

# Hough search for continuous gravitational waves using LIGO S4 data

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Illes Balears





# Content



- Introduction
- The Hough transform
  - Overview of S2 analysis
  - Improvements for S4
  - Preliminary S4 results
- Future work

Search can detect any periodic source.

Upper limits are set on gravitational-wave amplitude,  $h_0$ , of rotating triaxial ellipsoid. →

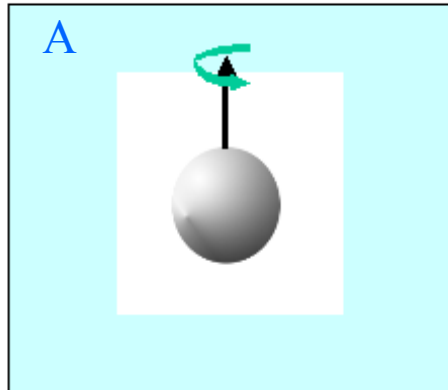
Credits:

A. image by Jolien Creighton; LIGO Lab Document G030163-03-Z.

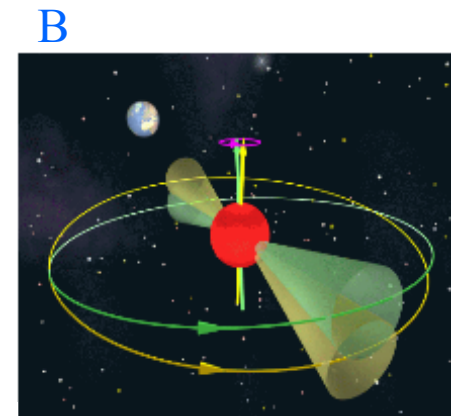
B. image by M. Kramer; Press Release PR0003, University of Manchester - Jodrell Bank Observatory, 2 August 2000.

C. image by Dana Berry/NASA; NASA News Release posted July 2, 2003 on Spaceflight Now.

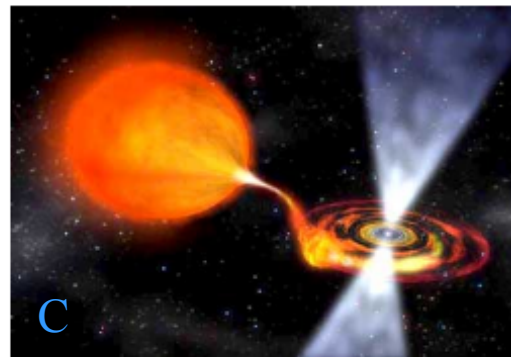
D. image from a simulation by Chad Hanna and Benjamin Owen; B. J. Owen's research page, Penn State University.



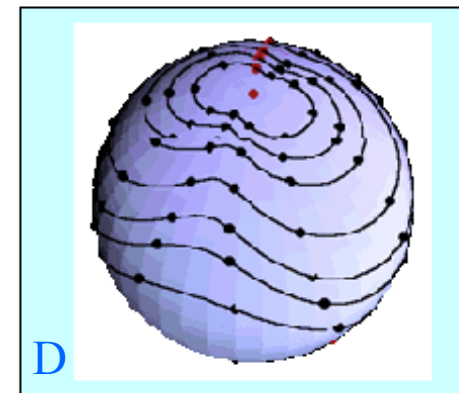
**Mountain on neutron star**



**Precessing neutron star**



**Accreting neutron star**



**Oscillating neutron star**

- Expected waveform from an isolated spinning NS is sinusoidal with small spin-down:

$$h(t) = F_+(t, \psi)h_+(t) + F_\times(t, \psi)h_\times(t)$$

$$h_+ = A_+ \cos \Phi(t)$$

$$h_\times = A_\times \sin \Phi(t)$$

$$\Phi(t) = \phi_0 + 2\pi \sum_{n=0}^{\infty} \frac{f_{(n)}}{(n+1)!} (t-t_0)^{n+1}$$

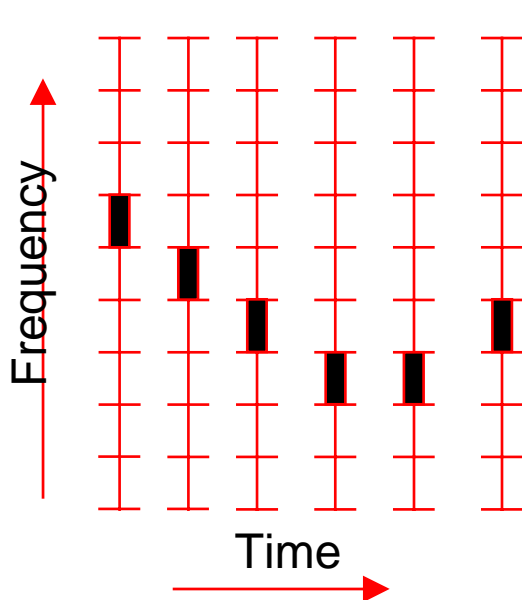
- Doppler frequency modulation due to motion of Earth and amplitude modulation due to detector antenna pattern.
- For setting upper limits only, we assume the emission mechanism is due to deviations of the pulsar's shape from perfect axial symmetry,  $f_{GW} = 2f_r$

$$A_+ = \frac{1}{2} h_0 (1 + \cos^2 \iota)$$

$$A_\times = h_0 \cos \iota$$

$$h_0 = \frac{16\pi^2 G}{c^4} \frac{I_{zz} \varepsilon f_r^2}{d}$$

- Three methods have been developed to search for cumulative excess power from a hypothetical periodic gravitational wave signal by examining successive spectral estimates:



**StackSlide**

**PowerFlux**

**Hough**

They are all based on breaking up the data into segments, FFT each, producing Short (30 min) Fourier Transforms (SFTs) from  $h(t)$ , as a coherent step (although other coherent integrations can be used if one increasing the length of the segments), and then track the frequency drifts due to Doppler modulations and  $df/dt$  as the incoherent step.

- Other fully coherent methods:
  - Frequency domain match filtering/maximum likelihood estimation
  - Time domain Bayesian parameter estimation



## Differences among the incoherent methods



What is exactly summed?

- **StackSlide** – Normalized power (power divided by estimated noise)  
→ Averaging gives expectation of 1.0 in absence of signal
- **Hough** – Weighted binary counts (0/1 = normalized power below/above SNR), with weighting based on antenna pattern and detector noise
- **PowerFlux** – Average strain power with weighting based on antenna pattern and detector noise  
→ Signal estimator is direct excess strain noise  
(circular polarization and 4 linear polarization projections)



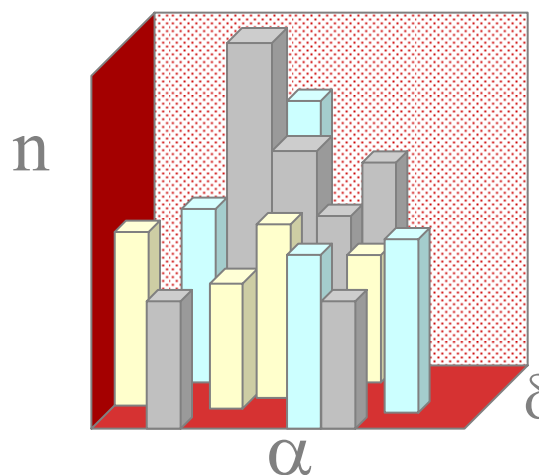
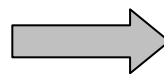
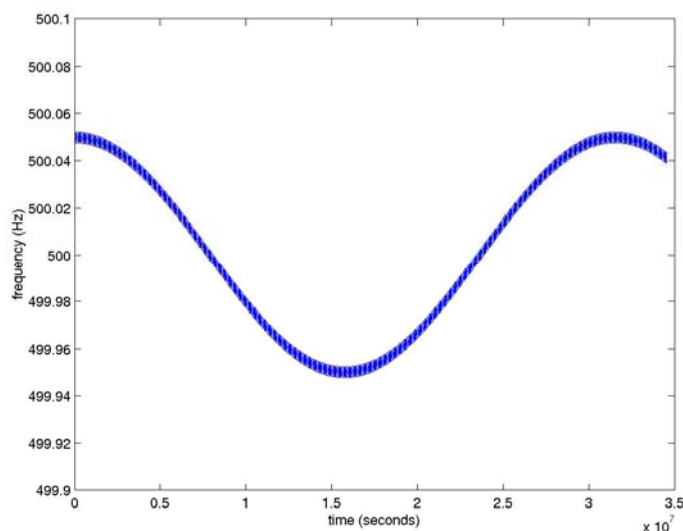
The *Hough transform* is a general method for pattern recognition that was developed and patented many years ago.

We use the *Hough transform* to find a pattern produced by the Doppler modulation & spin-down of a GW signal in the time-frequency plane of our data.

For isolated NS the expected pattern depends on:  $\{\alpha, \delta, f_0, \dot{f}_n\}$

$$f(t) - \hat{f}(t) = \hat{f}(t) \frac{\mathbf{v}(t) \cdot \mathbf{n}}{c}$$

$$\hat{f}(t) = f_0 + \dot{f}(t - t_0)$$





# Review of S2 Hough analysis



- Start with 1800s SFTs for each detector
- Select frequency bins by setting a threshold on normalized power  
– gives time-frequency collection of 0s and 1s
- For N SFTs, the final number count for a given parameter space point is

$$n = \sum_{i=1}^N n_i \quad \text{where} \quad n_i = \begin{cases} 0 \\ 1 \end{cases}$$

- Using 0s and 1s leads to gain in computational efficiency and it is more robust against large transient power artifacts



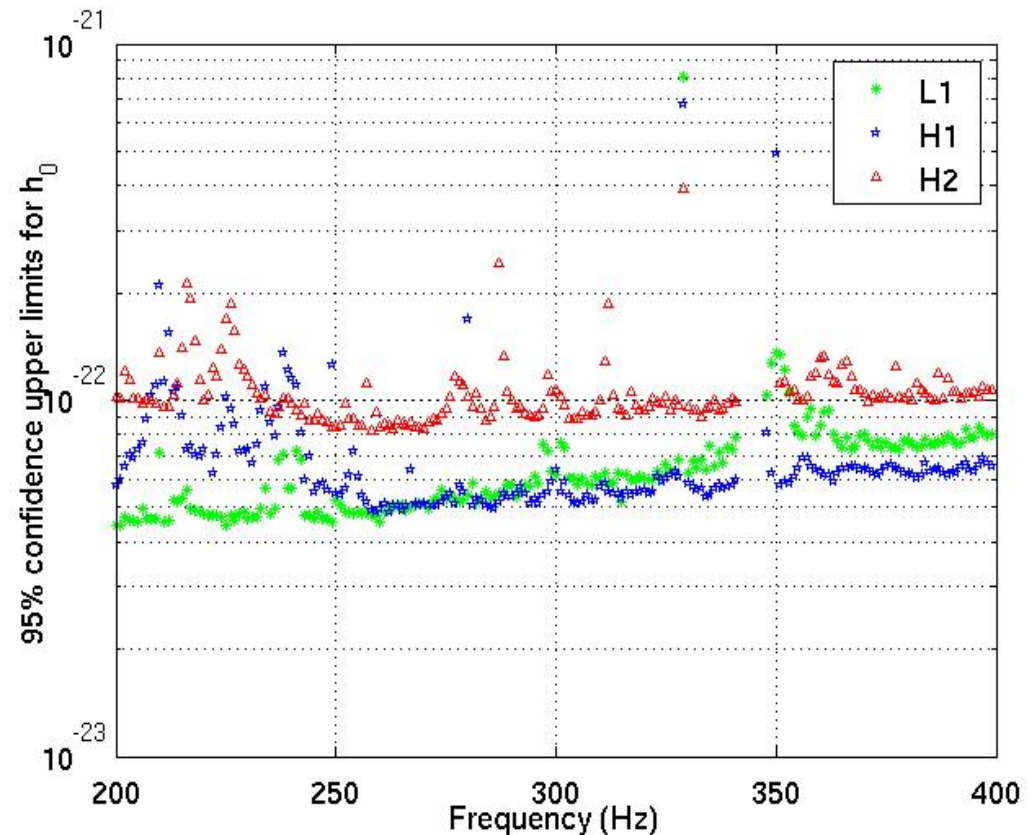


# Hough S2: UL Summary

Feb.14-Apr.14,2003



- S2 analysis covered 200-400Hz, over the whole sky, and 11 values of the first spindown ( $\Delta f = 5.55 \times 10^{-4}$  Hz,  $\Delta f_1 = -1.1 \times 10^{-10}$  Hz s<sup>-1</sup>)
- Templates: Number of sky point templates scales like (frequency)<sup>2</sup>
  - $1.5 \times 10^5$  sky locations @ 300 Hz
  - $1.9 \times 10^9$  @ 200-201 Hz
  - $7.5 \times 10^9$  @ 399-400 Hz
- Three IFOs analyzed separately
- No signal detected
- Upper limits obtained for each 1 Hz band by signal injections: Population-based frequentist limits on  $h_0$  averaging over sky location and pulsar orientation



Detector	L1	H1	H2
Frequency (Hz)	200-201	259-260	258-259
$h_0^{95\%}$	$4.43 \times 10^{-23}$	$4.88 \times 10^{-23}$	$8.32 \times 10^{-23}$

- Perform the Hough transform for a set of points in parameter space  $\lambda = \{\alpha, \delta, f_0, f_i\} \in \mathbf{S}$ , given the data:

$$\mathbf{HT}: \mathbf{S} \rightarrow \mathbf{N}$$

$$\lambda \rightarrow n(\lambda)$$

- Determine the maximum number count  $n^*$

$$n^* = \max (n(\lambda)): \lambda \in \mathbf{S}$$

- Determine the probability distribution  $p(n/h_0)$  for a range of  $h_0$
- The 95% frequentist upper limit  $h_0^{95\%}$  is the value such that for repeated trials with a signal  $h_0 \geq h_0^{95\%}$ , we would obtain  $n \geq n^*$  more than 95% of the time

$$0.95 = \sum_{n=n^*}^N p(n/h_0^{95\%})$$

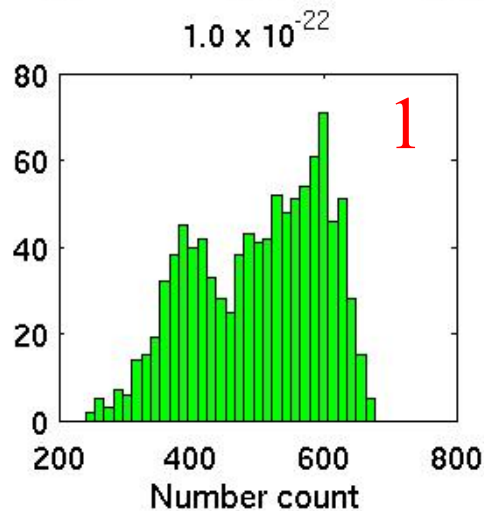
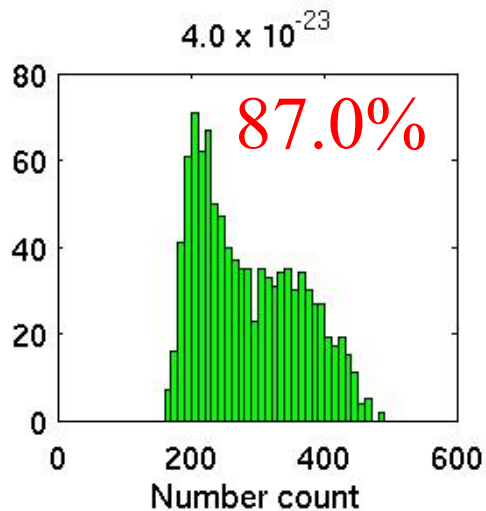
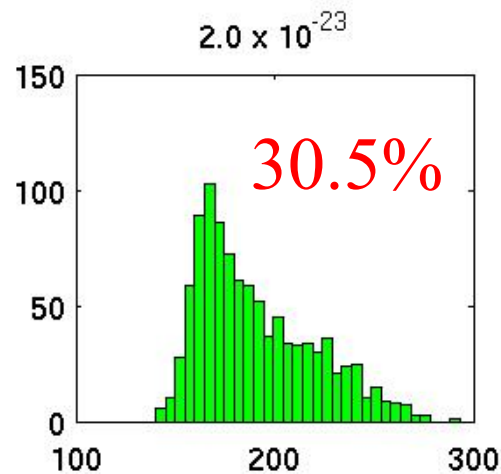
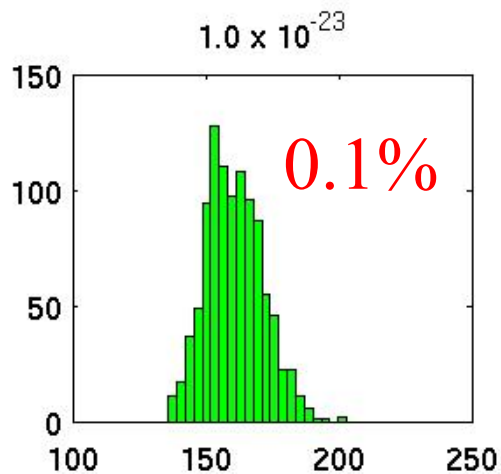
Compute  $p(n/h_0)$  via Monte Carlo signal injections



# Number count distribution for signal injections in S2 data



L1: 200-201 Hz,  $n^* = 202$ , 1000 injections



$p(n/h_0)$  ideally binomial for a target search, but:

- Non stationarity in the noise
- Amplitude modulation of the signal for different SFTs
- Different sensitivity for different sky locations and pulsar orientations
- Random mismatch between signal & templates

‘smear’ out the binomial distributions



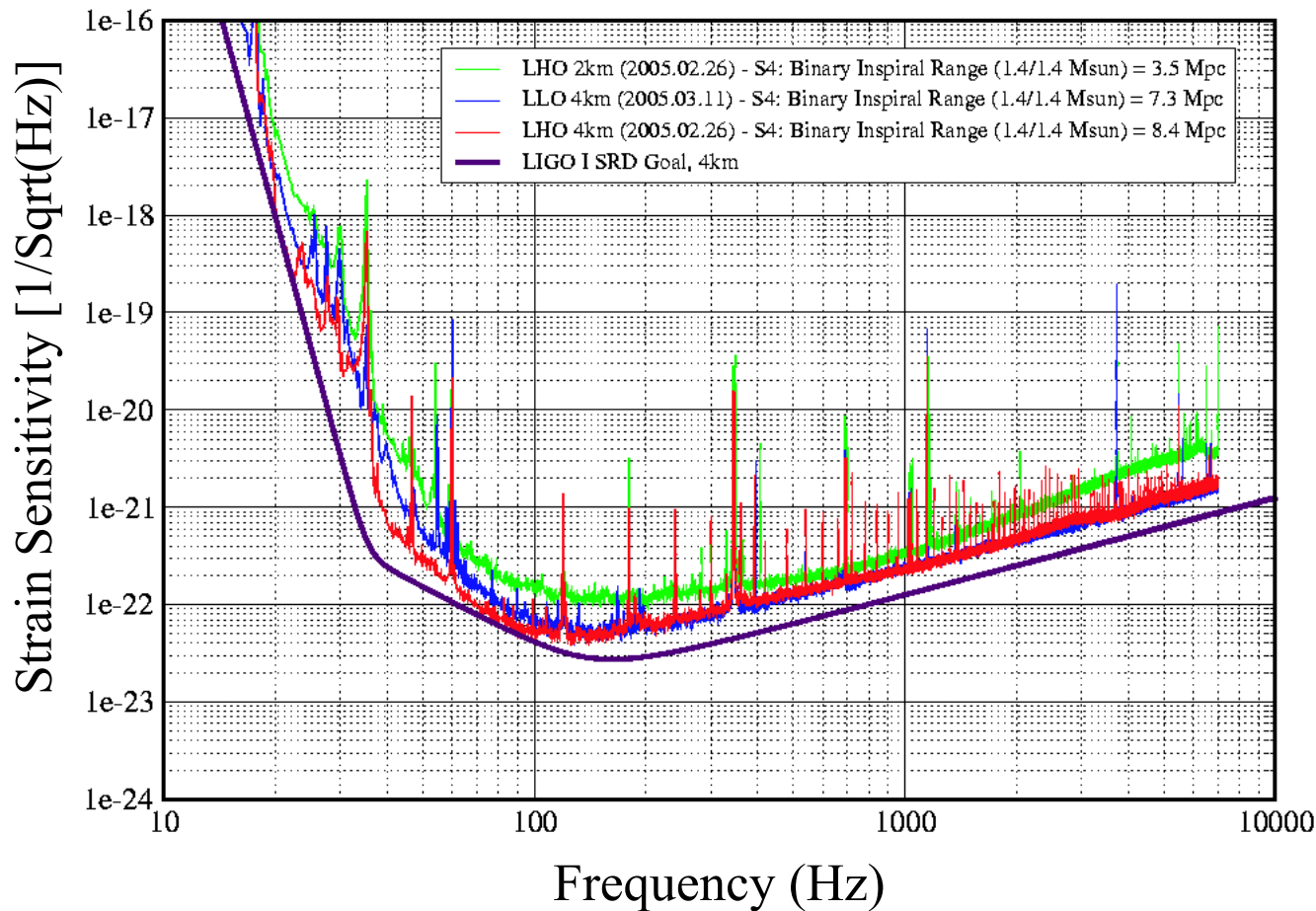
# Fourth Science Run (S4) Sensitivity

(February 22, 2005 – March 23, 2005)



## Strain Sensivities for the LIGO Interferometers

Best Performance for S4 LIGO-G050230-02-E



S5 is currently running at design sensitivity!



## Improvements for S4



- The S2 Hough search has been modified to take into account that the SFTs have different noise floors and the signal amplitude changes in time – SNR changes across SFTs
- We use a *weighted Hough* to give more weight to SFTs having greater SNR. Weights are proportional to the beam pattern functions and inversely proportional to the SFT noise floor.
- Weighting method applied to Hough was initially suggested by C.Palomba and S.Frasca at GWDAW-2004 and it is similar to the one used by the PowerFlux method.
- Number count  $n$  is not an integer anymore

$$n = \sum_{i=1}^N w_i n_i \quad \sum_{i=1}^N w_i = N$$

- Using the weights does not lead to any loss in computational efficiency or robustness
- It has also been generalized to the Multi-IFO case

- Nominal sensitivity for given FA and FD assuming a perfectly matched template averaged over sky, orientations and polarization angles:
  
- Improved Sensitivity:  
Assumes template is perfectly matched to signal, and average over all pulsar orientations and polarization angles (but not over sky-positions)
  
- Optimal choice of weights is:
  
- Optimally weights should be calculated at same sky-location as signal

$$h_0 = 5.34 \frac{S^{1/2}}{N^{1/4}} \sqrt{\frac{S_n}{T_{coh}}}$$

$$S = \text{erfc}^{-1}(2\alpha_H) + \text{erfc}^{-1}(2\beta_H)$$

$$h_0 = 3.38 S^{1/2} \left( \frac{\|\vec{w}\|}{\vec{w} \cdot \vec{X}} \right)^{1/2} \sqrt{\frac{\langle S_n^{(i)} \rangle}{T_{coh}}}$$

$$X_i = \langle S_n^{(i)} \rangle \frac{(F_+^{(i)})^2 + (F_\times^{(i)})^2}{S_n^{(i)}}$$

$$w_i \propto \frac{(F_+^{(i)})^2 + (F_\times^{(i)})^2}{S_n^{(i)}}$$

- Gain in sensitivity is large if standard deviation of SFT noise floors is large or if signal amplitude changes rapidly across SFTs
- Mean number count is unchanged due to normalization of weights:

$$\langle n \rangle = N\alpha = N \exp(-\rho_{th})$$

- Standard deviation always increases:

$$\sigma_n = \|\vec{w}\| \sqrt{\alpha(1-\alpha)}$$

- Number count threshold for a given false alarm:

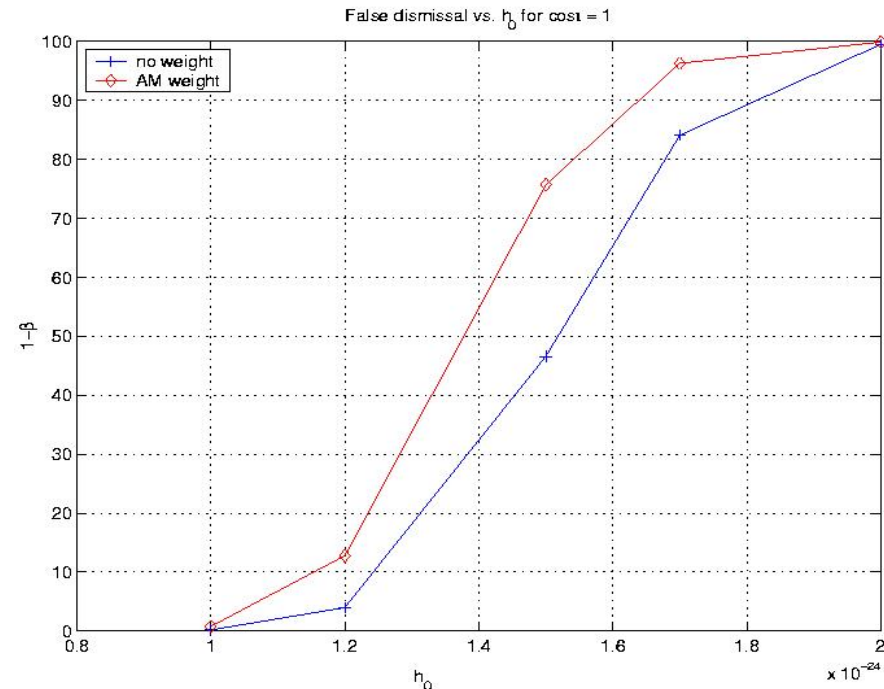
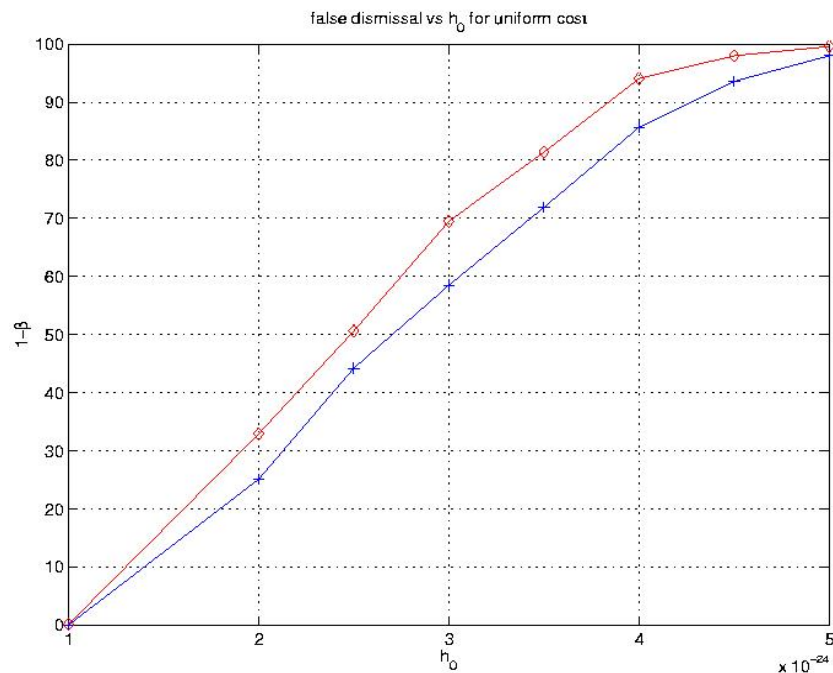
$$n_{th} = N\alpha + \sqrt{2\|\vec{w}\|^2 \alpha(1-\alpha)} \operatorname{erfc}^{-1}(2\alpha_H)$$



# Improvement in detection efficiency



Signal injections in fake data, 250-260Hz, random sky-position and polarization angles.  
Number count threshold set for  $\alpha_H=10^{-10}$



- Improvement in sensitivity at 90% efficiency is roughly 10% in signal amplitude for a perfectly matched template and stationary noise. The gain depends on pulsar orientation
- Will be somewhat degraded when searching in a sky-patch because of a mismatch and also because we will use a single set of weights for the whole sky-patch (calculated at the center), but sensitivity can also improve in case of non-stationary noise.





## The S4 Hough search



- As before, input data is a set of  $N$  1800s SFTs (no demodulations)
- Weights allow us to use SFTs from all three IFOs together: 1004 SFTS from H1, 1063 from H2 and 899 from L1
- Search frequency band 50-1000Hz
- 1 spin-down parameter. Spindown range  $[-2.2, 0] \times 10^{-9}$  Hz/s with a resolution of  $2.2 \times 10^{-10}$  Hz/s
- All sky search
- Sky is broken up into 92 patches  $0.4 \text{ rad} \times 0.4 \text{ rad}$  wide
- Line cleaning used to remove known narrow spectral lines
- All-sky upper limits set in 0.25 Hz bands
- Multi-IFO and single IFOs have been analyzed

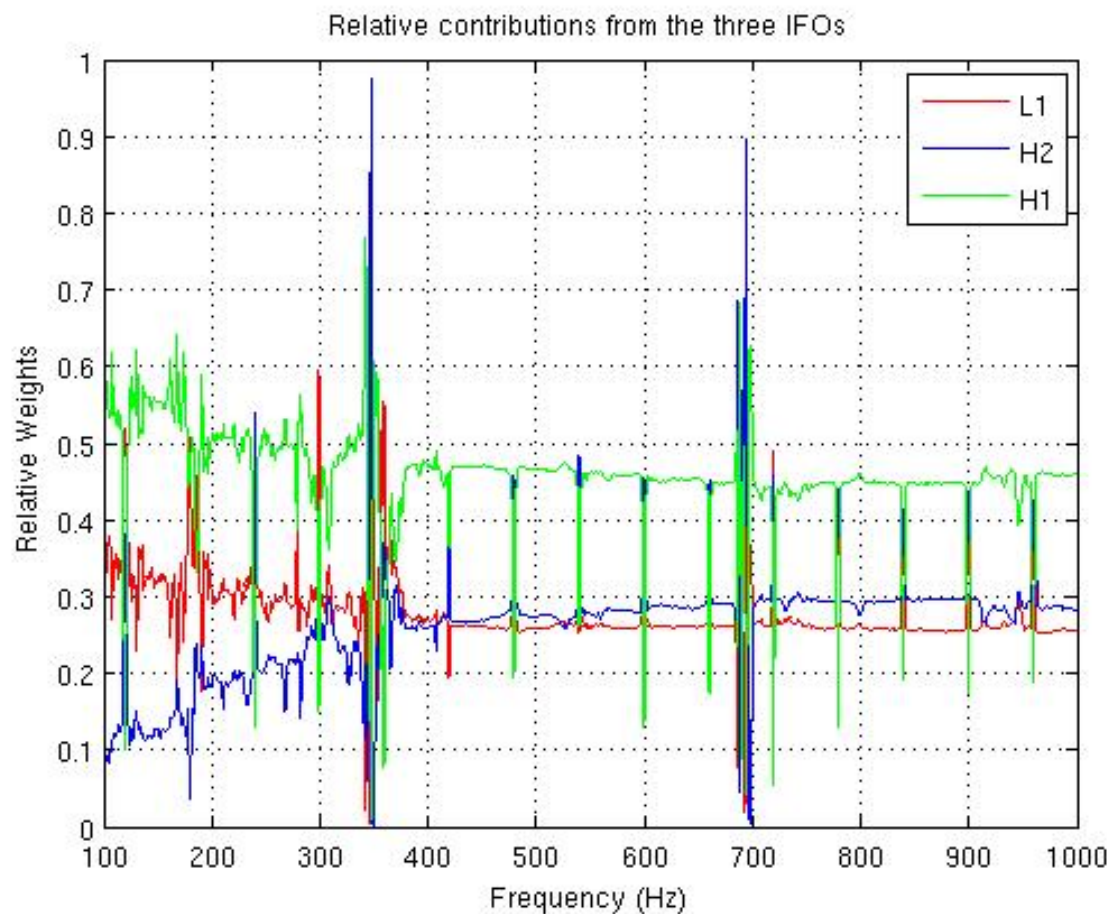


Figure plots the relative noise weights from H1, H2 and L1

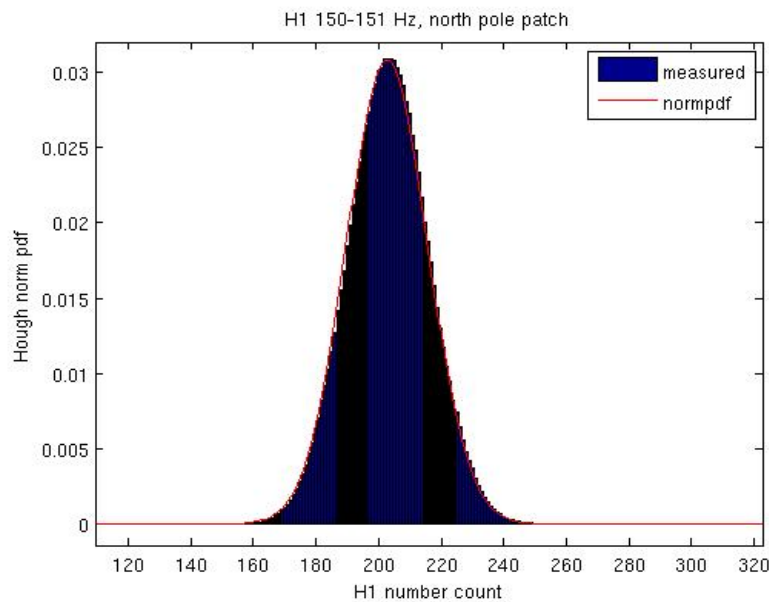
$$\frac{\sum_{IFO} w_i}{\sum_{All} w_i}$$



# Histogram of Hough number counts for the H1 detector

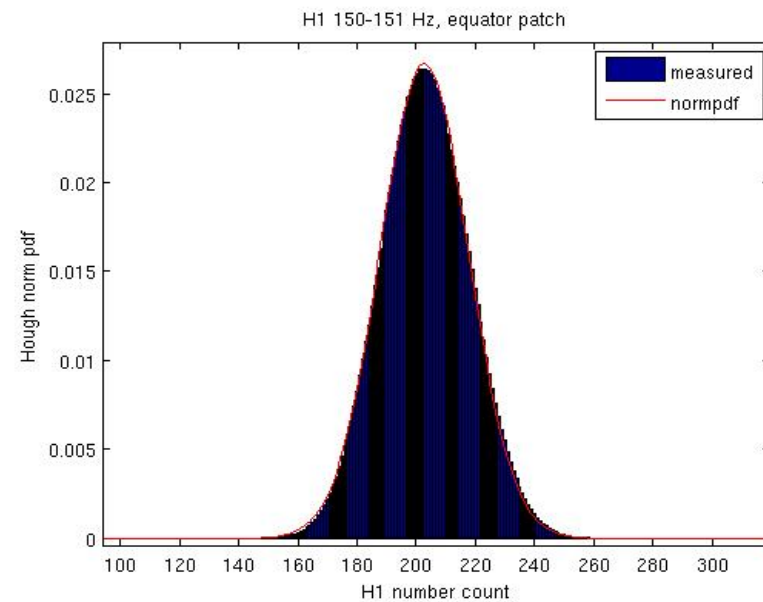


Histogram of the Hough number count compared to a Gaussian distribution for the H1 detector (1004 SFTs) in the frequency band 150-151 Hz. Number of templates analyzed in each sky patch  $\sim 11 \times 10^6$



$$\langle n \rangle = 202.7$$

$$\sigma = 12.94 \text{ (obtained from the weights)}$$

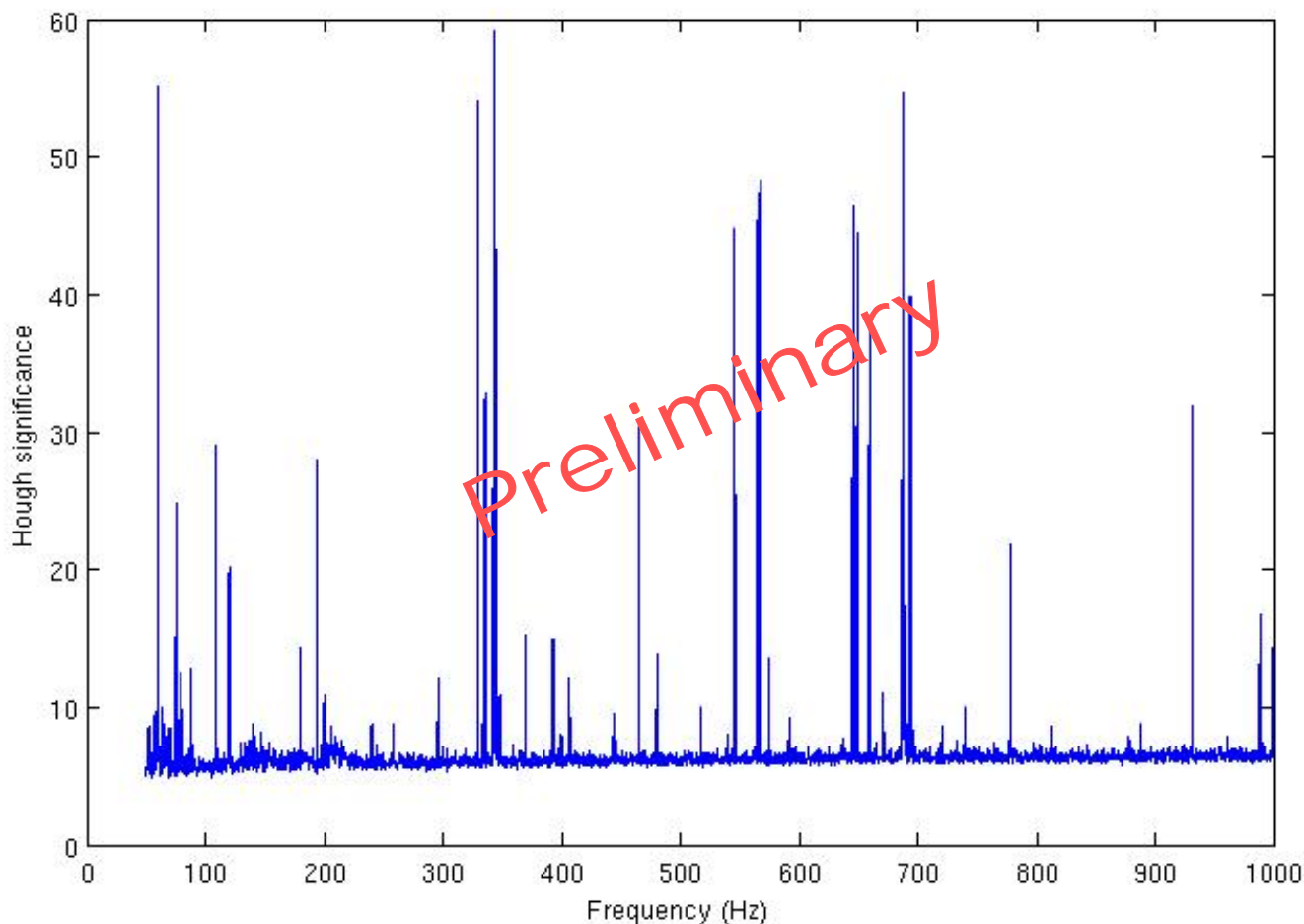


$$\langle n \rangle = 202.7$$

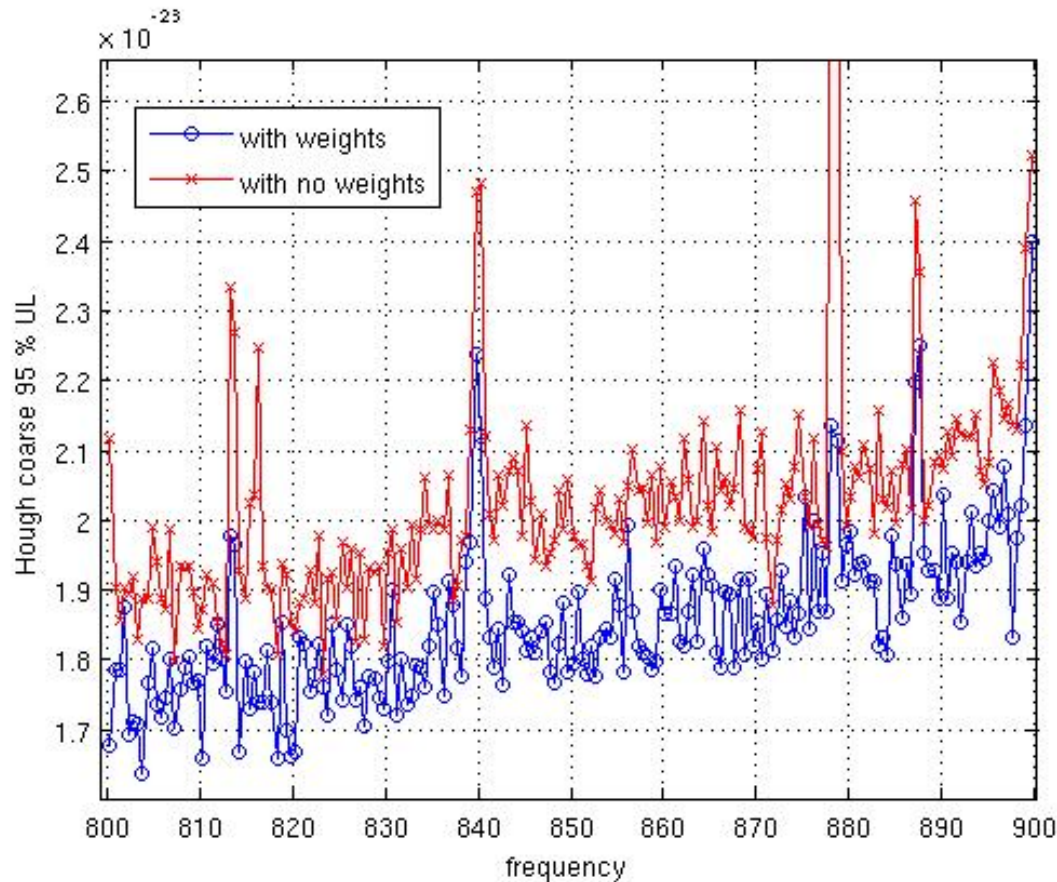
$$\sigma = 14.96$$



# Loudest events for every 0.25 Hz Multi interferometer case 50-1000 Hz

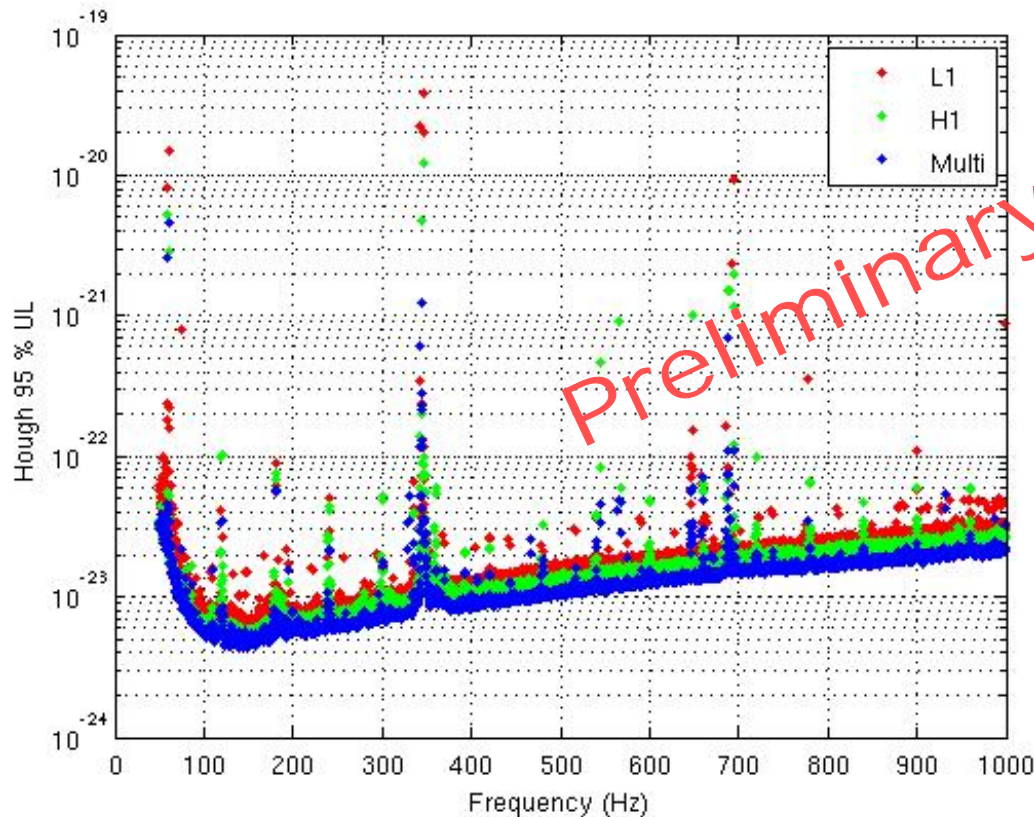


Significance defined as  $s = (n_{max} - \langle n \rangle) / \sigma$



Comparison of the All-sky 95% upper limits obtained by Monte-Carlo injections for the multi-IFO case.

The average improvement by using weights in this band is 9.25% for the multi-IFO case, but only ~6% for the single IFO



Best UL

for L1:  $5.9 \times 10^{-24}$

for H1:  $5.0 \times 10^{-24}$

for Multi:  $4.3 \times 10^{-24}$

It turns out that UL can be fitted by

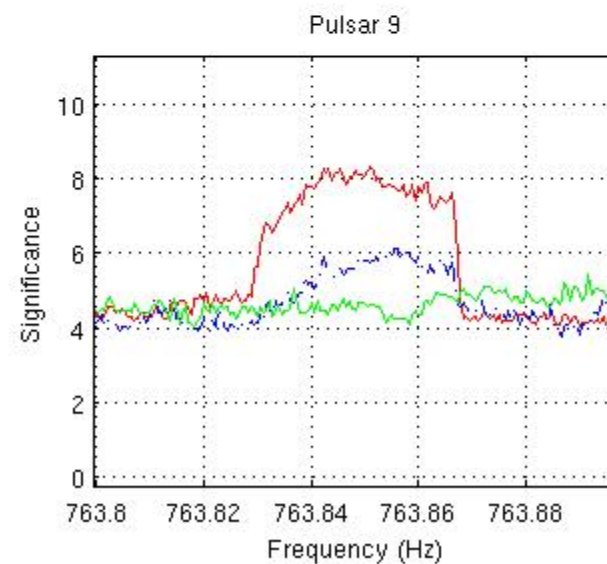
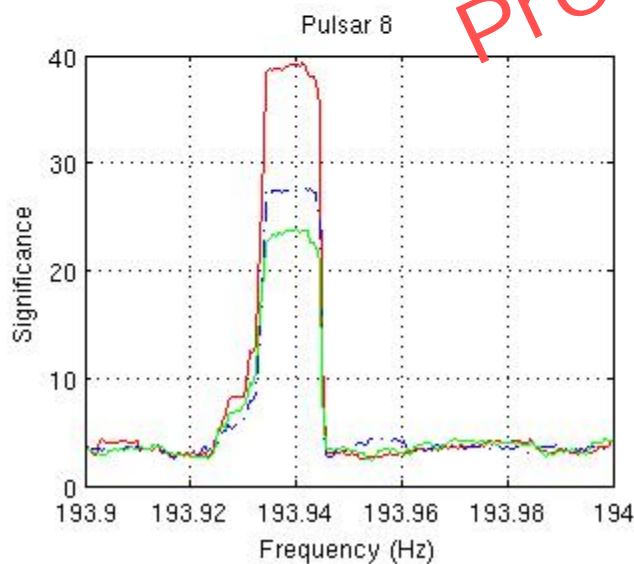
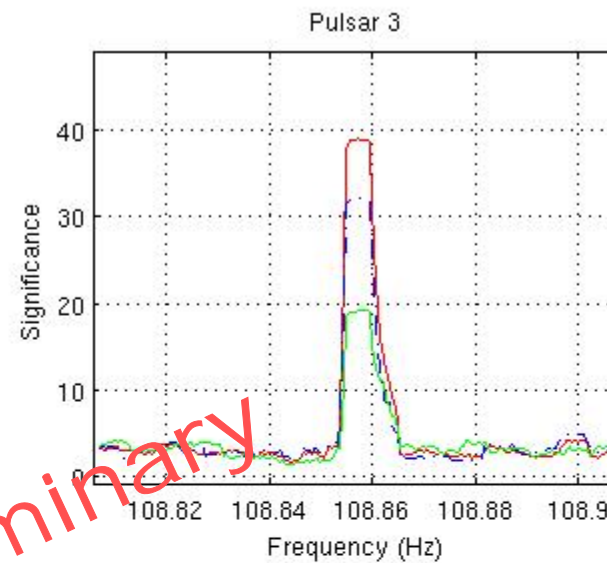
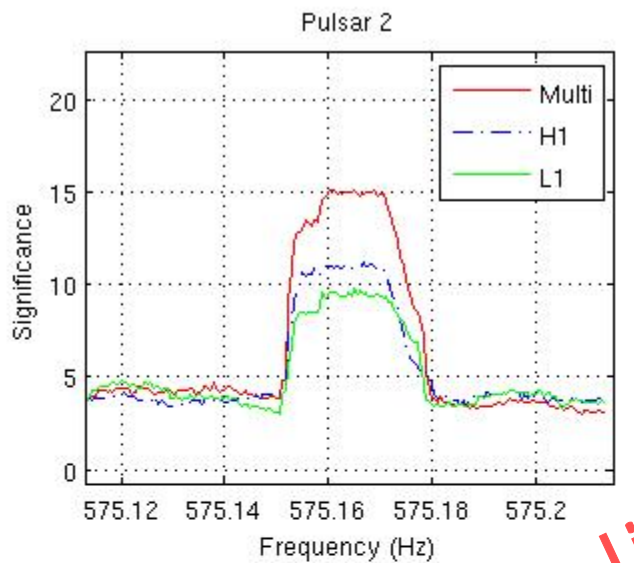
$$h_0^{95\%} = \frac{C}{\left( \sum_{i=1}^N \frac{1}{(S_n^{(i)})^2} \right)^{1/4}} \sqrt{\frac{S^{1/2}}{T_{coh}}},$$

with  $C=11.0 \pm 0.5$

$$S = \frac{\max(\text{significance})}{\sqrt{2}} + \text{erfc}^{-1}(0.1)$$



# Analysis of Hardware Injections



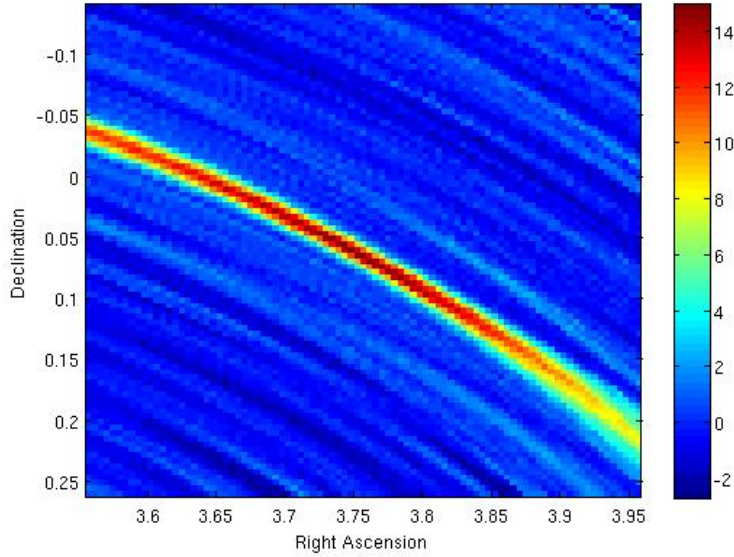
Preliminary



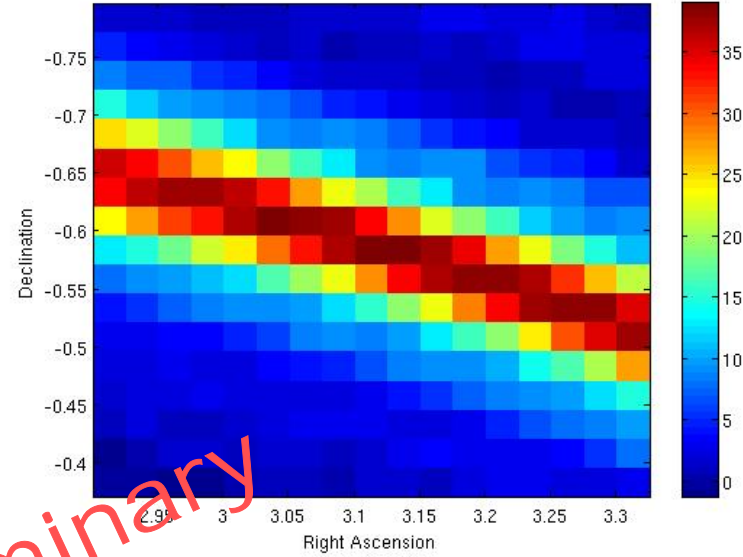
# Analysis of Hardware Injections



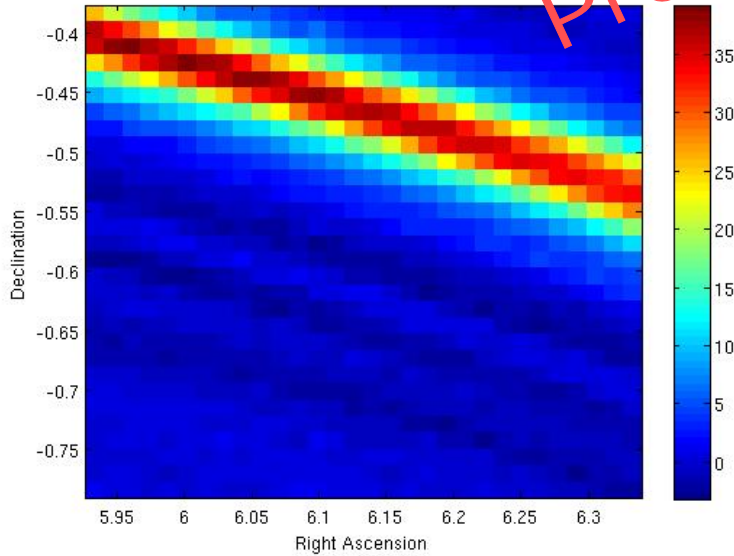
Pulsar 2, Freq =575.1633 Hz  $\dot{s}$  =0 Hz/s



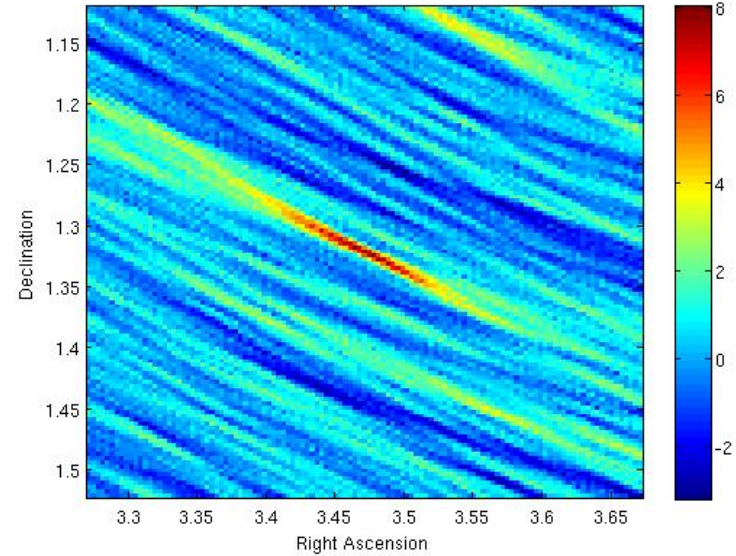
Pulsar 3, Freq =108.8572 Hz  $\dot{s}$  =0 Hz/s



Pulsar 8, Freq =193.9411 Hz  $\dot{s}$  =-8.3968e-09 Hz/s



Pulsar 9, Freq =763.8472 Hz  $\dot{s}$  =0 Hz/s



Preliminary





## Conclusions & Future work

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- Results are interesting, but much better results are on the way!
- Improvements for S5:
  - S5: ~2x better sensitivity, 12x or more data
  - Increase the time of the coherent step
  - Hough on F-statistic segments from multiple IFOs
  - Ongoing development of Hierarchical pipeline that combines coherent and semi-coherent searches