horizon, and optionally the seeing. Now, marginalizing over the unknown source location, this becomes

$$p(\text{EM}_i | \gamma_i, \text{GW}) \equiv \int b_i(\gamma_i^{-1}\omega) v_i(\omega) s(\omega) d\Omega. \quad (1)$$

The optimal pointing is

$$\gamma_i^* \equiv \arg \max_{\gamma_i} p(\text{EM}_i | \gamma_i, \text{GW}).$$



Conclusion

We have demonstrated that coordinating observations amongst multiple telescopes is a promising strategy for EM followup of GW events. We have developed an observation planning code that is *fast* enough to be used for extensive simulation campaigns and *flexible* enough to accommodate any network of telescopes. Our code aims to be *scalable* enough to produce observing plans in near real-time on a multi-core machine.

This project lays the groundwork for future multi-messenger studies that will account for a mix of telescopes with different limiting magnitudes, slew times, in addition to fields of view.

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