



# Violin Modes in Fused Silica Suspensions

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# Why Keep Studying Violin Modes?

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- Choice between ribbons & fibers will require better knowledge of noise performance.
- Suspension thermal noise is not expected to dominate optical noise in Advanced LIGO, but predicted low loss not yet demonstrated.
- Low frequency searches can reduce optical noise by reducing laser power and/or removing signal extraction mirror, thus exposing suspension thermal noise.
- Violin resonances can interfere with locking and control.

# Suspensions Apparatus

Automated fiber pulling lathe

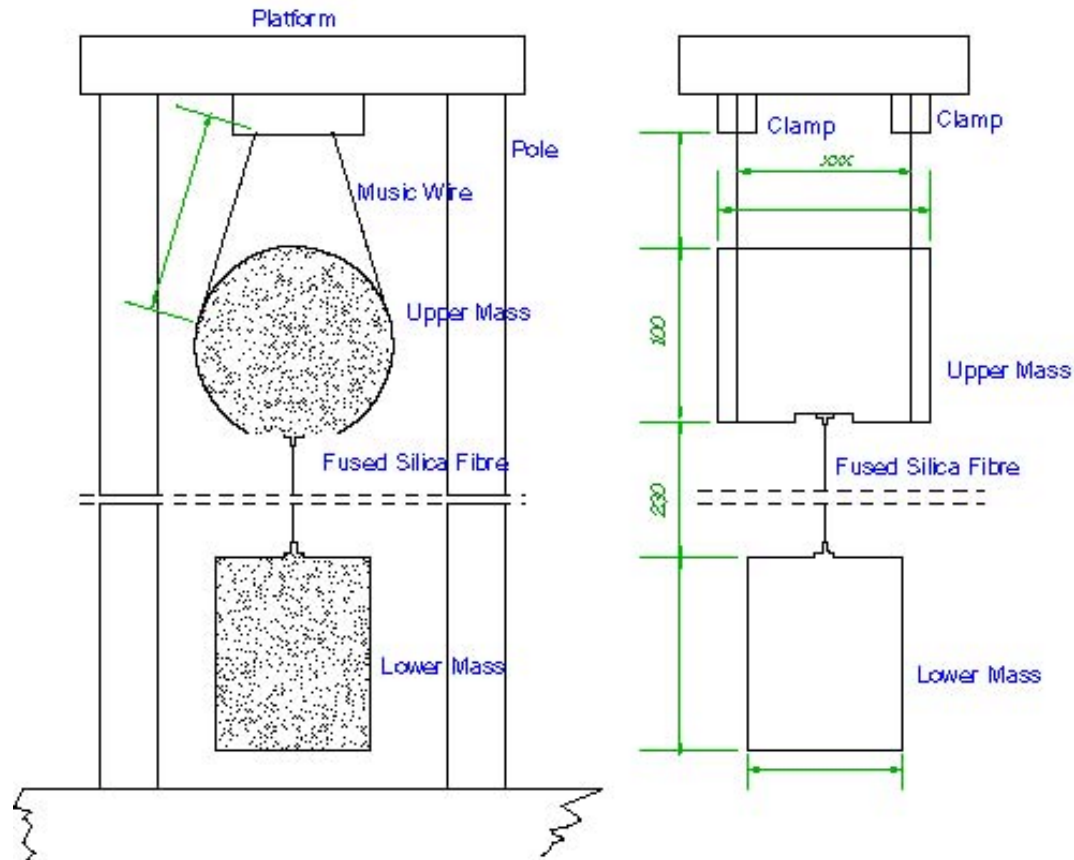


LIGO-G010168-00-D



Q-measurement rig

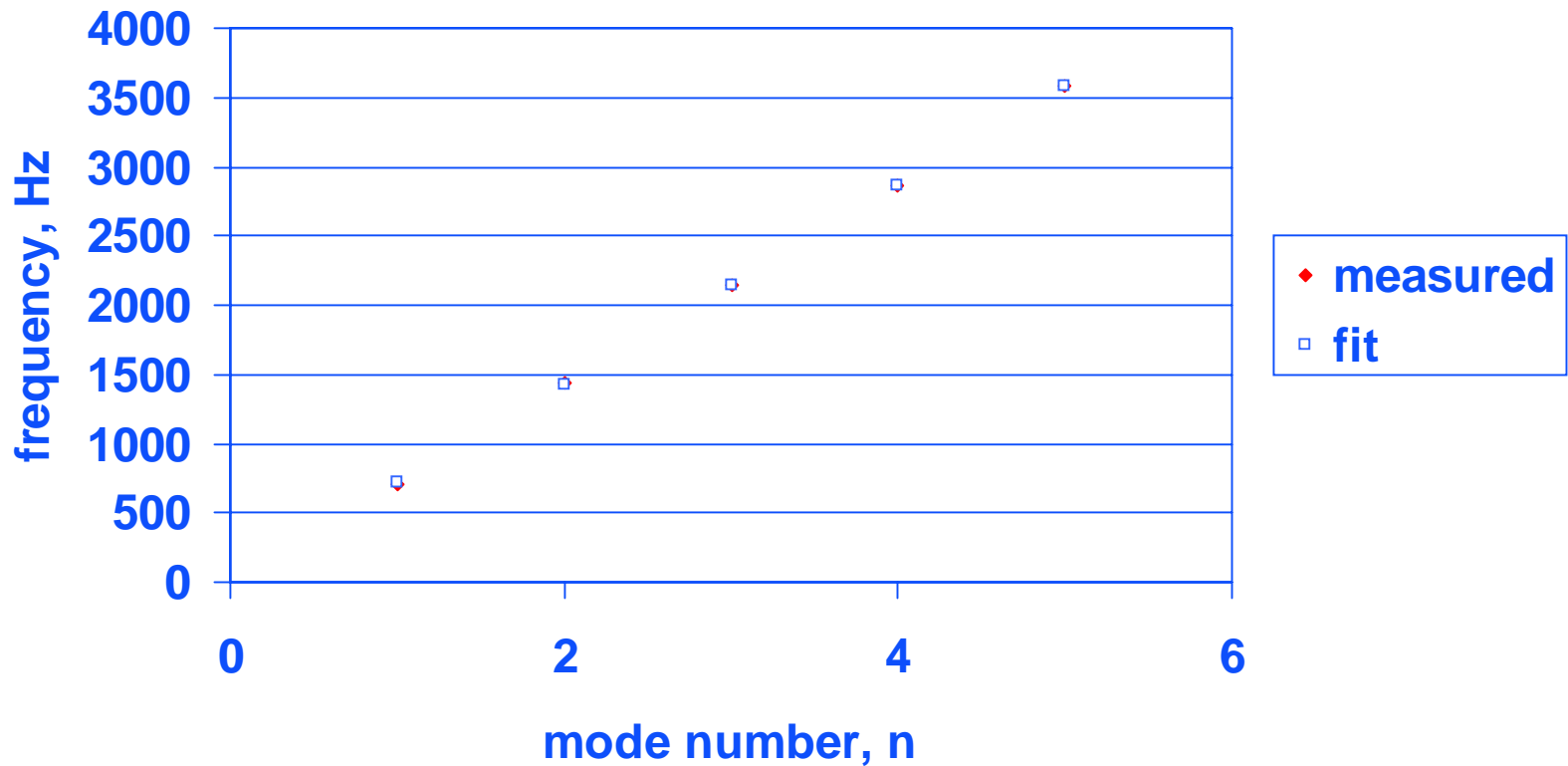
# Suspensions Apparatus



# Violin Mode Frequencies

$$f_n = \frac{n}{2L} \sqrt{\frac{P}{\rho_L}} \left[ 1 + \frac{2}{k_e L} + \frac{1}{2} \left( \frac{n\pi}{k_e L} \right)^2 \right]$$

# Measured Mode Frequencies

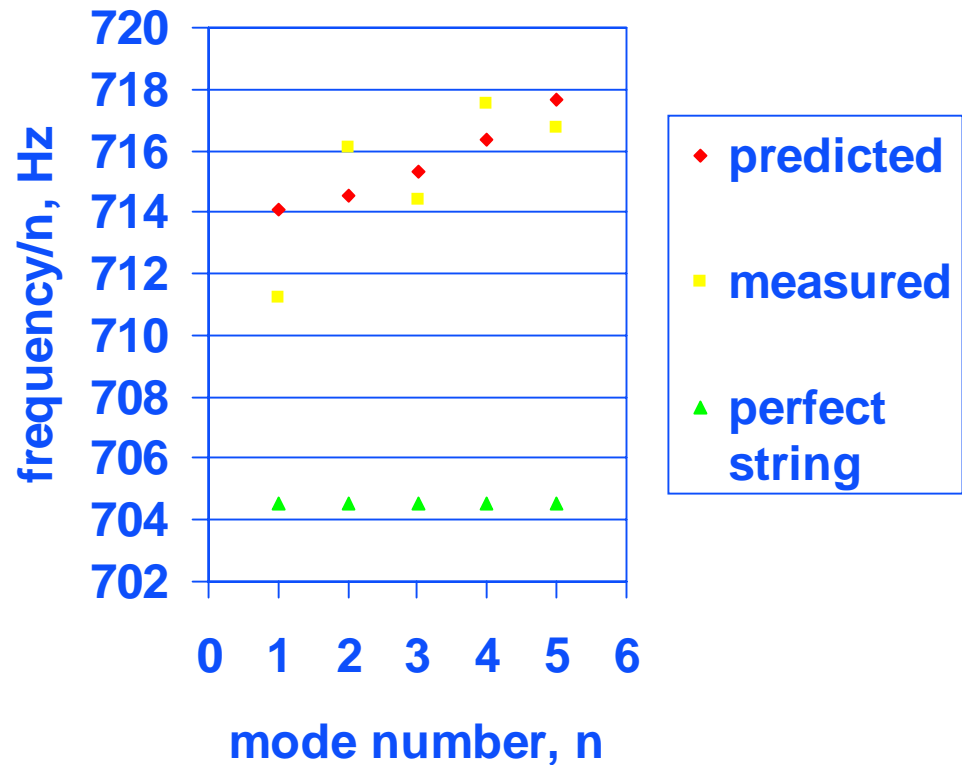


# Measured Mode Frequencies

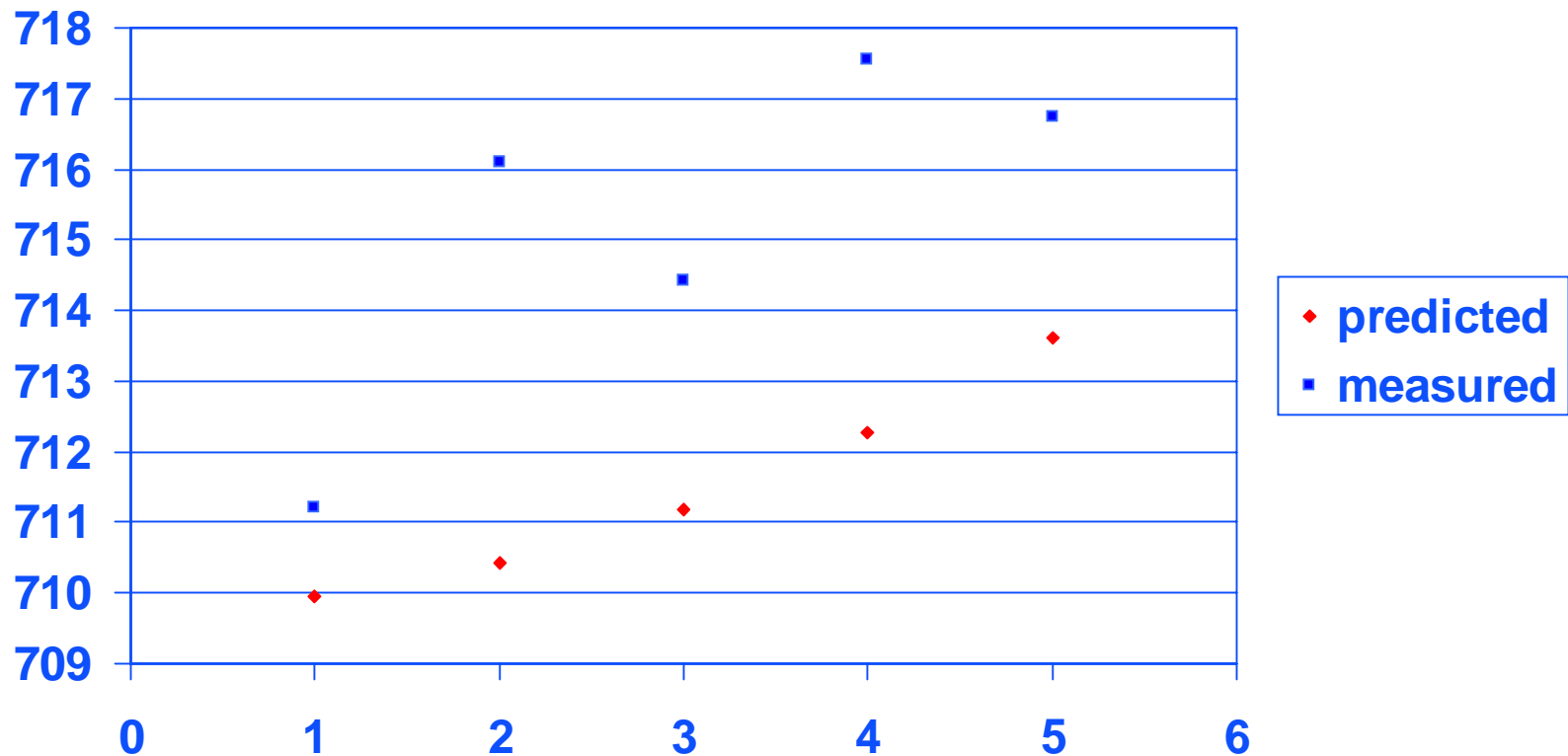
'Fit' of measured mode frequencies to model with radius as free parameter gives

$$r_{fiber} = 167.4 \mu m$$

with <.5% error for all modes

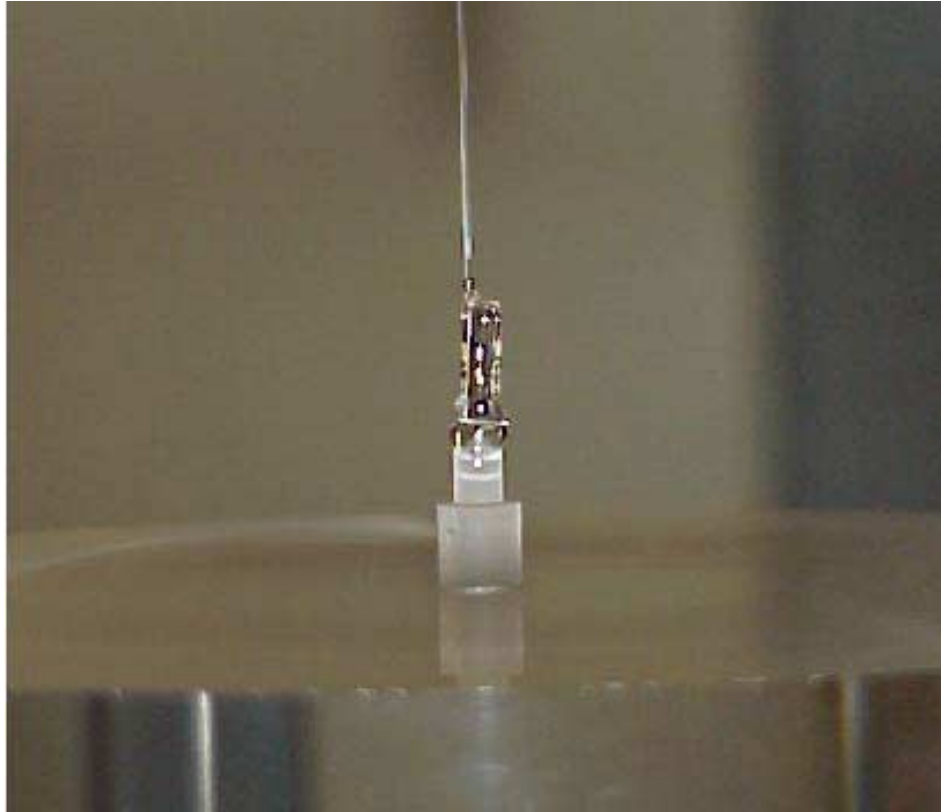


# How the frequencies change for $1\mu\text{m}$ change in fiber radius

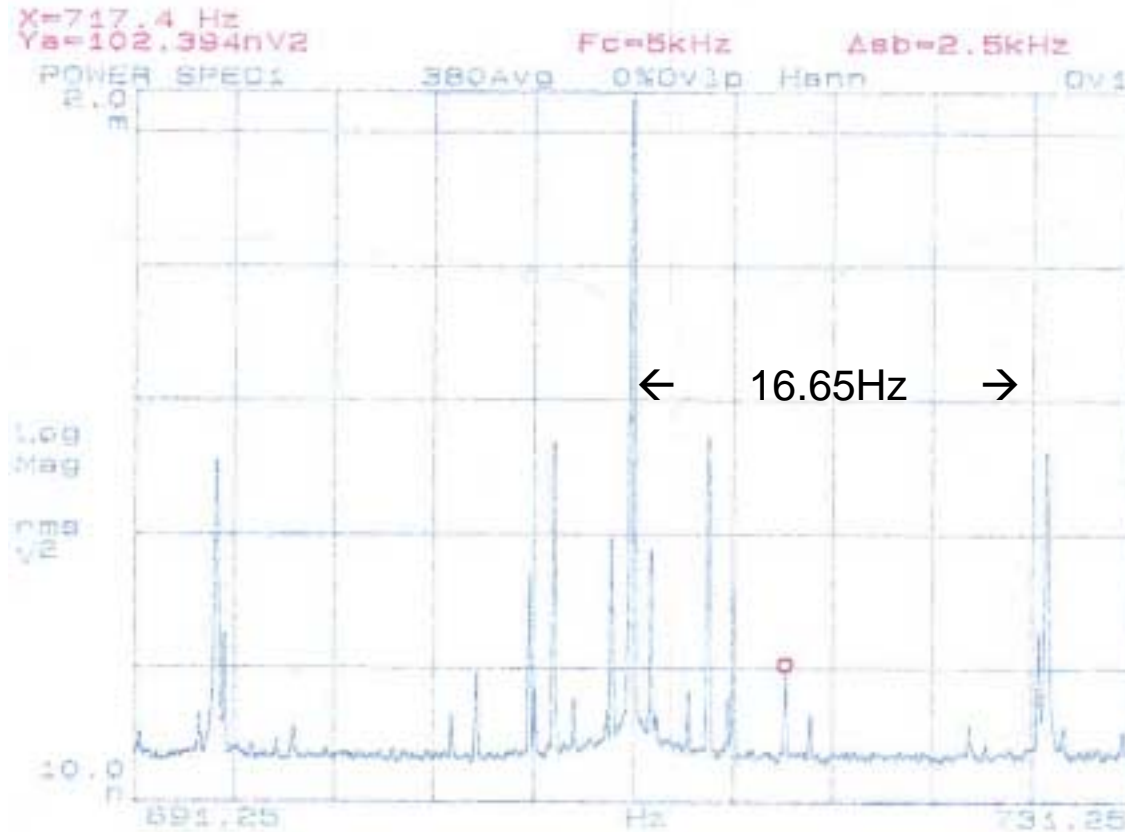




# Welded Fiber End



# Vertical Bounce Modulation



# Vertical Bounce and Violin Vibrato

Violin mode amplitude:

$$A(t) \approx A_0 \sin\left(\frac{n\pi}{L} \sqrt{\frac{P}{\rho_L}} t\right) = A_0 \sin(\omega_n t)$$

Strain amplitude:

$$P = P_0 + \delta P \sin(\Omega t)$$

$\Omega$ =bounce frequency

$$A(t) \approx A_0 \sin\left(\frac{n\pi}{L} \sqrt{\frac{P_0 + \delta P \sin(\Omega t)}{\rho_L}} t\right) \approx A_0 \sin\left(\omega_n \left[1 + \frac{\delta P}{2P_0} \sin(\Omega t)\right] t\right)$$

$$A(t) \approx A_0 \{J_0(m) \sin(\omega_n t) + J_1(m) \cos((\omega_n + \Omega)t) + J_1(m) \cos((\omega_n - \Omega)t)\}$$

where  $m = \delta P / 2P_0$  (in this case  $\delta P / P_0 \sim .1$ )

# Nonlinear Thermoelastic Damping

Loss function  $\phi$ :

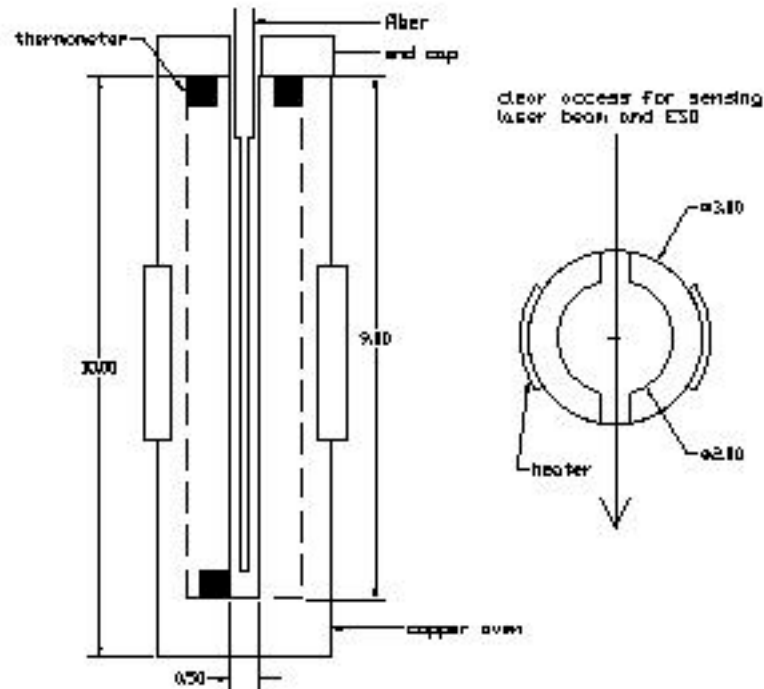
$$\phi_{\text{NTE}} = \frac{\text{TE}}{\rho C_V} \left( \alpha - \frac{dE}{dT} \frac{u}{E} \right)^2 \frac{\omega\tau}{1 + (\omega\tau)^2}$$

$$\tau_{\text{ribbon}} = \frac{t^2}{\pi^2} \frac{\rho C_V}{k} \quad \tau_{\text{fiber}} = \frac{.0737 \rho C_V r^2}{4k}$$

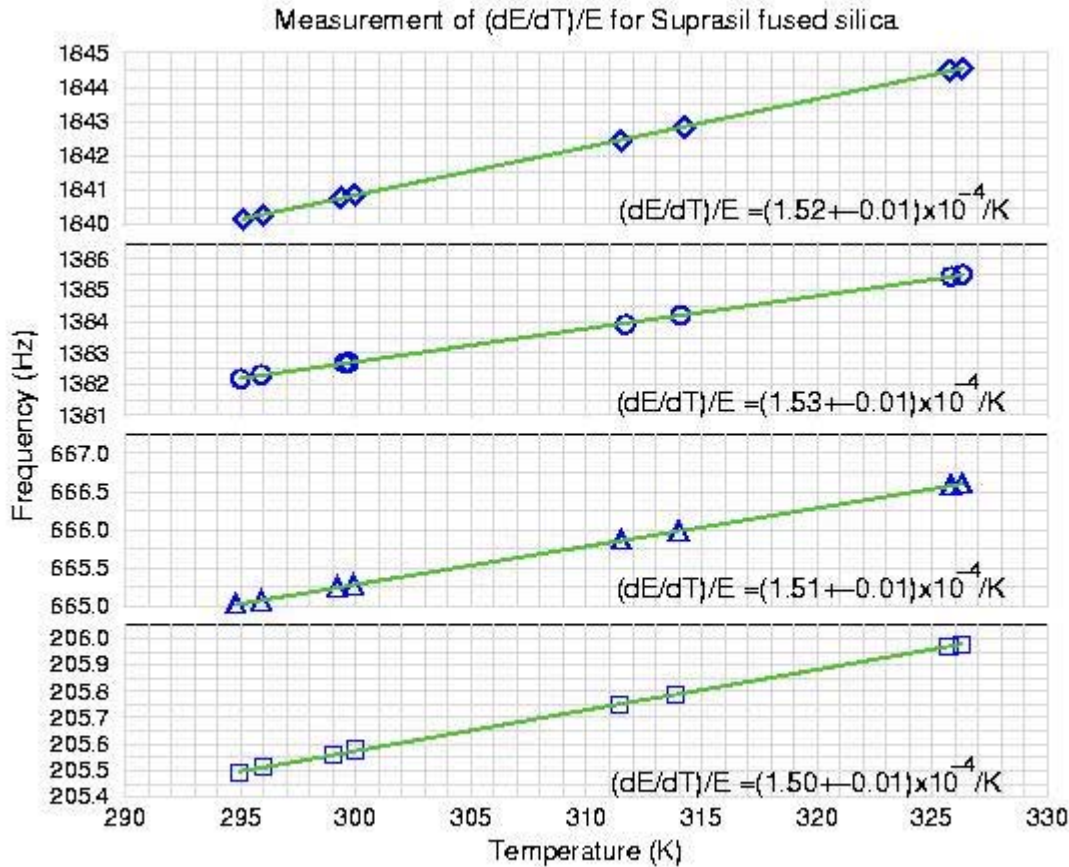
Measurement of  $\frac{dE/dT}{E}$ 

- Of all the parameters that determine  $\phi_{\text{NTE}}$ , the temperature dependence of the Young's modulus is most uncertain: published values vary by factor of 3.
  - » Measurements made from ~10kHz to GHz frequencies
  - » Measurements made using various mechanical & optical techniques
  - » Measurements made using many different types of fused silica
- Need exists for measurements on our fused silica (Suprasil 2), in our frequency range, and in our mode of oscillation (bending).

# $\frac{dE/dT}{E}$ Measurement Apparatus



# $\frac{dE/dT}{E}$ Data



Data consistent with  $(dE/dT)/E = 1.52e-4/K \pm 5\%$ .  
Main source of error is calibration of oven.

Note: B. Lunin measures  $2.2e-4/K$  for similar glass (Suprasil 300).

# Total Intrinsic Loss in Fibers

$$\phi_{total} = \phi_{bulk} + \phi_{surface} + \phi_{NTE}$$

$$\phi_{bulk} = 3 \times 10^{-8}$$

$$\phi_{surface} = \frac{3 \times 10^{-11}}{r}$$



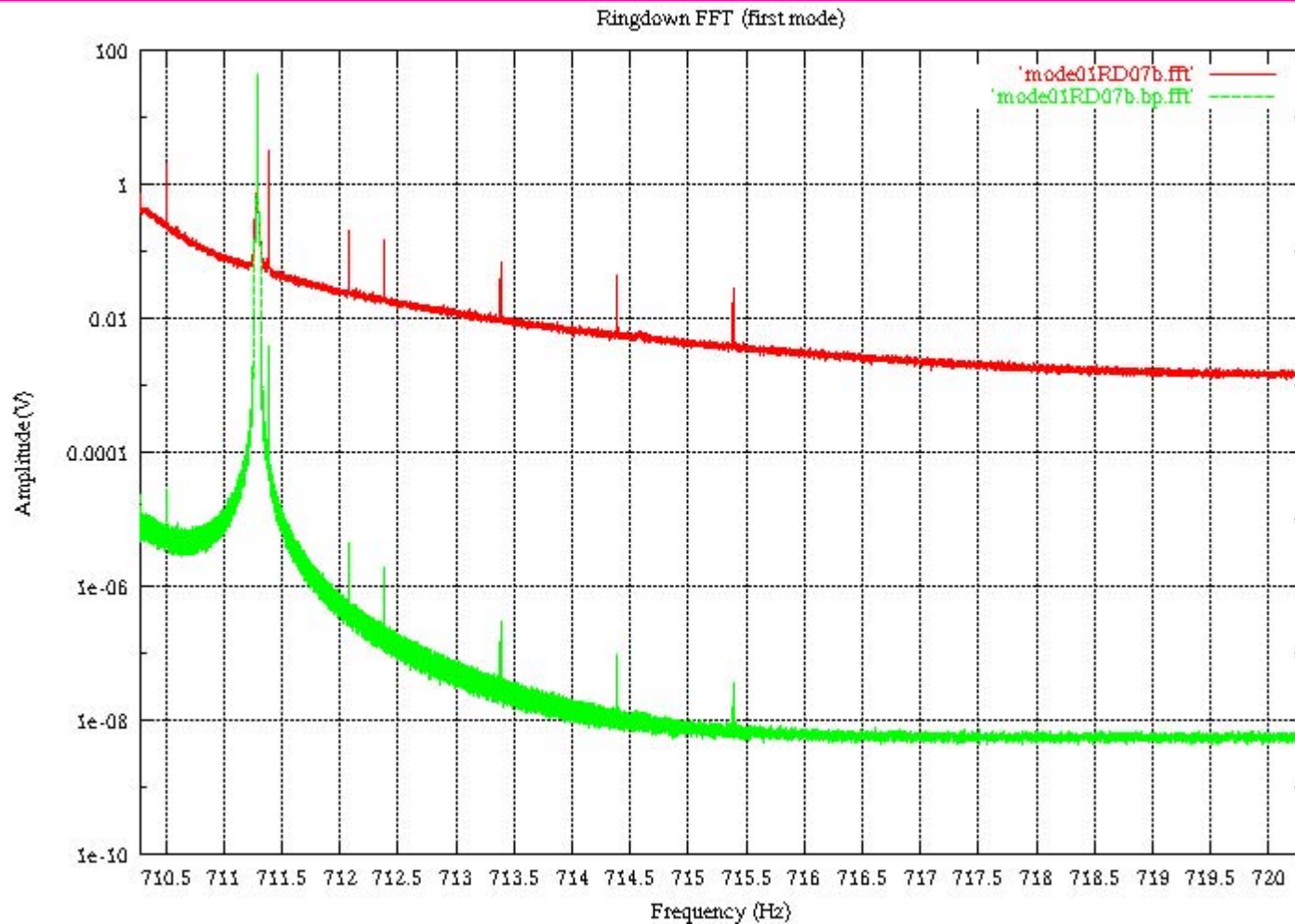
# Q of a Violin Mode

$$\frac{1}{Q_{violin}} = \phi_{total} \frac{2}{k_e L} \left[ 1 + \frac{(n\pi)^2}{2k_e L} \right]$$

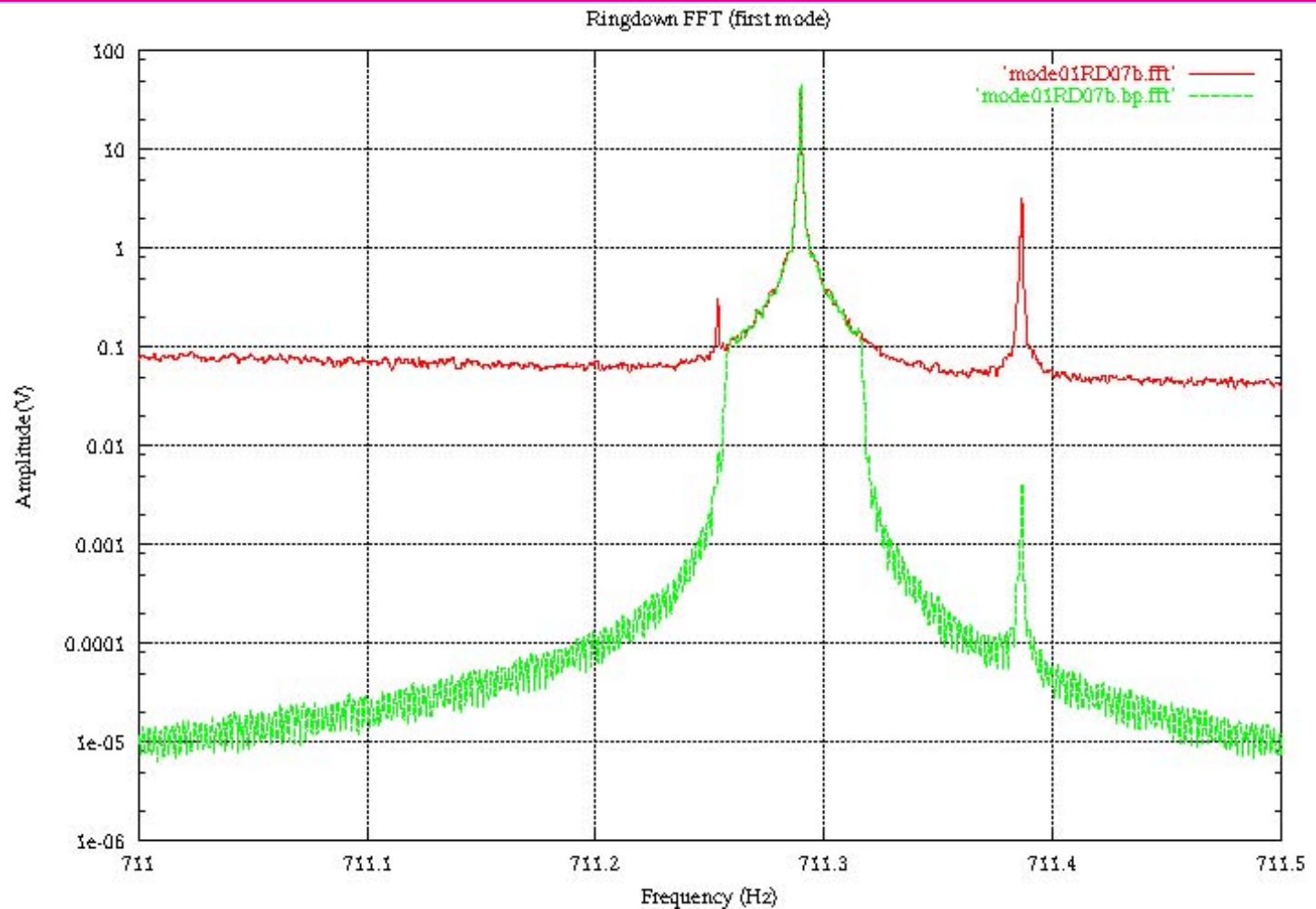
$n$  = mode number

$$k_e = \sqrt{\frac{P + \sqrt{P^2 + 4EI\rho_L\omega^2}}{2EI}} = \text{elastic wavenumber}$$

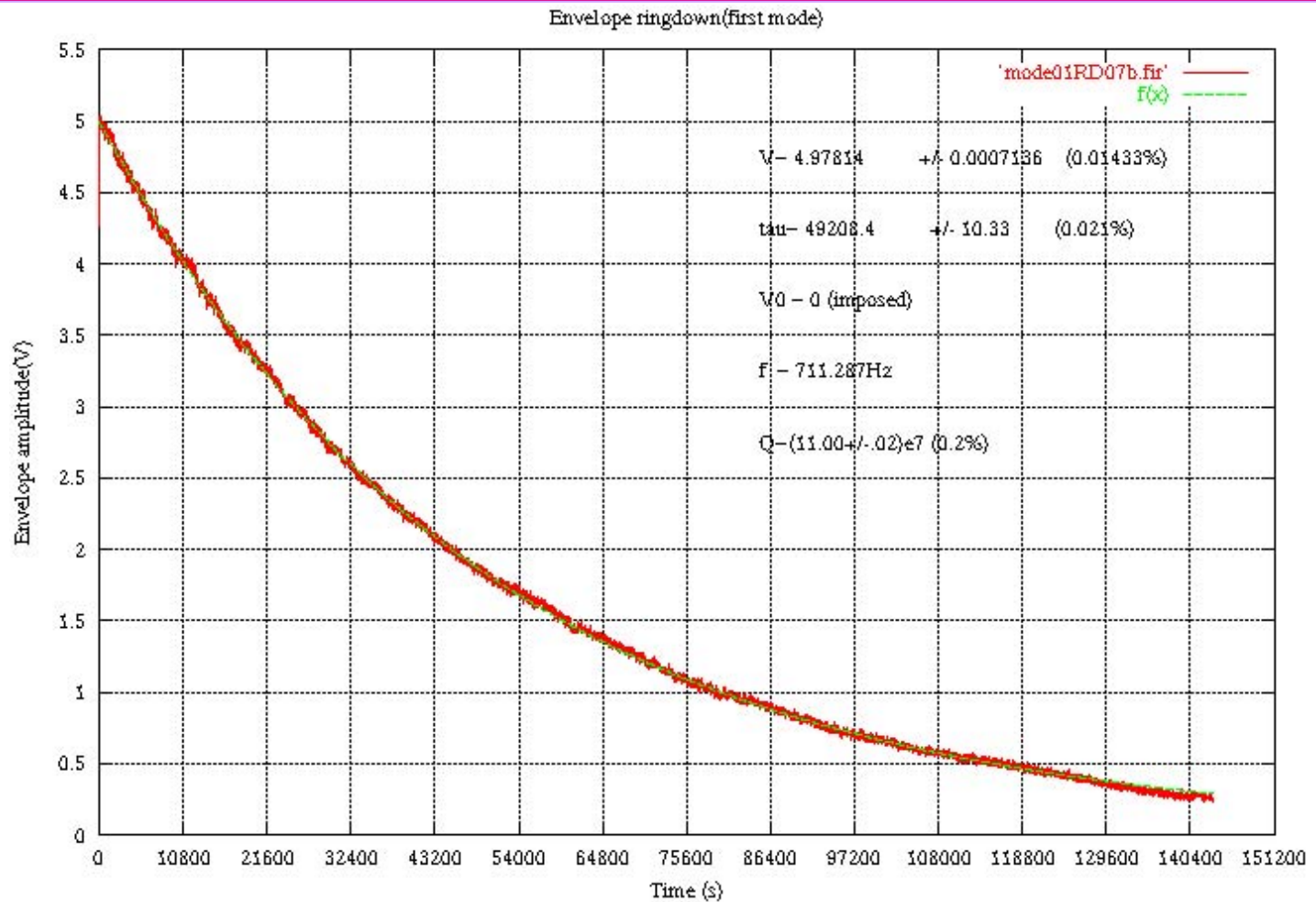
# Removal of Sidebands by Digital Filtering



# More Digital Filtering

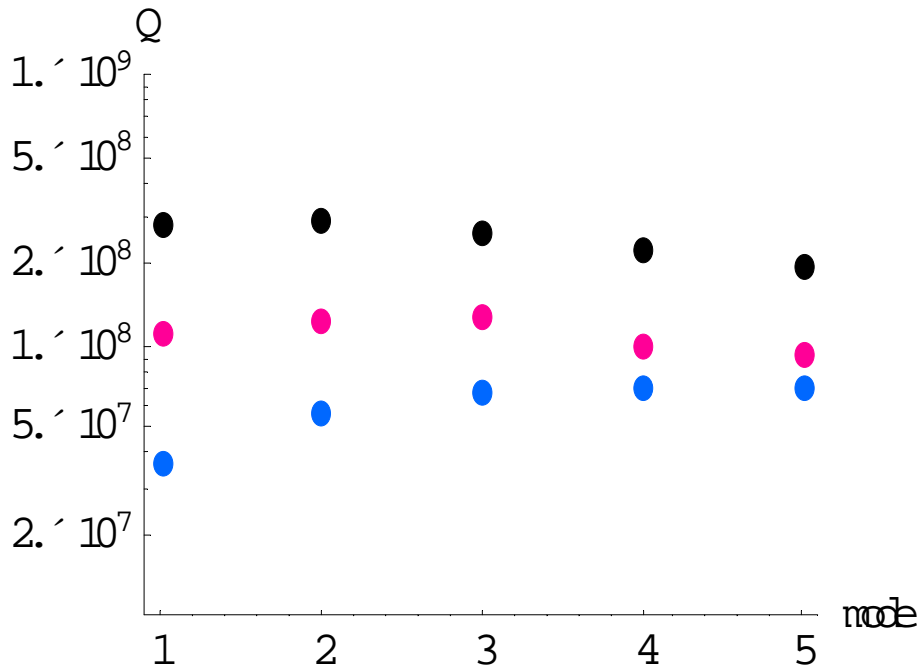


# Filtered Ringdown Data



# Our Data

- NTE
- LTE
- data



- Bear in mind that the predictions of the LTE model are *upper limits only*.
- ‘NTE’ prediction based on published values for  $\phi_{\text{bulk}}$ ,  $\phi_{\text{surface}}$ .
- We believe this to be the first observation of nonlinear thermoelastic damping.
- Future work:
  - » Measure  $\phi_{\text{bulk}}$ ,  $\phi_{\text{surface}}$  in unloaded fiber
  - » Measure Q of vertical bounce mode

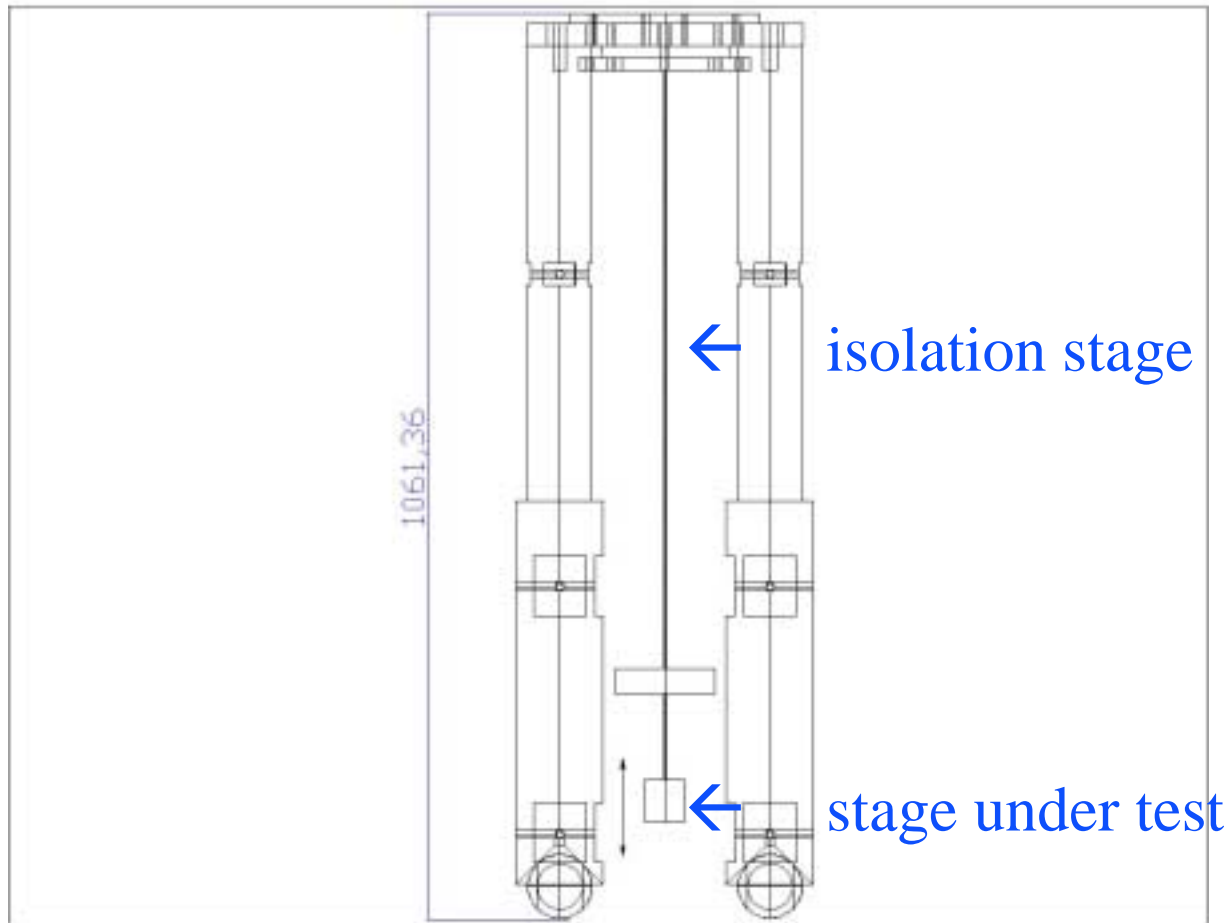
# Why Study Vertical Bounce Mode Q's?

- The vertical bounce mode probes the entire fiber uniformly, unlike the pendulum and violin modes which weight the endpoints more heavily.
- The vertical bounce mode has no dilution factor.
- The vertical bounce mode has no thermoelastic damping.
- Thus the vertical bounce mode permits the clean study of the internal friction of fused silica at high strain.

# Internal Friction at High Strain in Fused Silica

- So far, no good data exists at strains relevant for Advanced LIGO ( $u \sim .005-.01$ ):
  - » Braginsky et al. measure  $Q$  consistent with  $\phi = 1.4e-6$  at  $u = .005$
  - » Data above consistent with  $\phi = 7.8e-7$  at  $u = .003$
  - » Rowan et al. measure data consistent with  $\phi = 7e-7$  at  $u = .0002$
  - » Rowan et al. measure data consistent with  $\phi = 1.5e-6$  at  $u = .002$
  - » **Advanced LIGO baseline assumes  $\phi = 2e-7$  at  $u = .01$**
- Discrepancies could be due to recoil losses or contamination due to welding, but stress-dependent  $\phi$  cannot yet be excluded.
- Fused silica is getting nonlinear at these strains... Young's modulus changes by  $\sim 1e-2$  at  $u = .003$  (note:  $\phi$  is only  $1e-6$ !).
- It is prudent therefore to check stress-dependence of  $\phi$

# Vertical Bounce Measurement Apparatus- in Preparation





# Expected Sensitivity

- Use of monolithic fused silica upper suspension should contribute negligibly to loss from that stage...
- Fundamental recoil limit should be through suspension mounting structure- previous rigidities achieved at Glasgow should allow  $1/Q_{\text{recoil}} \sim 1e-7$  for 50um fibers holding 200g
- Hopefully new ultraheavy pendulum laboratory will permit larger fibers to be tested.