

# The Threat of Parametric Instabilities in Advanced Laser Interferometer Gravitational Wave Detectors

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**LIGO-G050513-00-Z**



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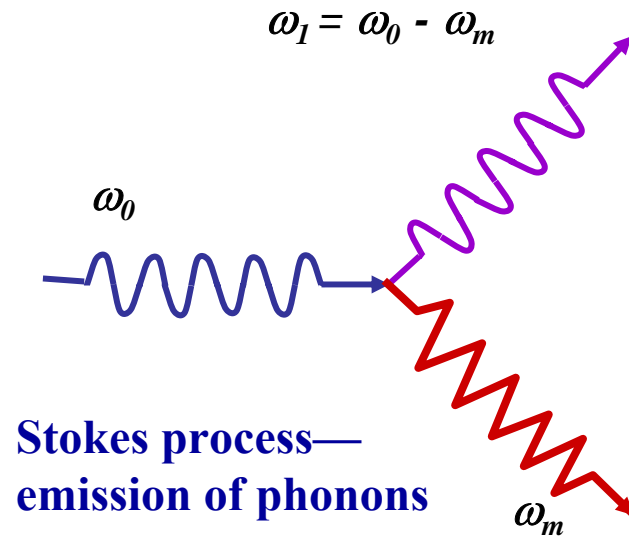
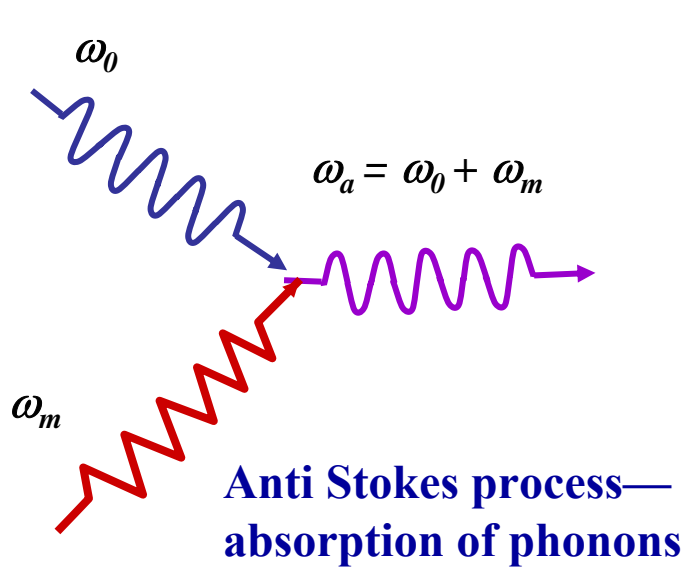
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- **Parametric instabilities**
- **Minefield for Advanced Detectors**
- **Suppression of instabilities**
  - **Thermal tuning**
  - **Q reduction**
  - **Feedback control**
- **Future work**

# When energy densities get high things go unstable...

- **Braginsky et al predicted parametric instabilities can happen in advanced detectors**
  - resonant scattering of photons with test mass phonons
  - acoustic gain like a laser gain medium

# Photon-phonon scattering

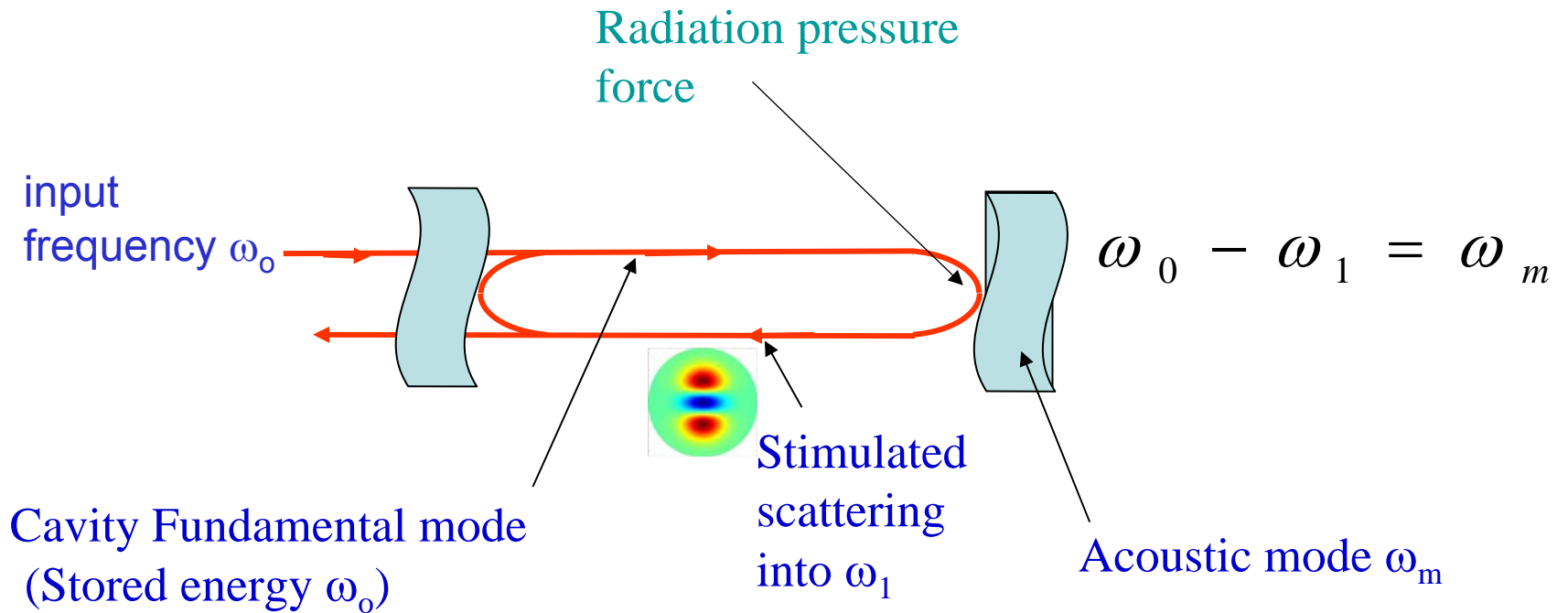


## Instabilities from photon-phonon scattering

- A test mass phonon can be **absorbed** by the photon, increasing the photon energy (**damping**);
- The photon can **emit** the phonon, decreasing the photon energy (**potential acoustic instability**).

# Schematic of Parametric Instability

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# Instability Condition

Parametric gain<sup>[1]</sup>

$$R \approx \frac{2PQ_m}{McL\omega_m^2} \left( \frac{Q_1\Lambda_1}{1 + \Delta\omega_1^2 / \delta_1^2} - \frac{Q_{1a}\Lambda_{1a}}{1 + \Delta\omega_{1a}^2 / \delta_{1a}^2} \right) > 1$$

Cavity Power      Mechanical Q      Stokes mode contribution      Anti-Stokes mode contribution

$$\Delta\omega_{1(a)} = |\omega_0 - \omega_{1(a)}| - \omega_m$$

$\Lambda$ —overlap factor

$$\delta_{1(a)} = \frac{\omega_{1(a)}}{2Q_{1(a)}}$$

Fundamental mode frequency

High order transverse mode frequency

Acoustic mode frequency

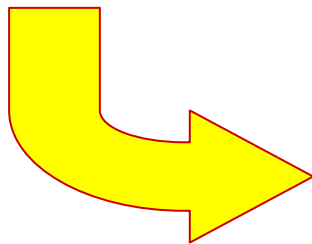
[1] V. B. Braginsky, S.E. Strigin, S.P. Vyatchanin, *Phys. Lett. A*, 305, 111, (2002)

# Instability conditions

- High circulating power  $P$
- High mechanical
- High optical mode  $Q$

+

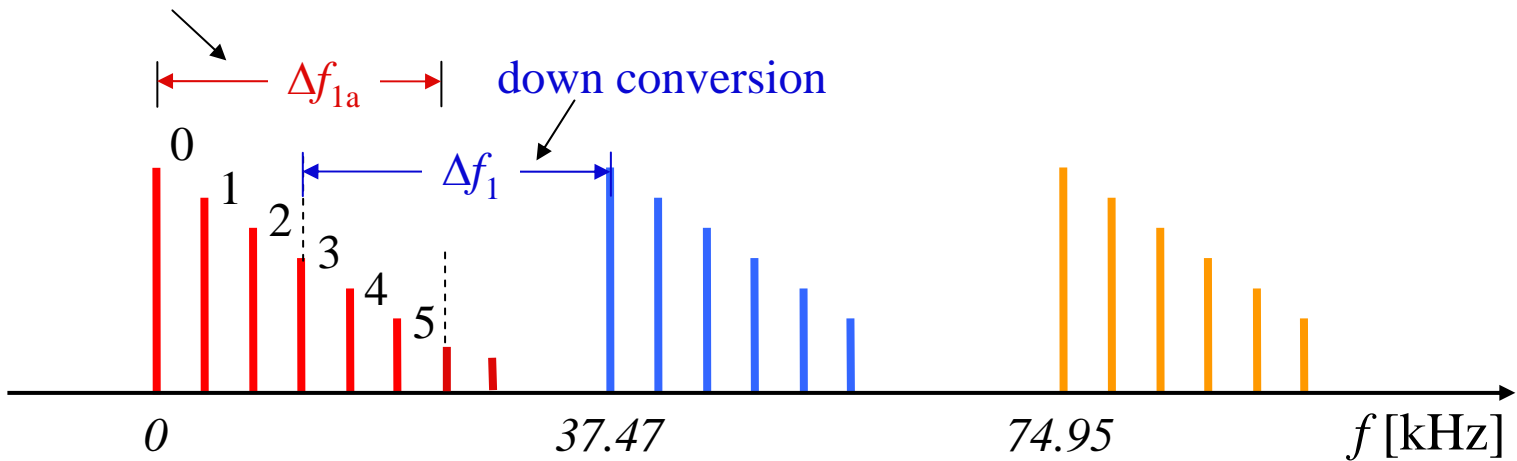
- Mode shapes overlap (High overlap factor  $\Lambda$ )
- Frequency coincidence— $\Delta\omega$  small



$R > 1,$   
Instability

# Mode Structure

upconversion



$$\left( \frac{Q_1 \Lambda_1}{1 + \Delta\omega_1^2 / \delta_1^2} - \frac{Q_{1a} \Lambda_{1a}}{1 + \Delta\omega_{1a}^2 / \delta_{1a}^2} \right)$$

$$\begin{aligned} \Lambda_1 &\neq \Lambda_{1a} \\ \delta_1 &\neq \delta_{1a} \end{aligned}$$

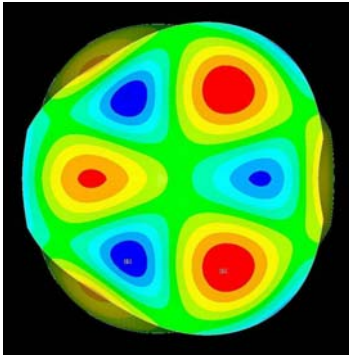


- Stokes & anti-Stokes modes contributions do not usually compensate

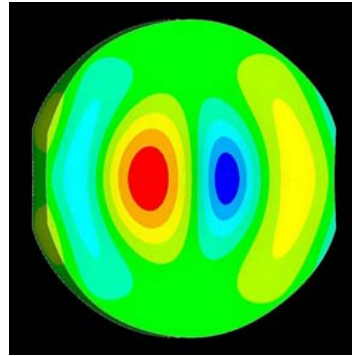


# Example of acoustic and optical modes for Al2O3 AdvLIGO

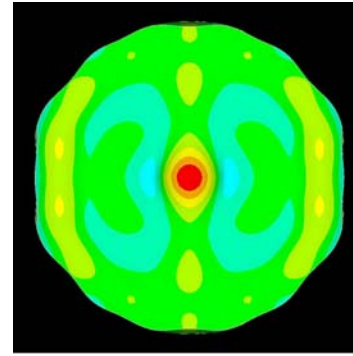
44.66 kHz



47.27 kHz

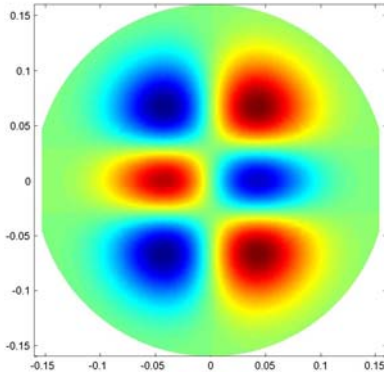


89.45kHz

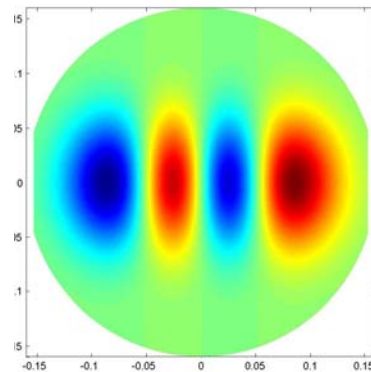


acoustic mode

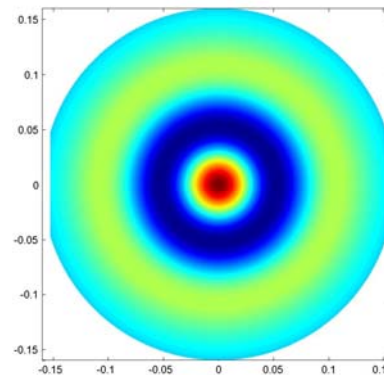
HGM12



HGM30



LGM20



optical mode

$\Lambda$

0.203

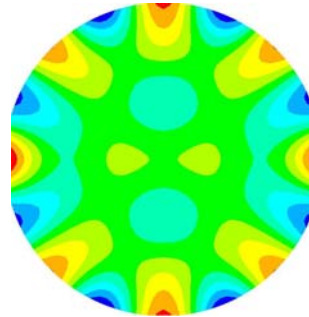
0.800

0.607

$\Lambda$  overlapping parameter

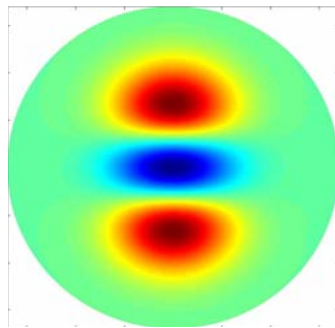
Summing over diagrams: multiple Stokes modes can drive a single acoustic mode.

## Example

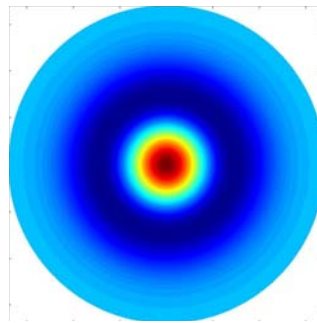


Mechanical mode shape  
( $f_m=28.34\text{kHz}$ )

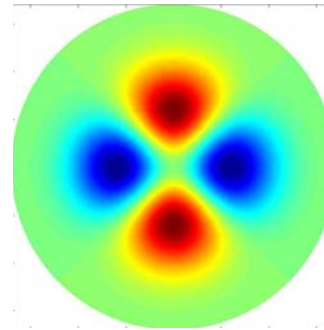
## Optical modes



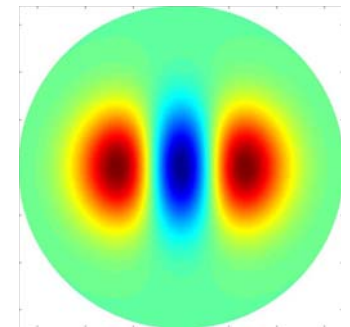
$\Lambda=0.007$   
 $R=1.17$



$\Lambda=0.019$   
 $R=3.63$



$\Lambda=0.064$   
 $R=11.81$



$\Lambda=0.076$   
 $R=13.35$

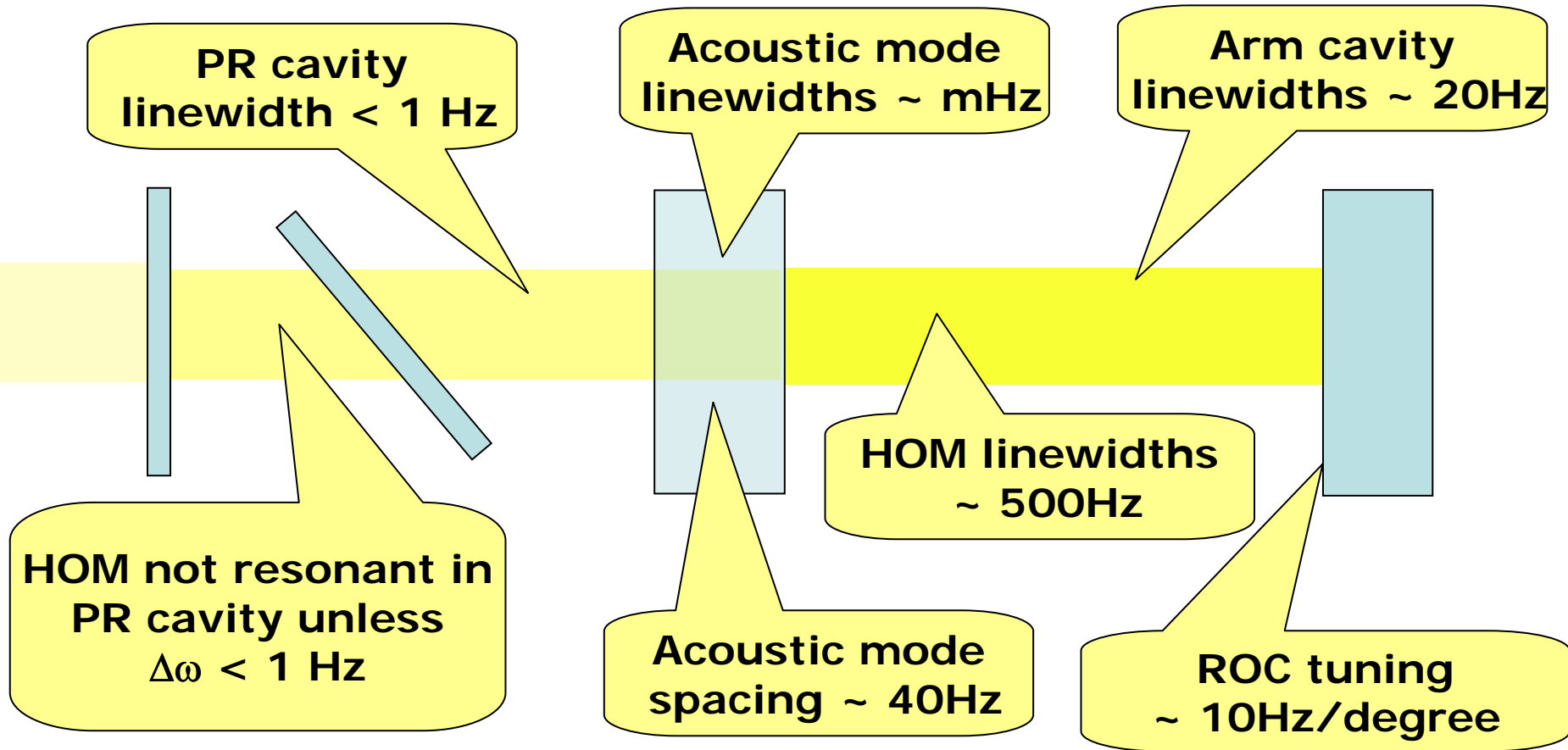
# Parametric gain— multiple modes contribution

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- Many Stokes/anti-Stokes modes can interact with single mechanical modes
- Parametric gain is the **sum** of all the possible processes

$$R = \frac{2PQ_m}{McL\omega_m^2} \left( \sum_{i=1}^{\infty} \frac{Q_{1i}\Lambda_{1i}}{1 + \Delta\omega_{1i}^2 / \delta_{1i}^2} - \sum_{j=1}^{\infty} \frac{Q_{1aj}\Lambda_{1aj}}{1 + \Delta\omega_{1aj}^2 / \delta_{1aj}^2} \right) > 1$$

# Influence of PR Cavity



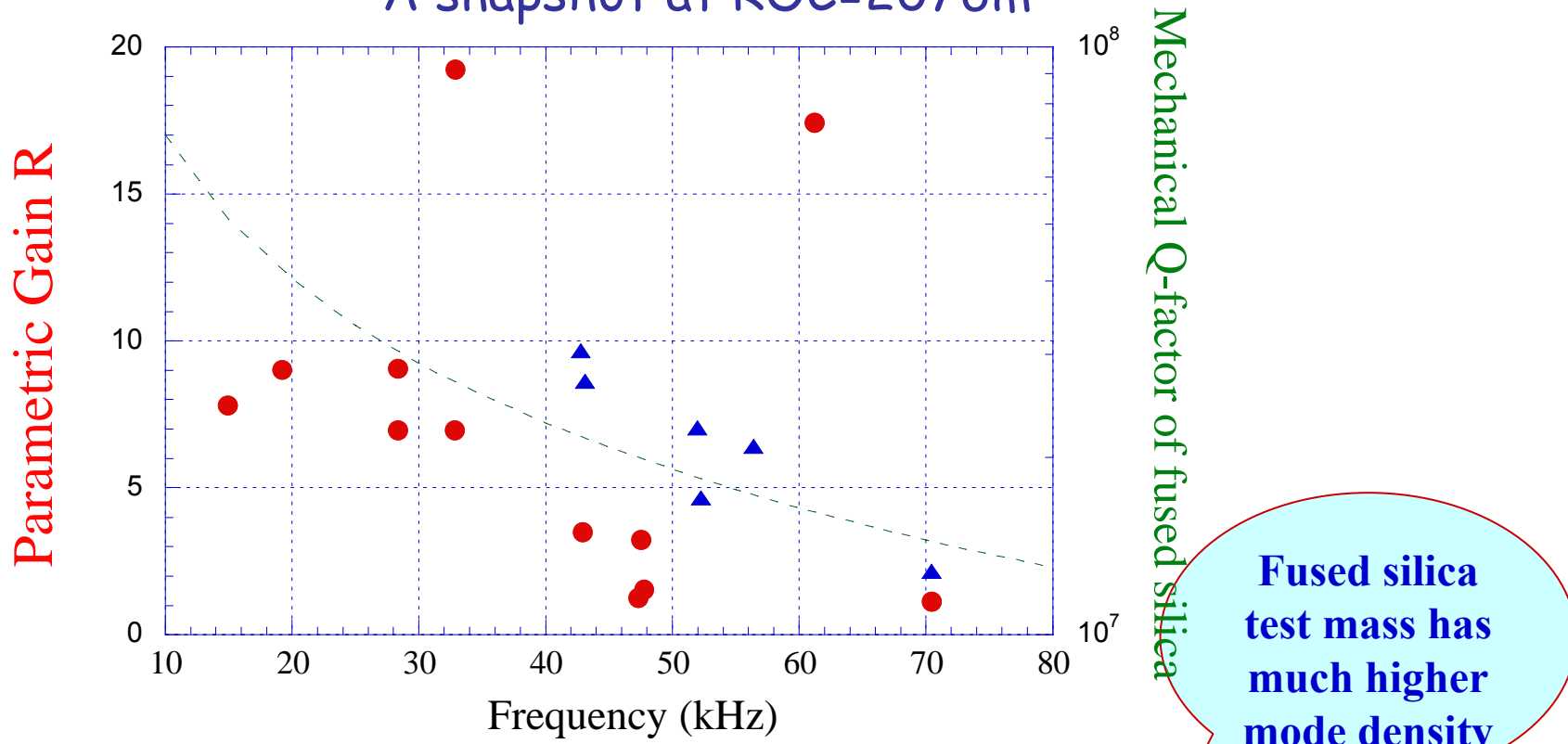
For  $\Delta\omega \gg 1$  Hz no recycling of HOM.

We calculate linewidths of HOMs from transmission + overlap loss of ideal modeshapes.

# Unstable modes for Adv/LIGO

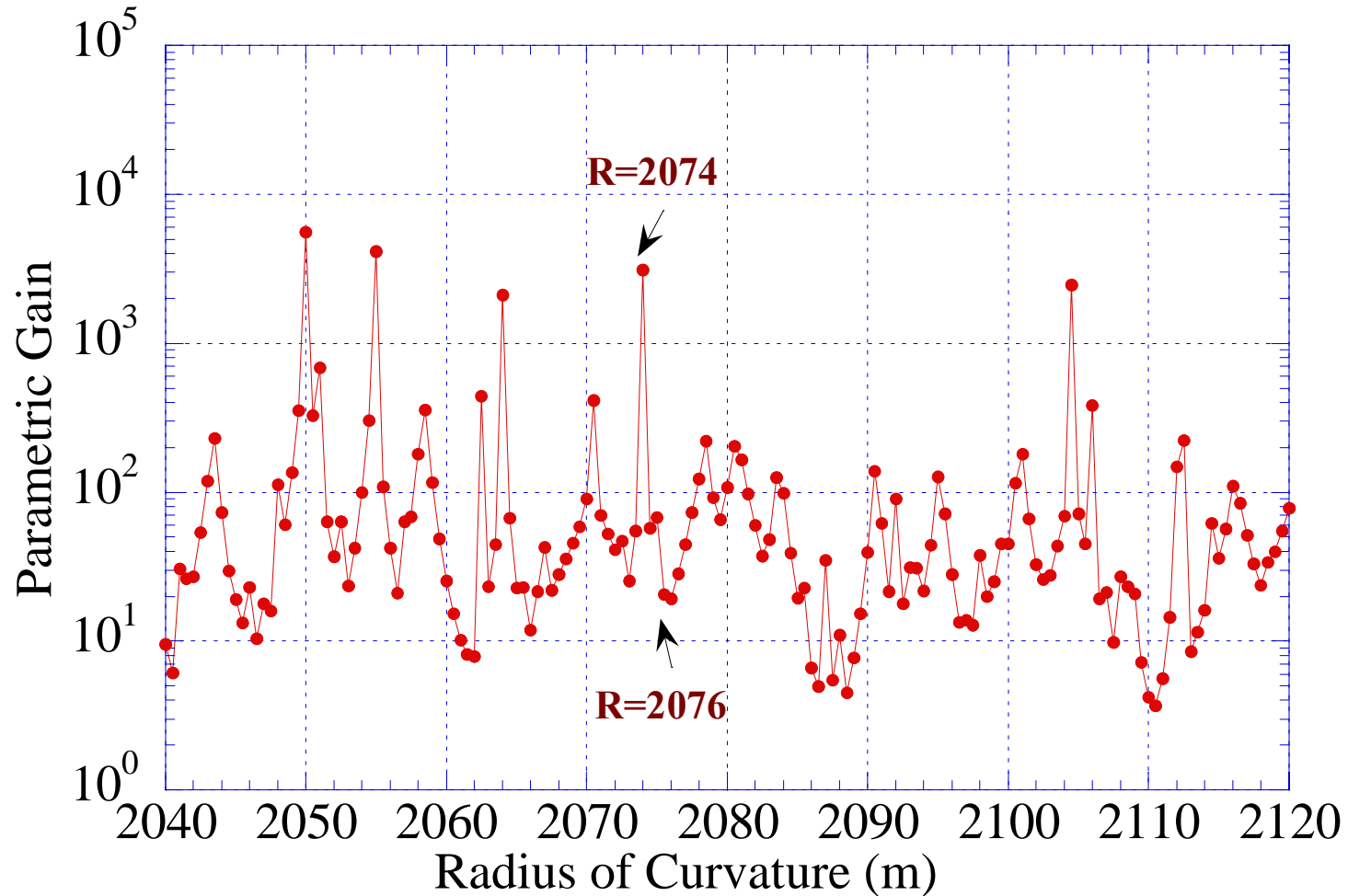
Sapphire & Fused silica nominal parameters

--A snapshot at ROC=2076m



- Sapphire —  $Q_m=10^8$ , 5 unstable modes (per test mass)
- Fused silica —  $Q_m(f)$ , 12 unstable modes (per test mass)

# Landmines! There is one at 2074!



# Instability Condition

Parametric gain<sup>[1]</sup>

$$R \approx \frac{2PQ_m}{McL\omega_m^2} \left( \frac{Q_1\Lambda_1}{1 + \Delta\omega_1^2 / \delta_1^2} - \frac{Q_{1a}\Lambda_{1a}}{1 + \Delta\omega_{1a}^2 / \delta_{1a}^2} \right) > 1$$

Stokes mode contribution      Anti-Stokes mode contribution

Cavity Power

Mechanical Q

$$\Delta\omega_{1(a)} = \left| \omega_0 - \omega_{1(a)} \right| - \omega_m$$

$\Lambda$ —overlap factor

$$\delta_{1(a)} = \frac{\omega_{1(a)}}{2Q_{1(a)}}$$

Fundamental mode frequency

High order transverse mode frequency

Acoustic mode frequency

[1] V. B. Braginsky, S.E. Strigin, S.P. Vyatchanin, *Phys. Lett. A*, 305, 111, (2002)

# Suppression of Parametric Instabilities

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- **Thermal tuning**
- **Mechanical Q-reduction**
- **Feedback control**



# Tuning Coefficients

HOM Frequency Depends on ROC

For 2km ROC, typical ROC tuning  $dR/dT \sim 1\text{m/K}$  for FS,  
10m/K for sapphire

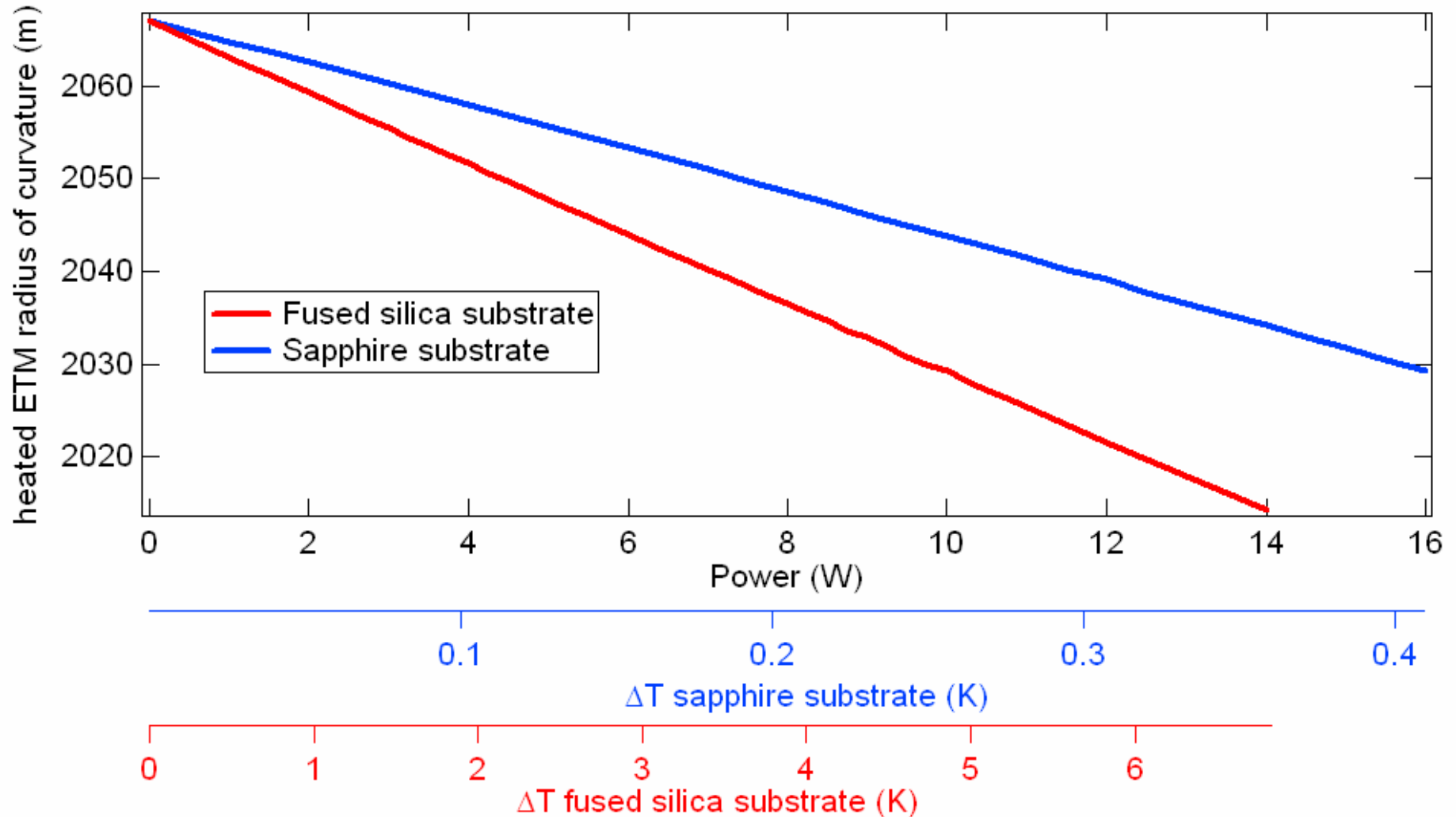
HOM frequency changes:  $df/dR \sim 10 \text{ Hz / m}$

Acoustic mode spacing:  $\sim 40\text{Hz}$  in fused silica

ROC uncertainty  $\sim 10\text{m}$  (?)

- **Change the curvature of mirror by heating**
- **Detune the resonant coupling**
- **How fast?**
- **How much R reduction?**

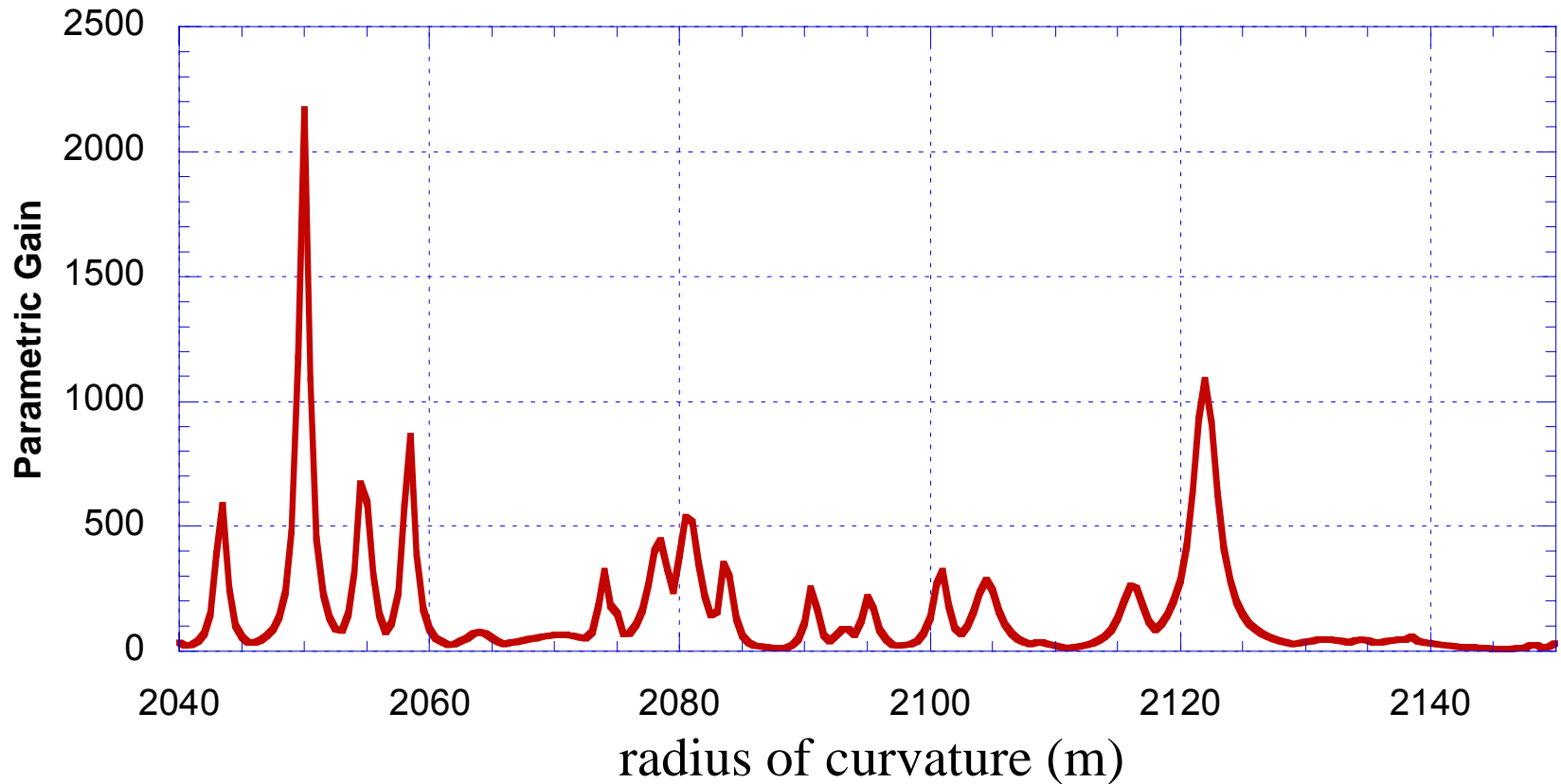
# ETM radius of curvature vs heating



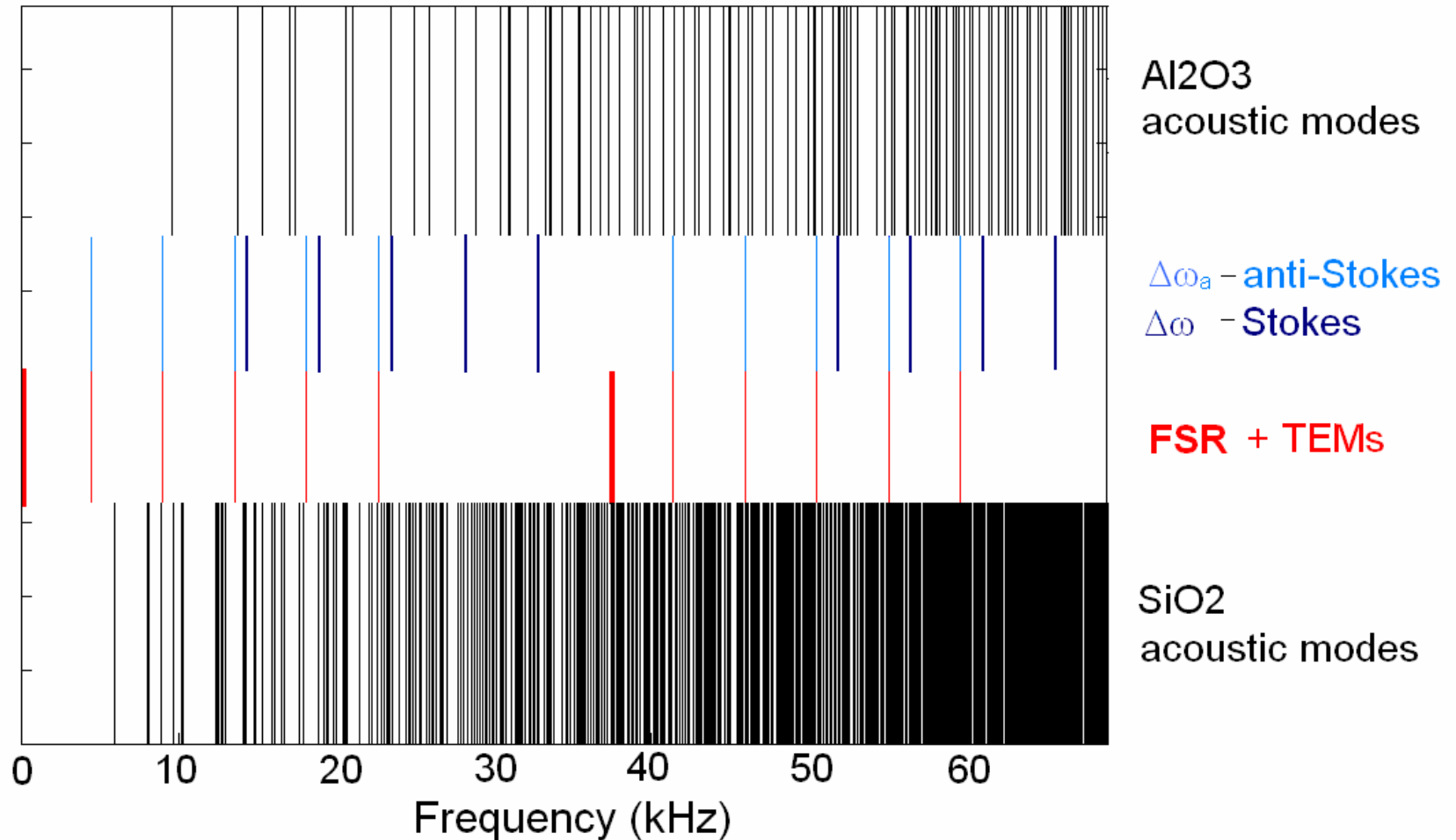
# Thermal tuning without PR Cavity

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**Fused silica**



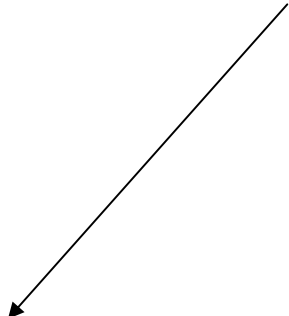
# Mode Structure for Advanced LIGO



If  $\Delta\omega - \omega_m < \text{optical linewidth}$  resonance condition may be obtained  
 $\Delta\omega = (n \cdot \text{FSR} - \text{TEM}_{mn})$  - frequency difference between the main and Stokes/anti-Stokes modes  
 $\omega_m$  - acoustic mode frequency,  $\delta$  - relaxation rate of TEM

# Instability Ring-Up Time

**Mechanical  
ring down  
time  
constant**



- For  $R > 1$ , ring-up time constant is  $\sim \tau_m / (R - 1)$

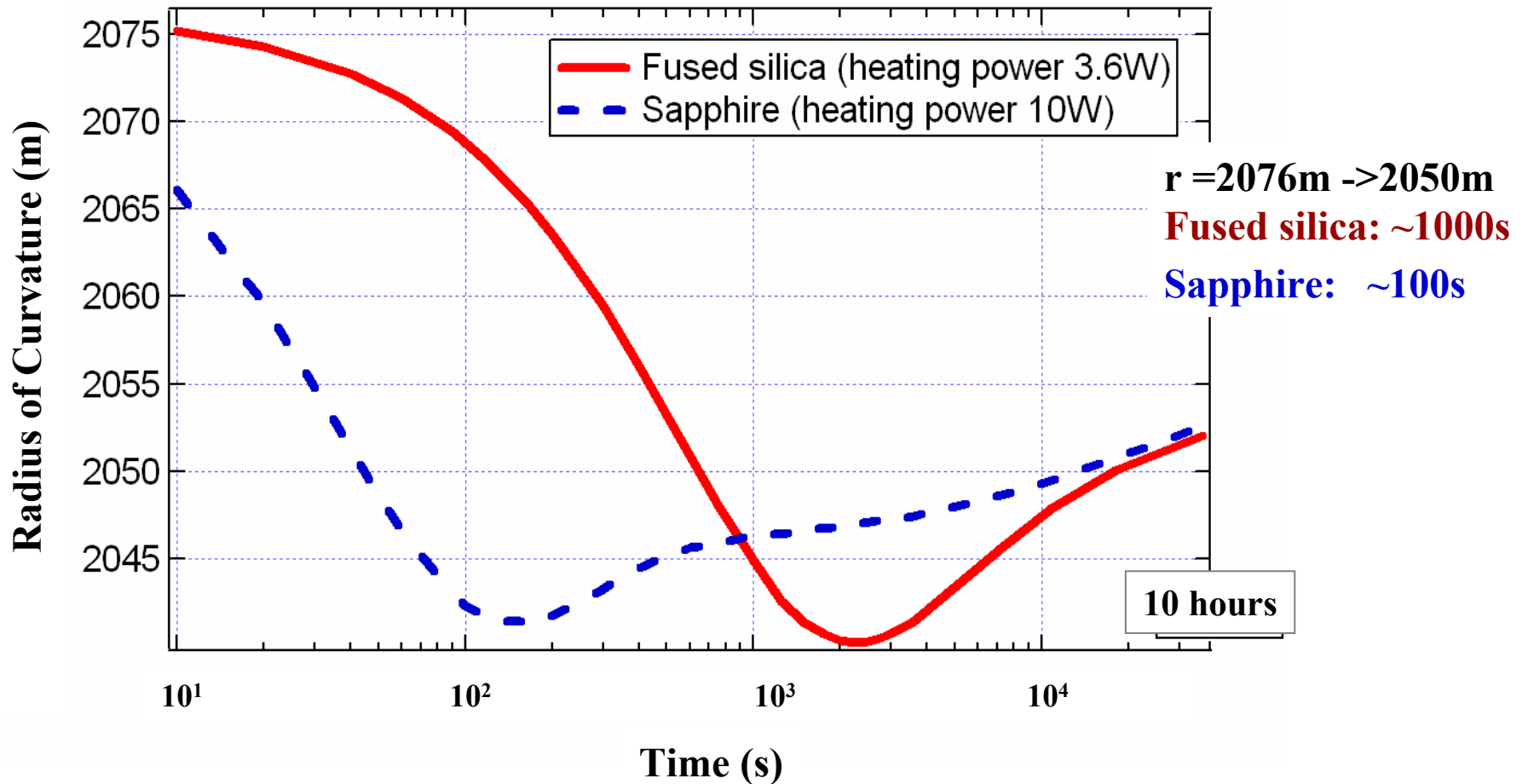
Time to ring from thermal amplitude to cavity position bandwidth ( $10^{-14}$  m to  $10^{-9}$  m) is

**$\sim 100-1000$  sec.**

- To prevent breaking of interferometer lock, cavities must be controlled within  $\sim 100$  s or less

# Thermal tuning time—sapphire is faster

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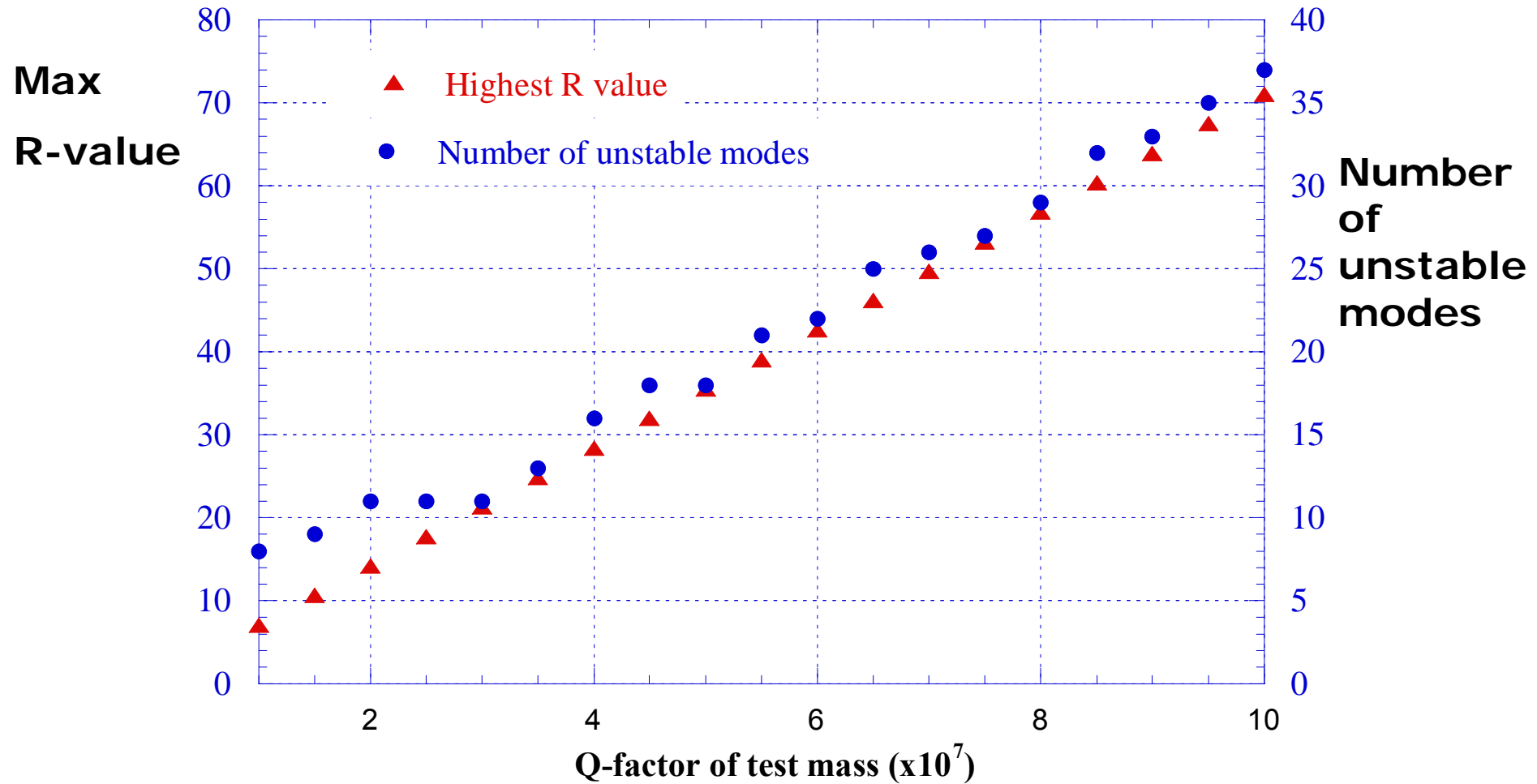
# Suppress parametric instabilities

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- Thermal tuning
- **Q-reduction**
- Feedback control

# Parametric instability and Q factor of test masses

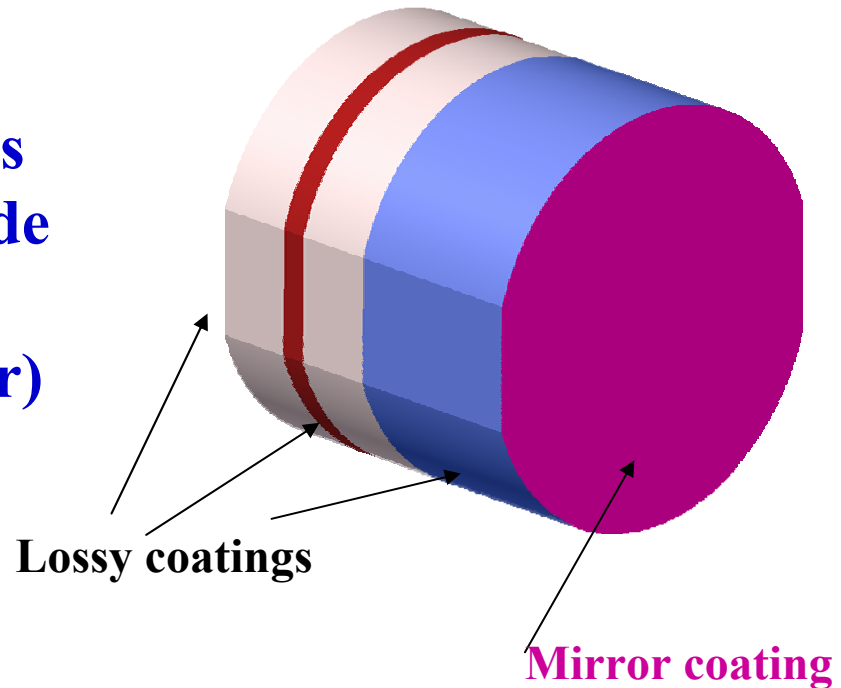
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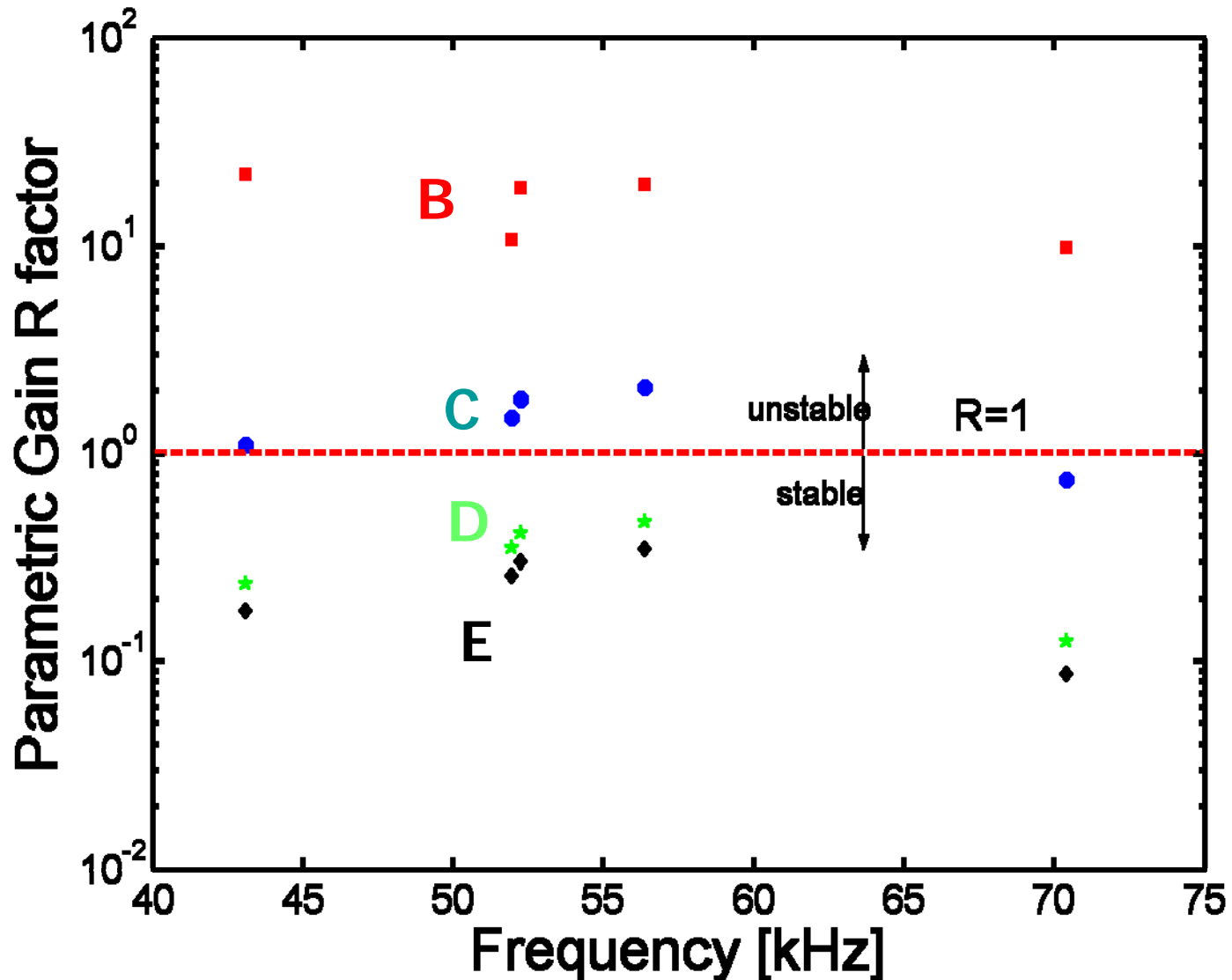


# Applying surface loss to reduce mode Q-factor

It is possible to apply lossy coatings ( $\phi \sim 10^{-4}$ ) on test mass to reduce the high order mode Q factors without degrading thermal noise (S. Gras poster)



# Parametric gain reduction

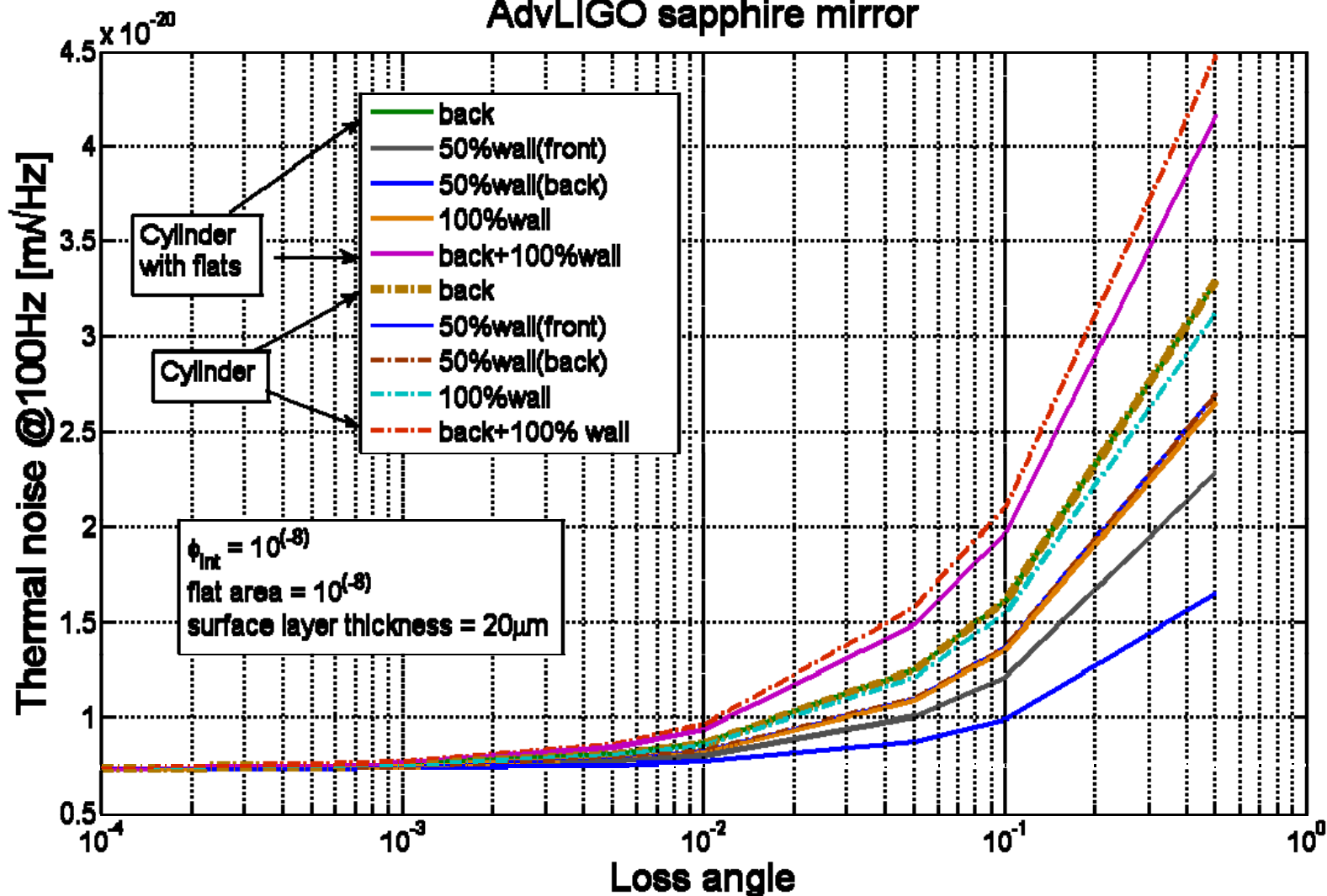


Unilateral stability for *nominal* AdvLIGO parameters

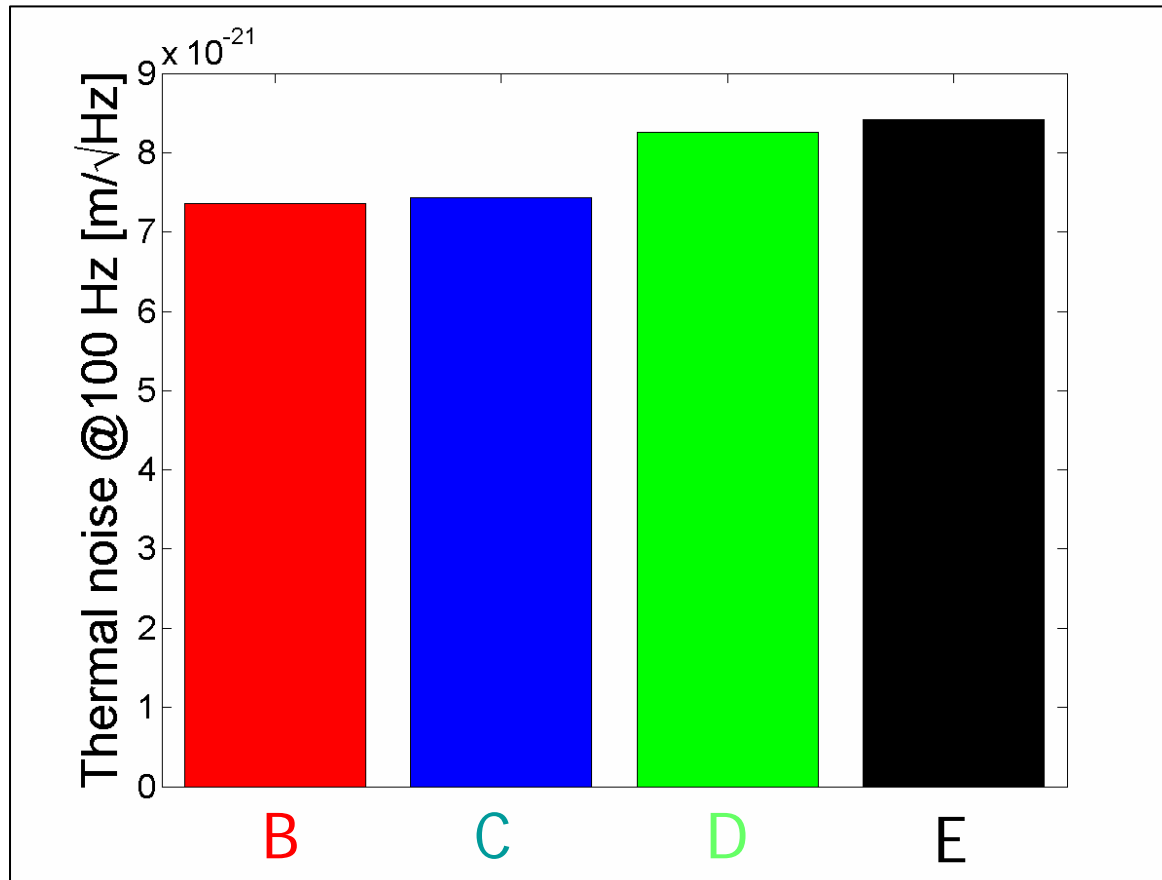
# Effect of localised losses on thermal noise

Side and Back

AdvLIGO sapphire mirror



# Noise increase 14% to achieve stability



# Suppress parametric instabilities

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- Thermal tuning
- Q-reduction
- **Feedback control**
  - **Problem: if test masses are similar but not identical instabilities will appear as quadruplets and individual test mass will not be identified unless well mode mapped before installation**

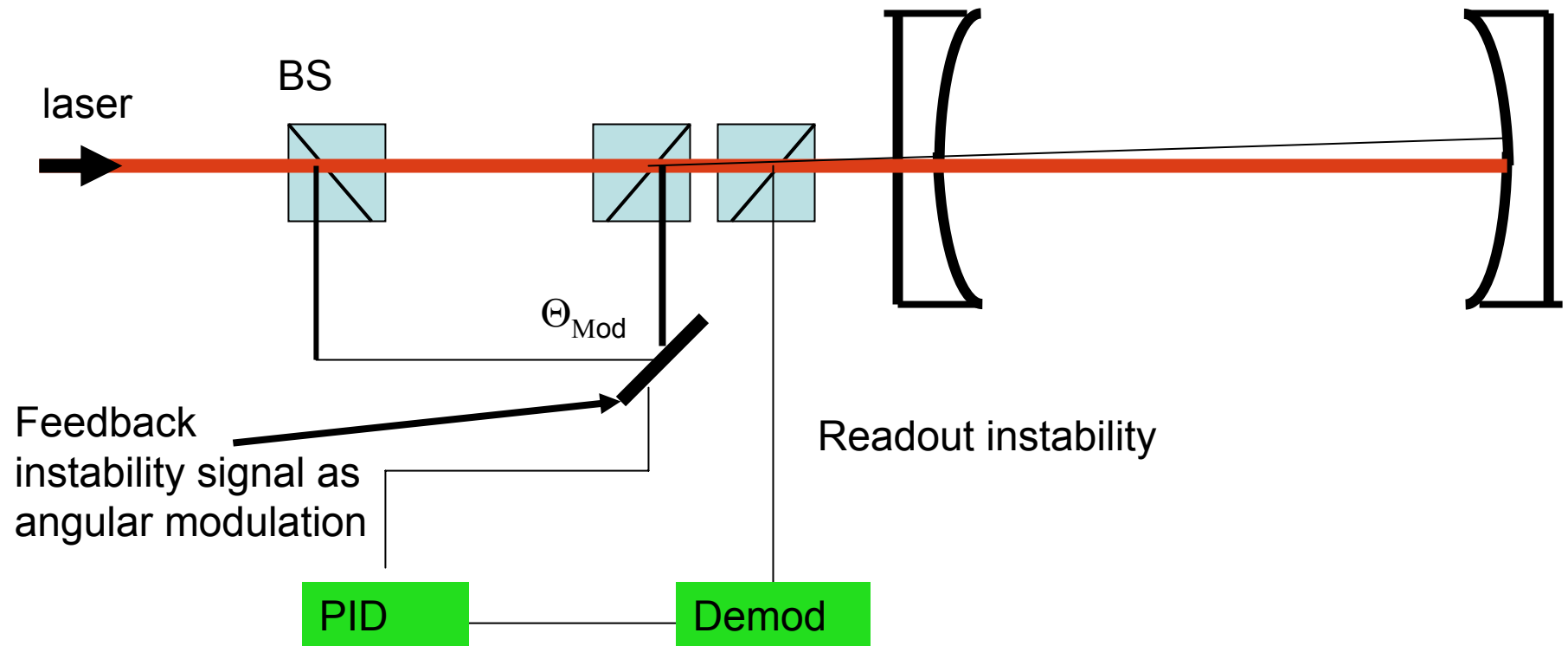
# Feedback control

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- Tranquiliser cavity (short external cavity )
  - Complex
- Direct force feedback to test masses
  - Capacitive local control or radiation pressure
  - Difficulties in distinguishing doublets/quadruplets
- Re-injection of phase shifted HOM
  - Needs external optics only
  - Multiple modes

# Direct Cold Damping by Feedback of HOM Signal

- HOM signal can by definition transmit in arm cavity



# Gingin HOPF Prediction

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- ACIGA Gingin high optical power facility 80m cavity should observe parametric instability effect with 10W power
- Expect to start experiment this year (Zhao's talk)



# Conclusions

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- Parametric instabilities are inevitable.
- FEM modeling accuracy/test masses uncertainties—  
precise prediction impossible
- Thermal tuning can minimise instabilities but can not  
completely eliminate instabilities.  
(*Zhao, et al, PRL, 94, 121102 (2005)*)
- Thermal tuning may be too slow in fused silica.
- Sapphire ETM gives fast thermal control and reduces  
total unstable modes (from ~64 to 43 (average))  
(*3 papers submitted to LSC review*)
- Instability may be actively controlled by various schemes
- Gingin HOPF is an ideal test bed for these schemes.