
New Folder Name Cryo Pump Infrared

Shielding Considerations

T950087

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CALIFORNIA INSTITUTE OF TECHNOLOGY
Laser Interferometer Gravitational Wave Observatory (LIGO) Project

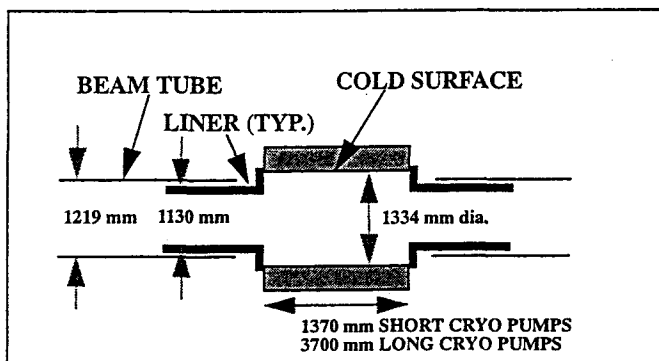
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Refer to: LIGO-T950087-00-E
Date: 18 Nov 95

Subject: Cryopump Infrared Shielding Considerations

Background

PSI has proposed¹ to reduce the heat load (and LN₂ consumption rate) of the cryopumps through the use of cylindrical liners (shields) with low emissivity ($\epsilon = 0.06$), diffuse surfaces on both sides of the cold surface, as indicated in Figure 1. PSI's calculations indicated a reduction in heat load to the cold surface of the long cryopump of approx. 33% is possible by using 1.52 m long liners on both sides of the long cryopump. The reduction in heat load is important in reducing the frequency of LN₂ tank refilling, reducing the size and expense of the LN₂ tank and reducing the cost of LN₂ during operation.

Figure 1: Cryopump Liners (Shields)



The "long" cryopumps have a 3.7 m long cylindrical cold surface and the "short" cryopumps have a 1.37 m long cylindrical cold surface. The internal diameters of the cold surfaces of both pumps are 1.334 m. PSI's current cryo pump design constrains the maximum diameter of the liners to be 1.13 m (PSI drawing V049-4-004, Rev. P1). The lengths of the liners for the cryo pumps are constrained by the placement of the adjacent gate valves; There appears to be some freedom in locating these gate valves relative to the cryo pump so as to achieve liner lengths on the order of 1 to 2 meters in length.

1. Process Systems International (PSI), LIGO Vacuum Equipment Preliminary Design, Vol. 1, Attachment 1.4, "80K Pump Performance", PSI-VE001AA1A01, 19 Jun 95.

Cryopump Infrared Shield Performance Analysis

An independent analysis of the effectiveness of the cryopump liner was undertaken due to concerns about the validity of two assumptions by PSI:

- 1) the entire beam tube (including the tube in proximity to the shield) acts like a black body, and
- 2) one can obtain a surface which behaves like a diffuse (Lambertian) IR reflector with low emissivity

The Lambertian characteristic of the surface is crucial; If the surface behaved as a specular reflector, then it would in essence simply mirror the black body of the beam tube to the cold surface and it would have no effect.

As reported previously¹, a view factor and radiation balance calculation² confirmed PSI's calculations of the mitigation in radiative heat load to the cold surface due to diffusive, low emissivity liners with the following parameters:

- 1) long cryopump (3.7 m cold surface length, and 1.2 m diameter³)
- 2) beam tube emissivity of 0.5
- 3) liner emissivity of 0.06 (diffuse)
- 4) liner length of 1.52 m

These calculations were based upon view factors for a coarse discretization of the geometry (i.e. a single view factor for each component). The calculations confirm that the simplification of treating the beam tube as a black body is reasonable.

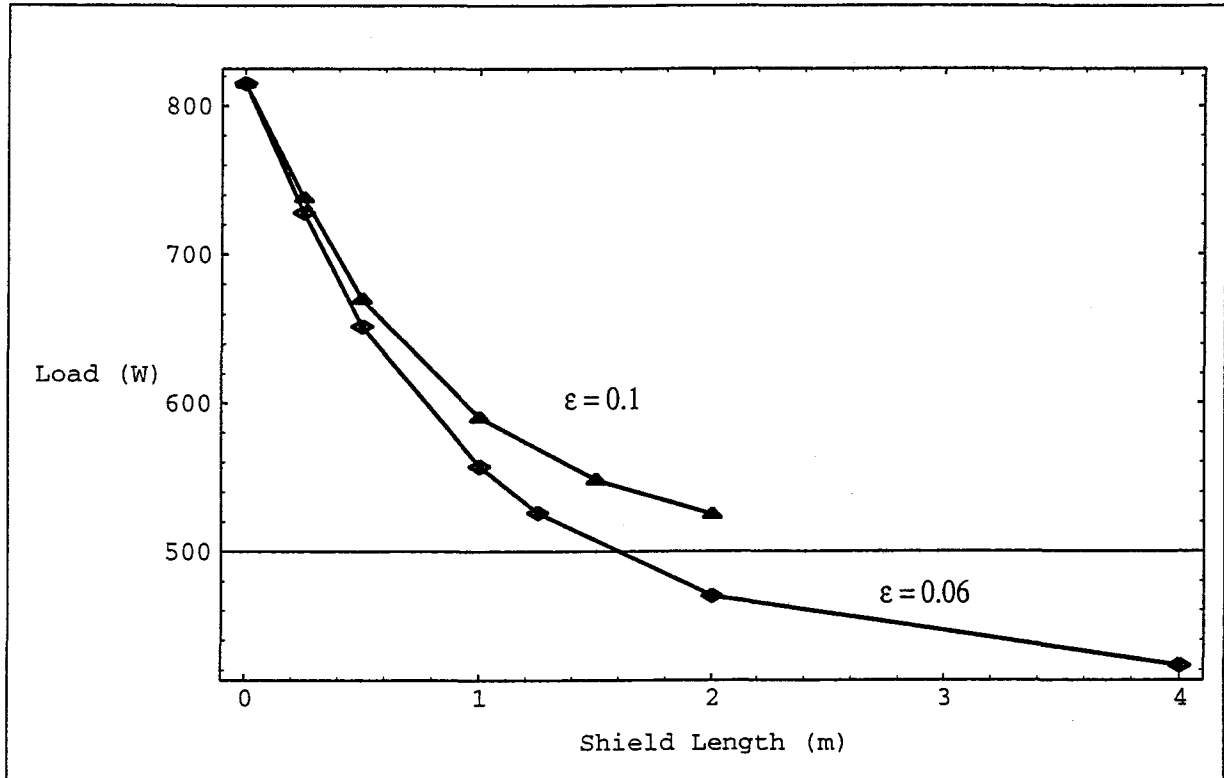
Further calculations have determined the effectiveness of the liners for the revised cryo pump diameter with a finer discretization to account for the gradient in radiosity. In the analysis the geometry of Figure 1 was approximated by assuming that the beam tube, liners and cryopump cold surface are all at the same diameter (the diameter of the beam tube, 1.2 m). The analysis will be updated for a subsequent revision of this memorandum: I did not want to delay release of this memo and the discussion which will ensue of the implications of the integration of the liner and laser light baffles. In addition, the analysis does not account for the interaction of the liners with the laser light baffles (to be discussed in the next section). On the basis of this analysis, the performance of the infrared shields for the long and short cryopumps as a function of shield length and parameterized by shield emissivity, are given in Figures 2 and 3 respectively. Although the performance will change somewhat when the diameters are revised, the results will not be substantially different.

Further extension to the radiative transport analysis of the liner effectiveness to account for the different diameters of the tube, liner and cold surface as well as inclusion of the laser light baffles, can be accomplished readily. As an alternative which would permit analysis of more complex geometries, I attempted to use the TMG module within the I-Deas Master Series 2 computer aided

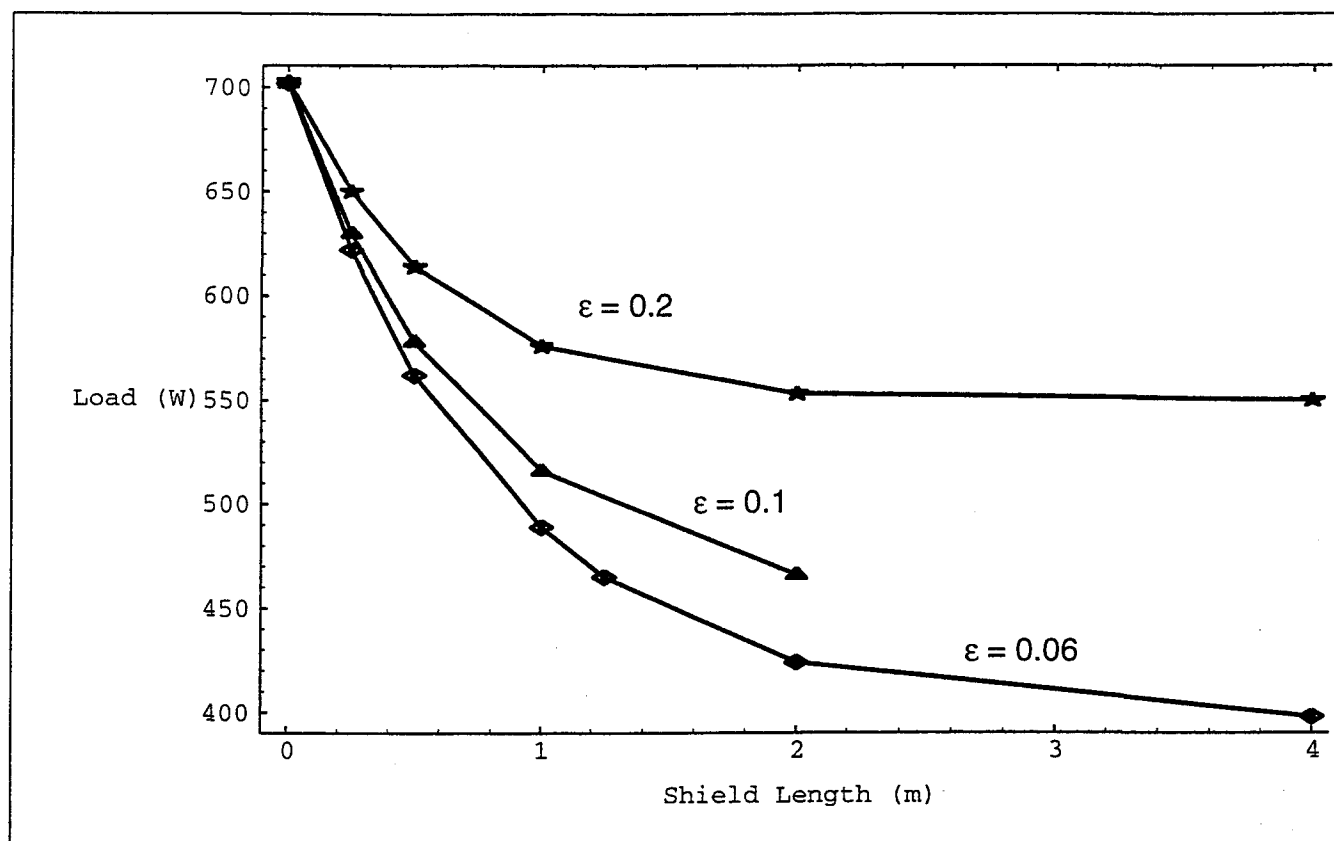
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1. Coyne, D., "Cryo Pump Shield Performance Analysis", e-mail to M. Zuker, A. Lazzarini and J. Worden, 2 Nov 95.
 2. Siegel, H. and Howell, J., Thermal Radiation Heat Transfer, 2nd ed., McGraw Hill, 1972, pp.236-248.
 3. The diameter of the cryopumps was 1.2 m at the time of PSI's initial PDR (June 95). The diameter has since been increased to 1.334 m based upon laser light baffling requirements as discussed in LIGO-L950593-01-E.

design and analysis package for analysis of specular as well as diffuse components of reflection. Although TMG is capable of handling combined diffuse and forward specular reflections which are independent of incidence angle, TMG cannot handle retro-reflection. Given that suitable nearly Lambertian surfaces are obtainable (as discussed in the next section), TMG may be adequate for any further analysis with more complex geometry (if necessary). Inclusion of a complete BRDF characterization in the analysis would require another code, perhaps TRASYS or CODE-V.

**Figure 2: Long Cryopump Liner (Infrared Shield) Performance
(for liner emissivities of 0.06 and 0.10)**



**Figure 3: Short Cryopump Liner (Infrared Shield) Performance
(for liner emissivities of 0.06, 0.10 and 0.20)**



Diffuse Infrared Shield Material

A quick check in the literature and discussions with Dr. Bob Breault (of Breault Research Organization, Inc.) indicates that Flame Sprayed Aluminum (FSA) or gold on FSA can meet the requirements for a low emissivity, diffuse reflector in the infrared. Attached are two references^{1,2}, with BRDF measurements on bare FSA and gold coated FSA. In the second reference, the MIL-STD used in the fabrication of the FSA surface is cited. The directional-hemispherical reflectance at 10.6 microns is approximately 0.85 for bare FSA and 0.95 for gold coated FSA. Both have very flat BRDFs indicating diffuse (Lambertian) behavior for low angles of incidence (as indicated in the attached). As the angle of incidence increases, these materials deviate more from a perfect Lambertian surface with a slight retro-component of reflection and a slight forward scatter (specular) component. However, the deviations from ideal Lambertian behavior are slight (about +/- 15%). The retro-reflection helps in reducing the infrared flux from the beam tube, but increases

1. Oppenheim, U., Turner, M., and Wolfe, W., "BRDF Reference Standards for the Infrared", Infrared Phys. Technol., Vol. 35, No. 7, pp. 873-879, 1994.
2. Brennan, W. (Hughes Electro-Optical & Data Systems Group) to Breault, R., letter with attached data on BRDF of flame sprayed aluminum at 5 deg. angle of incidence and 3 wavelengths (1.06, 3.39 and 10.6 microns), 23 Oct 91.

the coupling of emission from the absorptive side of the laser light baffles placed adjacent to the cold surface.

Considerations Related to the Integration of Laser Light Baffles and Infrared Shields

The minimum diameter of the LN₂ trap based upon considerations of optical shielding of the cold surface from view of the test masses has been addressed in a memorandum by Albert Lazzarini¹. This analysis does not account for the presence of the liners. For the same reasons outlined in this memorandum, the shields (or liners) should not be visible to the test masses, as indicated in Figure 5 (a to-scale drawing).

Stan Whitcomb has pointed out that PSI's proposal to use a diffusive shield in order to reduce the thermal load on the cryo-pump extends the effective length of the cryo-pump considerably resulting in some special considerations in combined IR and optical shielding. The dimensions indicated in the Figure 5 are for the most restrictive baffling situation, i.e. the short cryopumps at the mid- and end-stations. The diameters of the cryopump and its shields result in restricting the maximum length of the liners to about 0.8 m and 1.0 m (for the liner closest to, and farthest from, the test mass respectively) IF no laser light baffles are placed within the infrared liners.

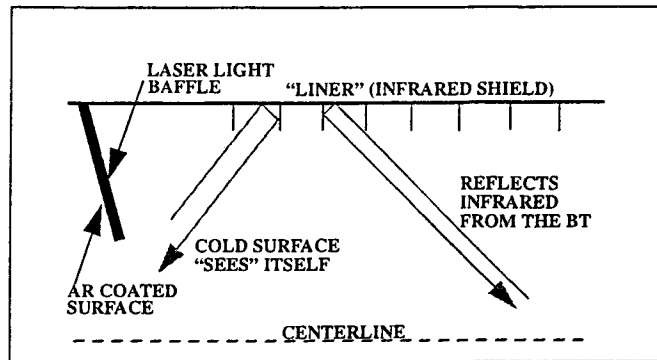
Due to the asymmetric nature of the problem, the two baffles placed between the cold surface and the test mass (Figure 5) serve to absorb laser light on one side and act as a low emissivity infrared radiator on the opposite side. The baffle placed on the side of the cold surface most distant from the test mass (and used to block direct view of the far liner) must be absorptive at the laser wavelength (and therefore absorptive in the infrared as well). Consequently, this baffle will compromise the ideal performance of the liners. However, the view factor from the cold surface to this baffle is not large (as can be seen from the scale drawing, Figure 4). A single baffle to block view of the cold surface and the far liner is impractical (diameter > 1.45 m for a liner of only 0.5 m in length).

An alternative to the use of diffuse liners (suggested by Stan Whitcomb), is to use a retro-reflection surface, as indicated in Figure 4. The cold surface would see itself and the self-emission of the (low emissivity) surface of the liner; The infrared emission of the beam tube would be retro-

1. Lazzarini, A., "Determination of the minimum LN₂ trap inner diameter needed to shield it from direct line of sight from a test mass", LIGO-L950593-01-E, 2 Nov 95.

reflected back down the tube. The performance of this approach is yet to be determined.

Figure 4: Cryopump Liners (Shields)



In either approach for infrared shielding, the laser light baffles adjacent to the cryo pumps would have to be treated differently from the baffles in the beam tube. Stan suggests an broad-band AR coating (from 1 to 10 microns).

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Attachments:

- 1) Drawing D95~~xxx~~⁰¹⁶²-SK (11/18/95), Cryopump Liner and Baffle Considerations
- 2) Cryopump Radiant Interchange Analysis (derivation and Mathematica Notebook)
- 3) Oppenheim, U., Turner, M., and Wolfe, W., "BRDF Reference Standards for the Infrared", Infrared Phys. Technol., Vol. 35, No. 7, pp. 873-879, 1994.
- 4) Brennan, W. (Hughes Electro-Optical & Data Systems Group) to Breault, R., letter with attached data on BRDF of flame sprayed aluminum at 5 deg. angle of incidence and 3 wavelengths (1.06, 3.39 and 10.6 microns), 23 Oct 91.

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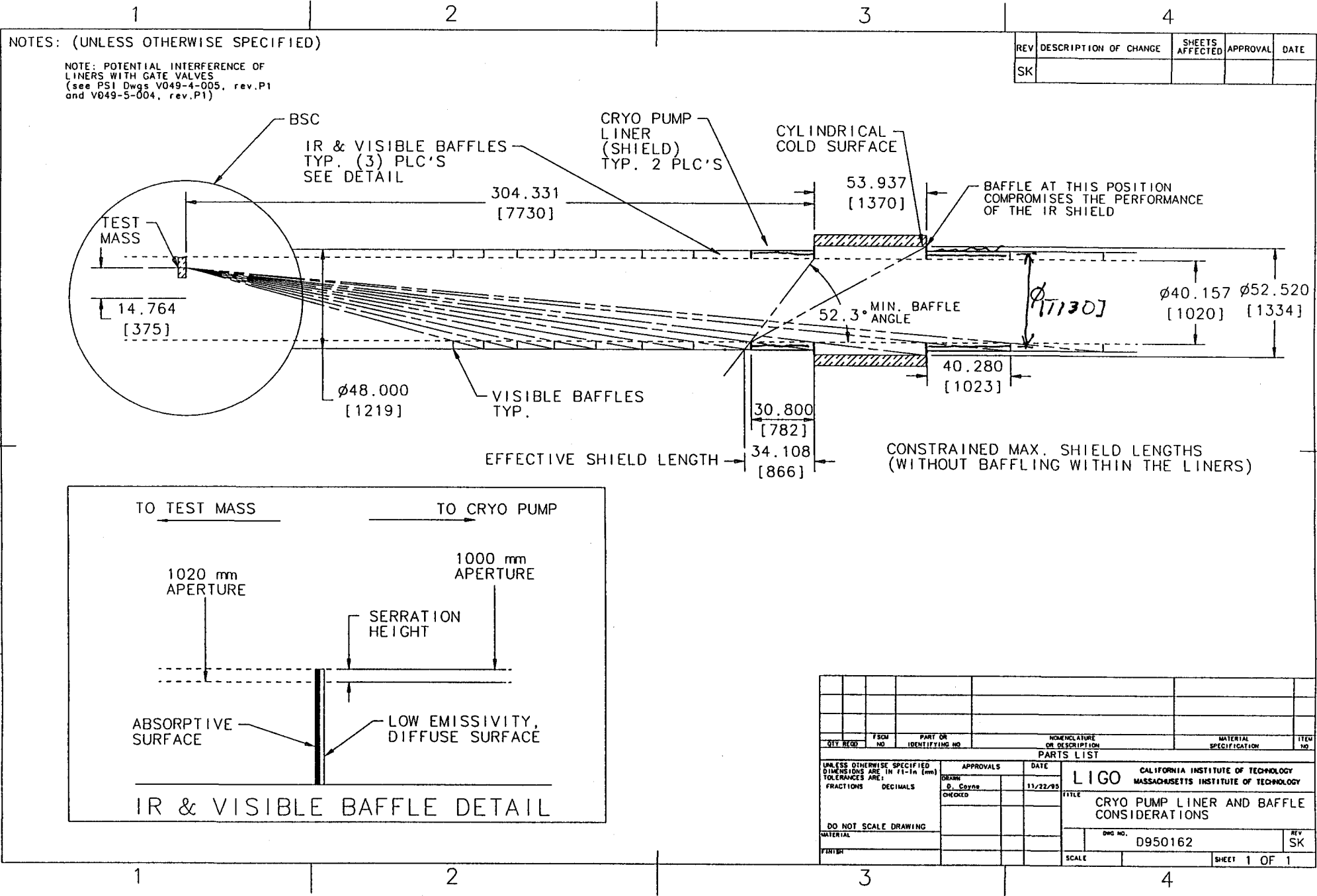
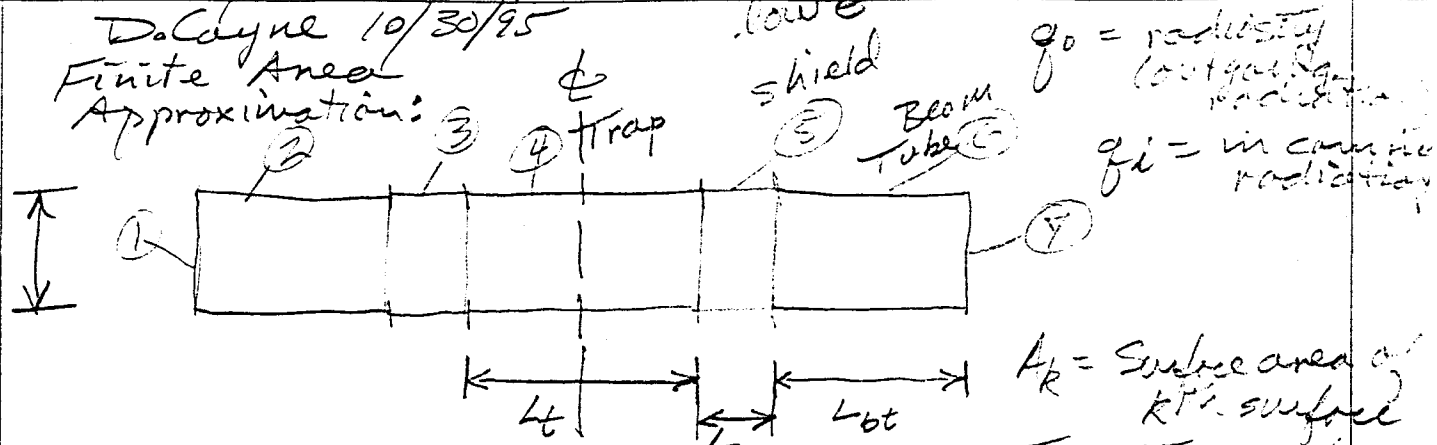


Figure 5.

CRYO PUMP RADIANT INTERCHANGE

DoCayne 10/30/95

Finite Area Approximation:



$q_0 =$ radiosity (outgoing radiation)
 $q_i =$ incoming radiation

$A_k =$ Subarea of k^{th} surface
 $T_k =$ Temperature of k^{th} surface

① $Q_k = \frac{q_k}{f_k} A_k = (q_{0,k} - q_{i,k}) A_k$

② $q_{i,k} = \epsilon_k \sigma T_k^4 + \rho_k q_{i,k} = \epsilon_k \sigma T_k^4 + (1 - \epsilon_k) q_{i,k}$

$$A_k q_{i,k} = A_1 q_{0,1} F_{1-k} + A_2 q_{0,2} F_{2-k} + \dots + A_j q_{0,j} F_{j-k} + \dots + A_k q_{0,k} F_{k-k} + \dots + A_N q_{0,N} F_{N-k}$$

or
 ③ $q_{i,k} = \sum_{j=1}^N F_{k-j} q_{0,j}$ (using the config factor reciprocity relation $A_j F_{j-k} = A_k F_{k-j}$)

from ②:
 $q_{i,k} = \left(\frac{1}{1 - \epsilon_k} \right) (q_{0,k} - \epsilon_k \sigma T_k^4)$
 combining with ①:

$$Q_k = A_k \left[q_{0,k} + \left(\frac{1}{1 - \epsilon_k} \right) (\epsilon_k \sigma T_k^4 - q_{0,k}) \right]$$

④ $Q_k = A_k \left(\frac{\epsilon_k}{1 - \epsilon_k} \right) [\sigma T_k^4 - q_{0,k}]$

combining ① and ③:

⑤ $Q_k = A_k \left(q_{0,k} - \sum_{j=1}^N F_{k-j} q_{0,j} \right)$

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Equations (2) and (3)

$$q_{o,k} - \sum_{j=1}^N F_{k-j} q_{o,j} = \left(\frac{\epsilon_k}{1-\epsilon_k} \right) [\sigma T_k^4 - q_{o,k}]$$

$$q_{o,k} - (1-\epsilon_k) \sum_{j=1}^N F_{k-j} q_{o,j} = \epsilon_k \sigma T_k^4$$

$$\textcircled{6} \quad \sum_{j=1}^N [\delta_{kj} - (1-\epsilon_k) F_{k-j}] q_{o,j} = \epsilon_k \sigma T_k^4$$

$$\textcircled{7} \quad [a] \{q_o\} = \{c\}$$

$$\textcircled{8} \quad \{q_o\} = [a]^{-1} \{c\}$$

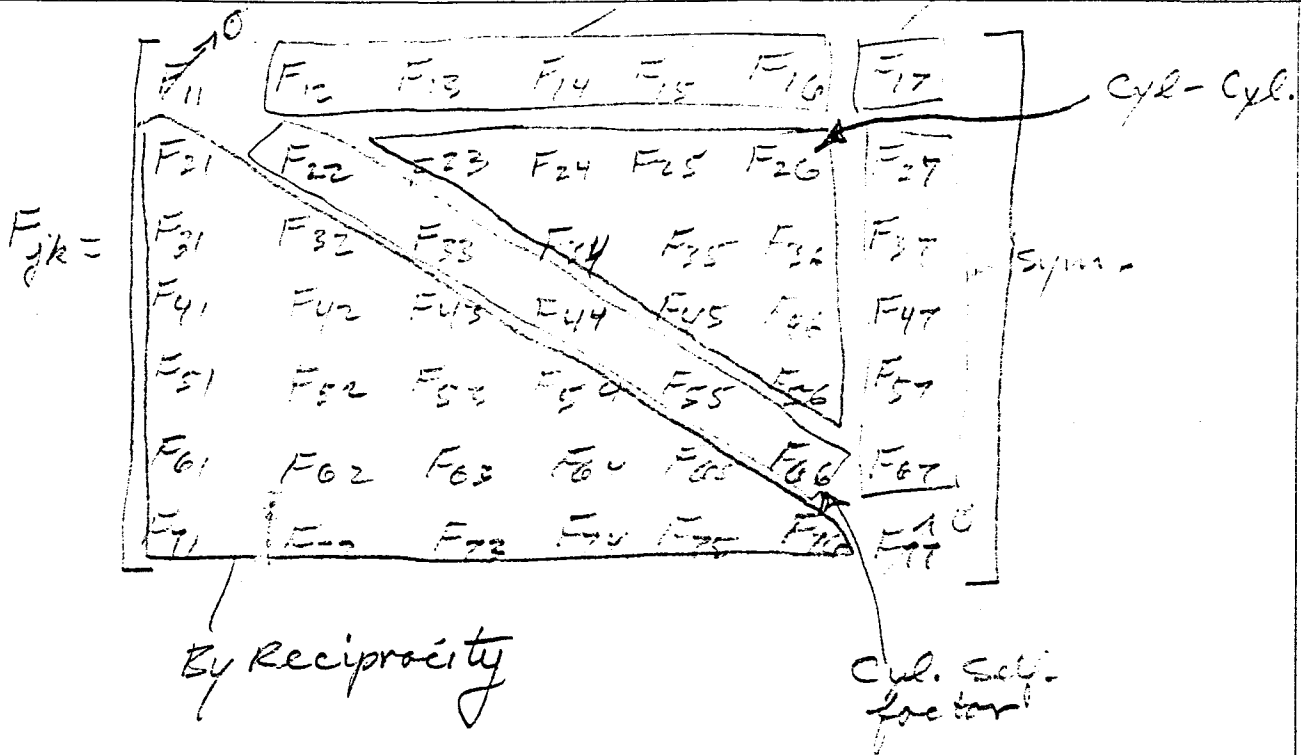
Use

$$\textcircled{5} \quad Q_k = A_k \left(q_{o,k} - \sum_{j=1}^N F_{k-j} q_{o,j} \right)$$

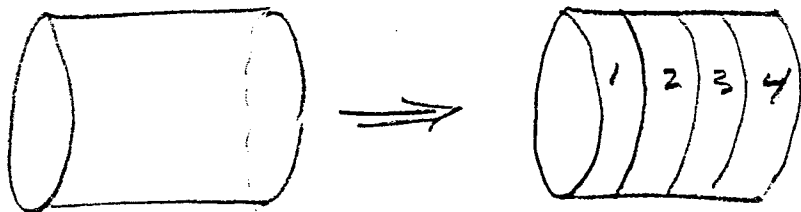
to project out the net heat balance at each surface

disk-cyl.

disk-disk



Discretization of the component geometries is straight forward:



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CRYO PUMP RADIANT INTERCHANGE

■ Cryo Pump Properties

■ Component Dimensions

In[2790]:=

```
Dia = 1.22;      (* m *)  
Ltrap = 3.7;    (*large cryo-pump, m *)  
Lshield = 1.5;  (* m *)  
Lbt = 20;      (* m *)
```

In[2794]:=

```
radius = Dia/2;  
L = Ltrap + 2 (Lshield + Lbt);    (* overall length, m *)
```

■ Component Emissivities

In[2796]:=

```
eend = 0.9;      (* emissivity of the bt ends *)  
ebt = 0.5;      (* emissivity of the beam tube *)  
eshield = 0.06; (* emissivity of the shields *)  
etrap = 1.0;    (* emissivity of the cold trap *)
```

■ Component Temperatures (K)

In[2800]:=

```
Tend = 295;  
Tbt = 295;  
Tshield = 295;  
Ttrap = 80;
```

■ Cryo Pump Model Parameters

■ Segmentation

In[2804]:=

```
ntrap = 10;      (* number of segments for the cold trap surface *)
nshield = 10;   (* number of segments for each shield surface *)
nbt = 20;      (* number of segments for each beam tube surface *)
n = 2 + 2 (nbt + nshield) + ntrap
```

Out[2807]=

72

■ Segmented Surface Area, Position, Length, Emissivity & Temperature

In[2808]:=

```
A = Table[0, {n}];
emiss = A;
T = A;
x = A;
s = A;
A[[1]] = Pi Dia^2/4;
A[[n]] = A[[1]];
emiss[[1]] = eend;
emiss[[n]] = emiss[[1]];
T[[1]] = Tend;
T[[n]] = T[[1]];
x[[1]] = 0;
x[[n]] = L;
s[[1]] = 0;
s[[n]] = 0;
```

In[2823]:=

```

As = Pi Dia Lbt/nbt;
es = ebt;
Ts = Tbt;
ss = Lbt/nbt;
Do [
  A[[i]] = As;
  emiss[[i]] = es;
  T[[i]] = Ts;
  s[[i]] = ss;
  x[[i]] = x[[i-1]] + s[[i-1]];
  A[[i+nbt+ntrap+2 nshield]] = As;
  emiss[[i+nbt+ntrap+2 nshield]] = es;
  T[[i+nbt+ntrap+2 nshield]] = Ts;
  s[[i+nbt+ntrap+2 nshield]] = ss;
  x[[i+nbt+ntrap+2 nshield]] = x[[i]] + Lbt + Ltrap + 2 Lshield,
  {i,2,nbt+1}
];
As = Pi Dia Lshield/nshield;
es = eshield;
Ts = Tshield;
ss = Lshield/nshield;
Do [
  A[[i]] = As;
  emiss[[i]] = es;
  T[[i]] = Ts;
  s[[i]] = ss;
  x[[i]] = x[[i-1]] + s[[i-1]];
  A[[i+nshield+ntrap]] = As;
  emiss[[i+nshield+ntrap]] = es;
  T[[i+nshield+ntrap]] = Ts;
  s[[i+nshield+ntrap]] = ss;
  x[[i+nshield+ntrap]] = x[[i]] + Lshield + Ltrap,
  {i,nbt+2,nbt+nshield+1}
];
As = Pi Dia Ltrap/ntrap;
es = etrap;
Ts = Ttrap;
ss = Ltrap/ntrap;
Do [
  A[[i]] = As;
  emiss[[i]] = es;
  T[[i]] = Ts;
  s[[i]] = ss;
  x[[i]] = x[[i-1]] + s[[i-1]],
  {i,nbt+nshield+2,nbt+nshield+ntrap+1}
];

```

■ View Factors

```
In[2838]:=
  F = Table[0, {n}, {n}];
```

■ Disk-Disk View Factors

N.B.: Concentric, r1 and r2 are radii, h = distance between disks

```
In[2839]:=
  Fdd[r1_, r2_, h_] := If[h == 0, 1,
  With[{X = 1 + (1 + (r2/h)^2) / (r1/h)^2}, 1/2 (X - Sqrt[X^2 - 4
  (r2/r1)^2])]];

```

Left End Disk(1) - Right End Disk(n):

```
In[2840]:=
  F[[1, n]] = Fdd[Dia/2, Dia/2, L];
```

■ Disk-Interior Cylinder View Factors

N.B.: Concentric, r1 = disk radius, r2 = cylinder radius, r1 <= r2, h1 = distance from disk to near end of cylinder, h2 = distance from disk to far end of cylinder

```
In[2841]:=
  Fdc[r1_, r2_, h1_, h2_] := Fdd[r1, r2, h1] - Fdd[r1, r2, h2];
```

Left End Disk(1) - Cylindrical Segments (2->n-1)
and by symmetry, Cylindrical Segments (2->n-1) - Right End Disk(n):

```
In[2842]:=
  Do [
    F[[1, i]] = Fdc[radius, radius, x[[i]], x[[i]] + s[[i]]];
    j = n - i + 1;
    F[[j, n]] = F[[1, i]] A[[n]] / A[[j]],
    {i, 2, n - 1}
  ];
```

■ Interior Cylinder Self-View Factors

N.B.: r = radius, l = cylinder length

```
In[2843]:=
  Fcself[r_, l_] := 1 - (r/l) Fdc[r, r, 0, l];
```

Cylindrical Segments (i) - Cylindrical Segments (i), for $i=\{2,n-1\}$:

```
In[2844]:=
Do [
  F[[i,i]] = Fcself[radius,s[[i]],
    {i,2,n-1}
];
```

■ Interior Equal-Diameter Concentric Cylinder View Factors

N.B.: r = radius, l_1 = cylinder 1 length, l_2 = cylinder 2 length, h = separation distance

```
In[2845]:=
Fceq[r_,l1_,l2_,h_] := (Fdc[r,r,h,l1+h] - Fdc[r,r,l2+h,l2+l1+h]) (r/(2 l
```

Cylindrical Segments (i) - Cylindrical Segments (j), for $i,j=\{2,n-1\}$:

```
In[2846]:=
Ftemp=Table[0,{n},{n}];

In[2847]:=
Do [
  Do [
    F[[i,j]] = Fceq[radius,s[[i]],s[[j]],x[[j]]-x[[i]]-s[[i]];
    Ftemp[[i,j]] = F[[i,j]],
    {j,i+1,n-1}
  ],
  {i,2,n-2}
];
```

■ Reciprocity

```
In[2848]:=
Do [
  Do [
    F[[i,j]] = F[[j,i]] A[[j]]/A[[i]],
    {j,1,i-1},
    {i,2,n}];
Do [
  Print[Sum[F[[i,j]],{j,1,n}],
    {i,1,n}];
```


■ Radiant Exchange Solution

```

In[2849]:=
  one = Table[1, {n}];
  a = IdentityMatrix[n] - Dot[DiagonalMatrix[(one - emiss)], F];

In[2851]:=
  stefanBoltzmann = 5.6696 10^-8; (* W/m^2/K^4 *)
  b = stefanBoltzmann emiss T^4;

In[2853]:=
  q0 = Inverse[a].b;

In[2854]:=
  q = q0 - F.q0 //N;

In[2855]:=
  Q = Dot[DiagonalMatrix[A], (q0 - F.q0)] //N;

In[2856]:=
  Qends = Q[[1]] + Q[[n]]

Out[2856]=
  0.302877

In[2857]:=
  Qbts = Sum[Q[[i]], {i, 2, nbt+1}] + Sum[Q[[i]], {i, nbt+ntrap+2,
  nshield+2, n-1}]

Out[2857]=
  363.647

In[2858]:=
  Qshields = Sum[Q[[i]], {i, nbt+2, nbt+nshield+1}]
  + Sum[Q[[i]], {i, nbt+nshield+ntrap+2, nbt+ntrap+2, nshield+1}]

Out[2858]=
  138.878

In[2859]:=
  Qtrap = Sum[Q[[i]], {i, nbt+nshield+2, nbt+nshield+ntrap+1}]

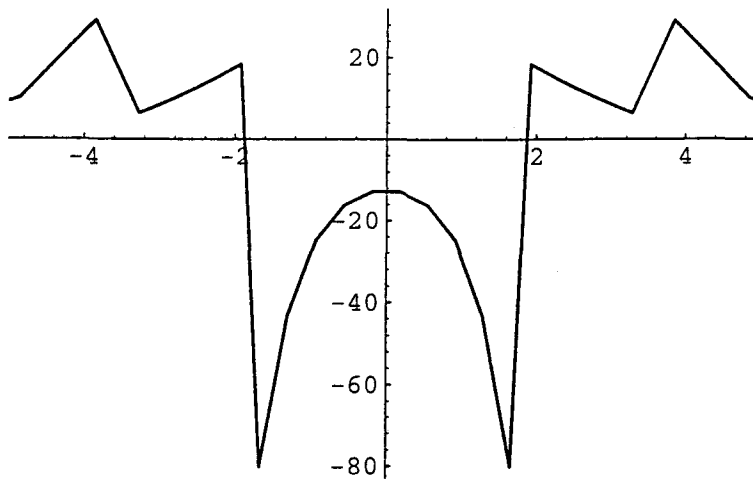
Out[2859]=
  -502.828

In[2860]:=
  xmid = x + s/2 - L/2;

```

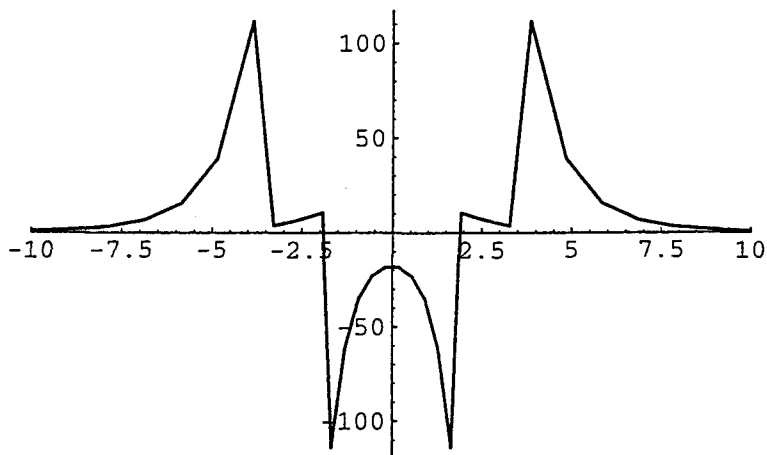
In[2861]:=

```
ListPlot[Transpose[Partition[Join[xmid,q],n]],PlotJoined->True,PlotRange->{{-5,5},All}];
```



In[2862]:=

```
ListPlot[Transpose[Partition[Join[xmid,Q],n]],PlotJoined->True,PlotRange->{{-10,10},All}];
```



■ Tabulated Results

■ Long Cryo-Pump

1.2 m diameter
 3.7 m long cold trap
 nominally 1.52 m long shields (liners)
 trap emissivity = 1.0 (ice)
 shield emissivity = 0.06 (diffuse)
 beam tube model length = 20 m

shield emissivity = 0.06

nbt = 20

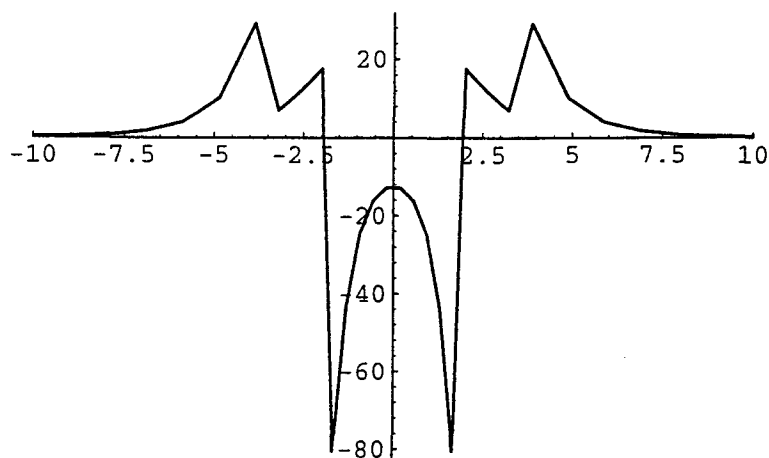
nshield = 5

ntrap = 10

n = 62

Qtrap = 504 W

q (W/m²) vs Xmid (m) :



shield emissivity = 0.5

nbt = 20

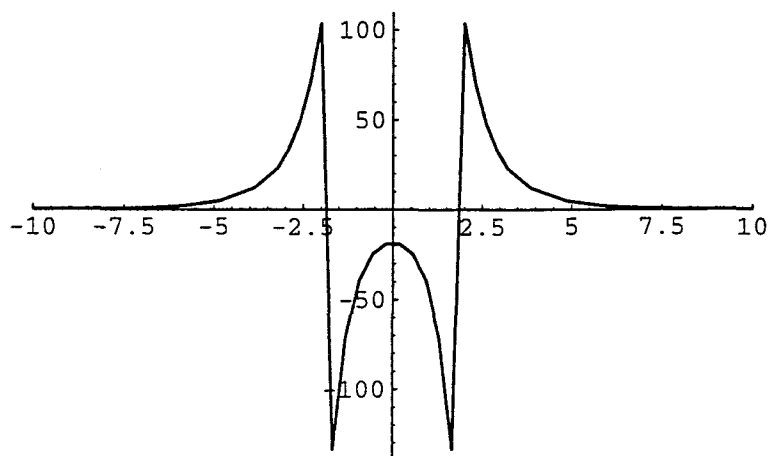
nshield = 5

ntrap = 10

n = 62

Qtrap = 817 W (823 W when Lshield == 3 m)

q (W/m²) vs Xmid (m) :



In[2863]:=

Needs ["Graphics`MultipleListPlot`"]

Used nshield = 10, for a no-shield resulting Q = 815 W in the following sensitivity calculations to shield emissivity and length:

```
In[2864]:=
```

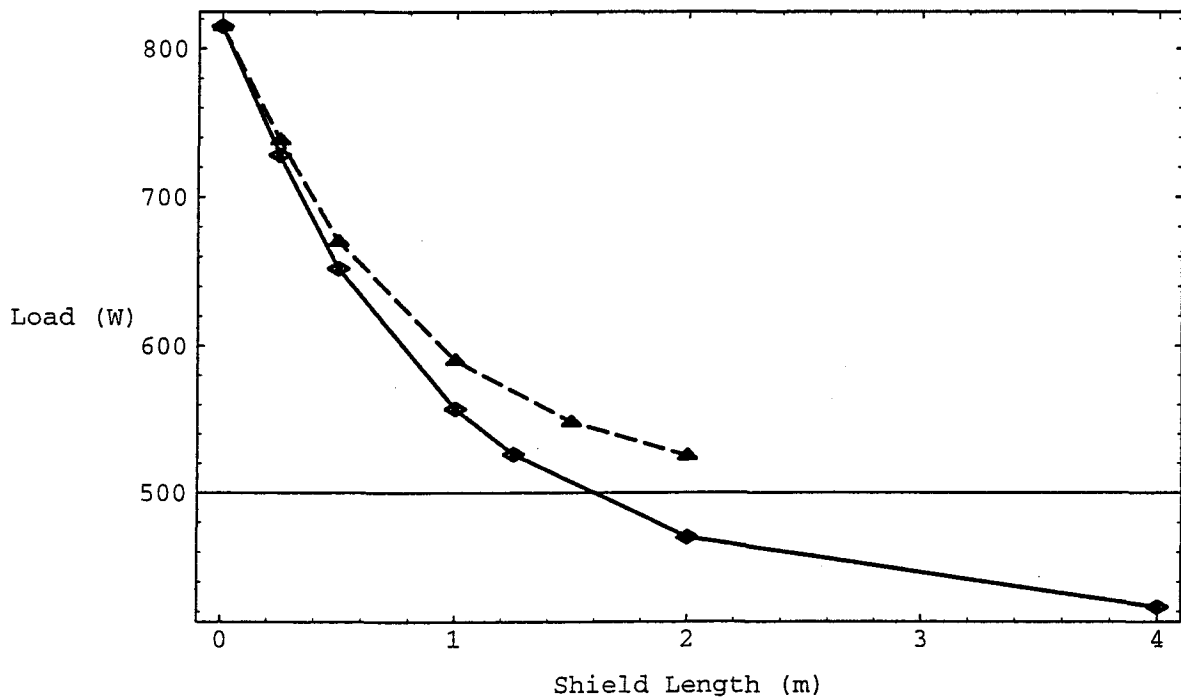
```
emissShield = {0.06, 0.10};
```

```
In[2865]:=
```

```
Qlong = {{{0,815}, {0.25,728}, {0.5,652}, {1.0,557},
{1.25,526}, {2,470}, {4,423}},
{{0,815},{0.25,738}, {0.5,670}, {1,590}, {1.5,548}, {2,525}}};
```

```
In[2866]:=
```

```
MultipleListPlot[Qlong[[1]],Qlong[[2]],PlotJoined-
>True,PlotRange->All,LineStyles-
>{{},{Dashing[{0.02,0.01]}},{Dashing[{0.03,0.01,0.01,0.01]}]},
Frame->True, RotateLabel->False, FrameLabel->{"Shield Length
(m)", "Load (W)"}];
```



Heat load (W) on the inner surface of the cryo-pump (cold trap) for the "long" (3.7 m) cryo-pump, as a function of shield length (m) parameterized by emissivity (0.06, 0.10 and 0.20).

■ Short Cryo-Pump

1.2 m diameter
 1.2 m long cold trap
 1 m long shields (liners)
 trap emissivity = 1.0 (ice)
 shield emissivity = 0.06 (diffuse)
 beam tube model length = 10 m

shield emissivity = 0.06

nbt = 20

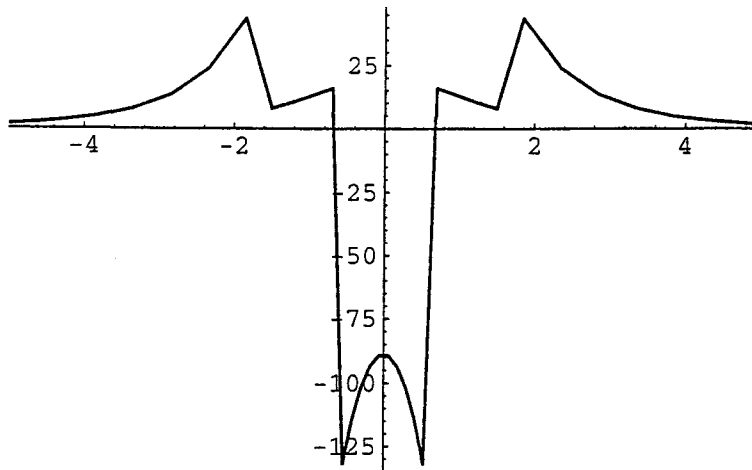
nshield = 5

ntrap = 10

n = 62

Qtrap = 489 W

q (W/m²) vs Xmid (m) :



shield emissivity = 0.5

nbt = 20

nshield = 5

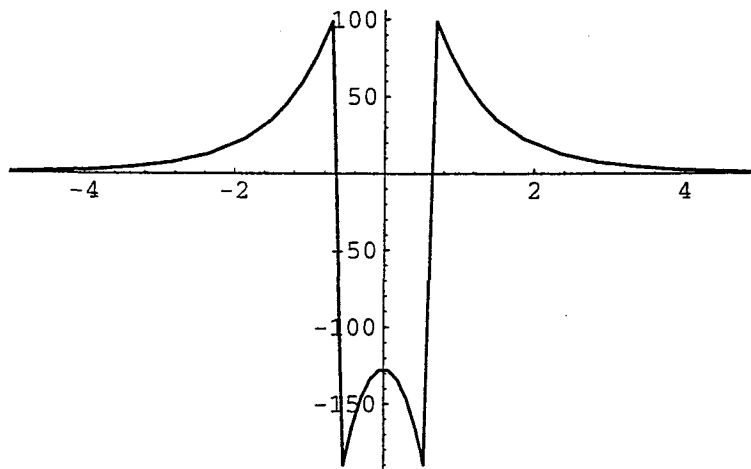
ntrap = 10

n = 62

Qtrap = 702 W (when Lshield == 1 m)

Qtrap = 708 W (when Lshield == 3 m)

q (W/m²) vs Xmid (m) :



Note: Used 10 segments for the shield when the emissivity of the shield was set to 0.2

In[2867]:=

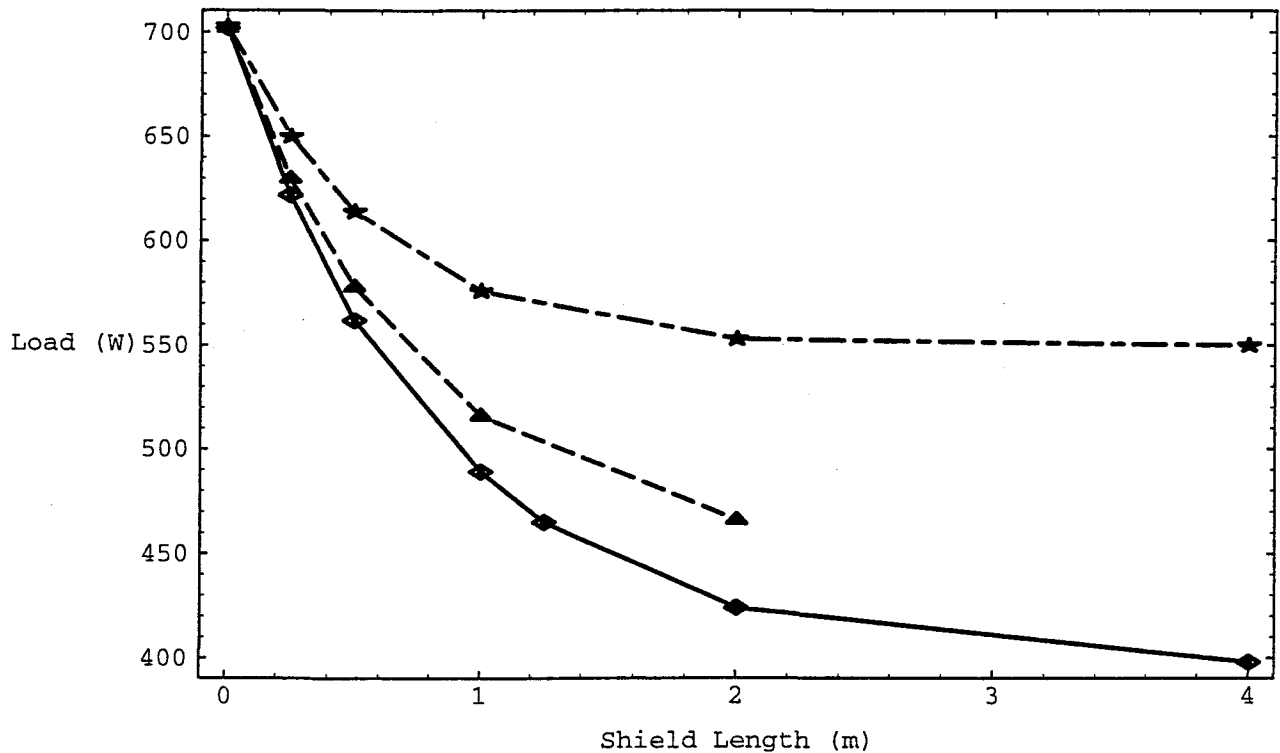
```
emissShield = {0.06, 0.10, 0.20};
```

In[2868]:=

```
Qshort = {{{0,702}, {0.25,622}, {0.5,562}, {1.0,489},
{1.25,465}, {2,424}, {4,398}},
{{0,702},{0.25,630}, {0.5,578}, {1,516}, {2,466}},
{{0,703}, {0.25,650}, {0.5,614}, {1,576}, {2,553}, {4,550}}};
```

In[2869]:=

```
MultipleListPlot[Qshort[[1]],Qshort[[2]],Qshort[[3]],PlotJoined-
->True,PlotRange->All,LineStyle-
>{{}, {Dashing[{0.02,0.01}]}, {Dashing[{0.03,0.01,0.01,0.01}]}}},
Frame->True, RotateLabel->False, FrameLabel->{"Shield Length
(m)", "Load (W)"}];
```



Heat load (W) on the inner surface of the cryo-pump (cold trap) for the "short" (1.2 m) cryo-pump, as a function of shield length (m) parameterized by emissivity (0.06, 0.10 and 0.20).

D. Coyne



Pergamon

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BRDF REFERENCE STANDARDS FOR THE INFRARED

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(Received 24 April 1994)

Abstract—A comparison between various recommended reference standards of diffuse reflectance in the IR is presented. It is shown that at a wavelength of 10.6 μm sulfur is the most Lambertian of the tested samples, although its powdery consistence makes it less suitable for use as a standard. Flame sprayed aluminum, with or without gold coating, also approaches a Lambertian surface and is suitable for use as a standard for BRDF measurements at 10.6 μm . Results for the BRDF of sulfur, gold-coated sandpaper, a commercial diffuse gold surface (by Labsphere) and flame sprayed aluminum are presented.

INTRODUCTION

Round Robin studies of the Bidirectional Reflectance Distribution Function (BRDF) of surfaces in the infrared^(1,2) have proved the need for better reference standards and improved measurement techniques in this region of the spectrum. The latest report⁽³⁾ shows differences in BRDF values of up to 50% between laboratories. There is a need for highly diffuse surfaces which are also highly reflecting in the IR and at the same time are sturdy and reproducible.

The present study is a comparison between several candidates for diffuse standards made with the AZSCAT [Arizona Scatterometer⁽⁴⁾] situated in the Optical Sciences Center of the University of Arizona. The source of radiation was a chopped cw CO₂ laser, polarized in the vertical direction which was always normal to the plane of incidence of the sample. The samples available for this study were sulfur, gold-coated sandpaper, a commercial diffuse gold surface and flame sprayed aluminum (bare and gold-coated). Angles of incidence were chosen at 10, 30, 45 and 60°. BRDF measurements for all these samples are presented.

Although no specular reflectance peak was observed for these samples (except the commercial sample), a distinct "retro" effect was observed, showing a hump in the BRDF curve at a backscatter angle equal to the angle of incidence. Slight departures from Lambertian behavior were observed at angles of incidence of 45 and 60°, although the surfaces were still highly diffuse.

MEASUREMENTS AND RESULTS

The following samples were available for the present study:

1. Flowers of sulfur, also known as sublimed sulfur powder. The sample was prepared by mixing the powder with acetone and compressing it in a tray, according to the "sulfur flooded" method described by Haner and Menzies.⁽⁵⁾ The sample had a diameter of 5 cm.
2. Gold-coated sandpaper. This surface was available from previous studies at the Optical Sciences Center and was described by Stuhlinger *et al.*⁽⁶⁾ A sample of 600 grit gold-coated sandpaper of 5 cm dia. was used.
3. A "certified reflectance standard" obtained from Labsphere (North Sutton, NH) of 5 cm dia., No. IRS-94-020. This consisted of a coarse sandblasted, gold electroplated surface.

[†]On ׁׂ׃ׄ (sabbatical leave) from the Technion-Israel Institute of Technology, Haifa, Israel.

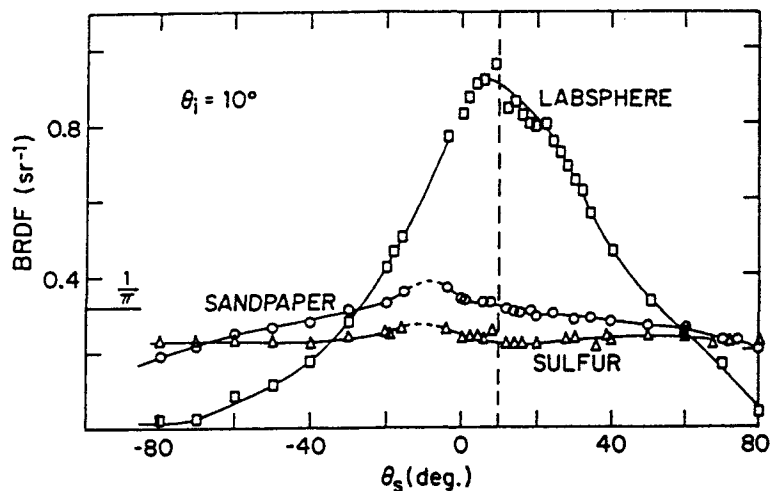


Fig. 1. BRDF of three samples at $10.6 \mu\text{m}$ for an angle of incidence of 10° .

4. Flame sprayed aluminum, made by flame spraying a flat aluminum surface, 5×5 cm in size, with an appearance that was much coarser than the Labsphere sample. Visual inspection of this surface showed many shiny dots which moved along the surface as the viewing angle was changed.
5. Gold-coated flame sprayed aluminum. This was obtained by having a sample of flame sprayed aluminum electroplated with pure gold. The plating was done by a process called "laser gold" by Epner Technology, Inc. (25 Division Place, Brooklyn, NY 11222).

All measurements were made with the AZSCAT instrument, which is fully automatic and is

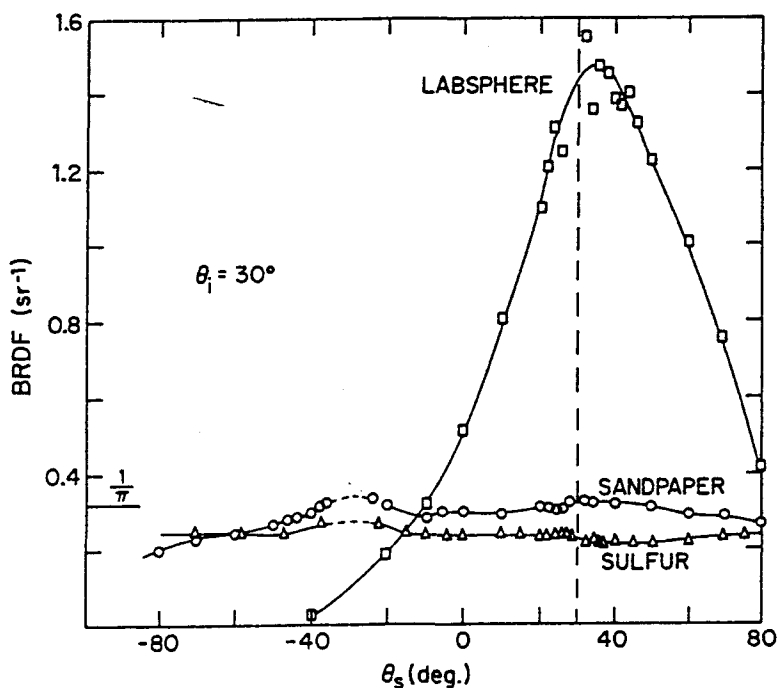


Fig. 2. BRDF of three samples at $10.6 \mu\text{m}$ for an angle of incidence of 30° .

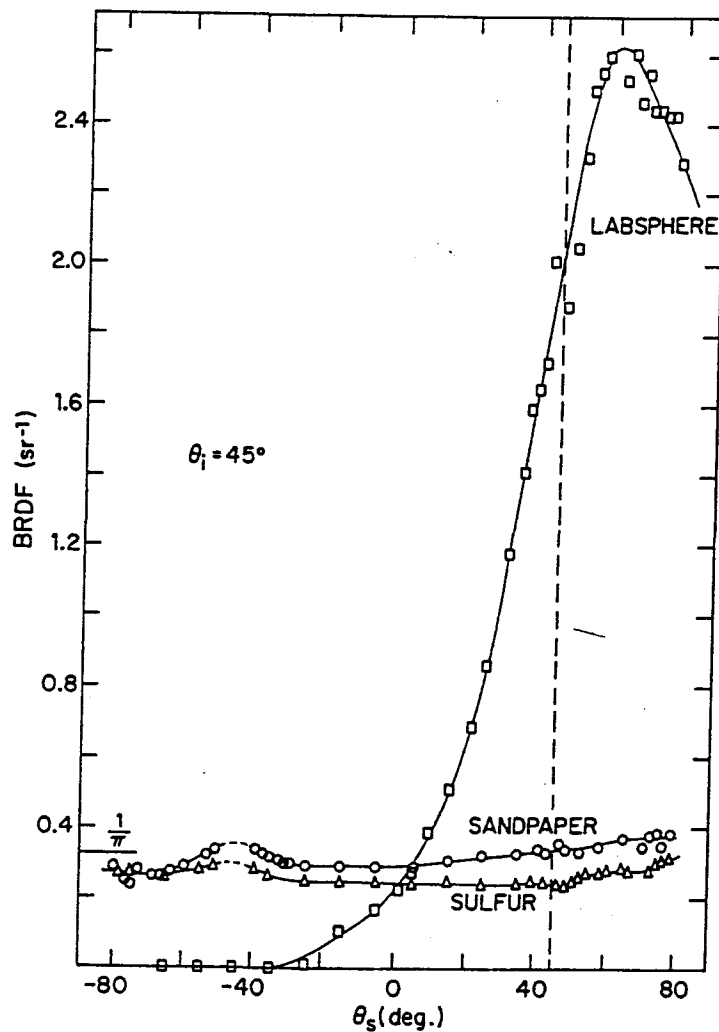


Fig. 3. BRDF of three samples at 10.6 μm for an angle of incidence of 45°.

based on the single reference method.⁽⁴⁾ The BRDF is calculated using the relation

$$\text{BRDF} = \frac{\rho}{\pi} \frac{V_i}{V_o \cos \theta_s}$$

where ρ is the directional-hemispherical reflectance, V_i is the output of the detector for the sample under study, V_o is the detector output for the reference sample and θ_s is the scattering angle with respect to the normal to the surface. Before a run of BRDF measurements of a certain sample was started, the gold-coated sandpaper was used as a reference by placing it in the beam at an angle of incidence of 10° and a scattering angle of 5°. The detector output in this position was taken as V_o in the runs that followed. It was assumed that the hemispherical reflectance of the gold sandpaper was 0.97 in this position, and this value was used in the above expression for the BRDF. Once this value was established the BRDF for any given sample and any pair of angles could be measured by the instrument.

All measurements were made with the incident and scattered beams in the plane of incidence of the sample, so that the azimuthal scattering angle was zero at all times. The instrument was capable of spinning the sample around an axis normal to the sample plane and this was done in all measurements, in order to eliminate speckle effects. The output of the laser was polarized in the vertical direction, but since the detector was not sensitive to the plane of polarization these effects were not taken into account in the present study.

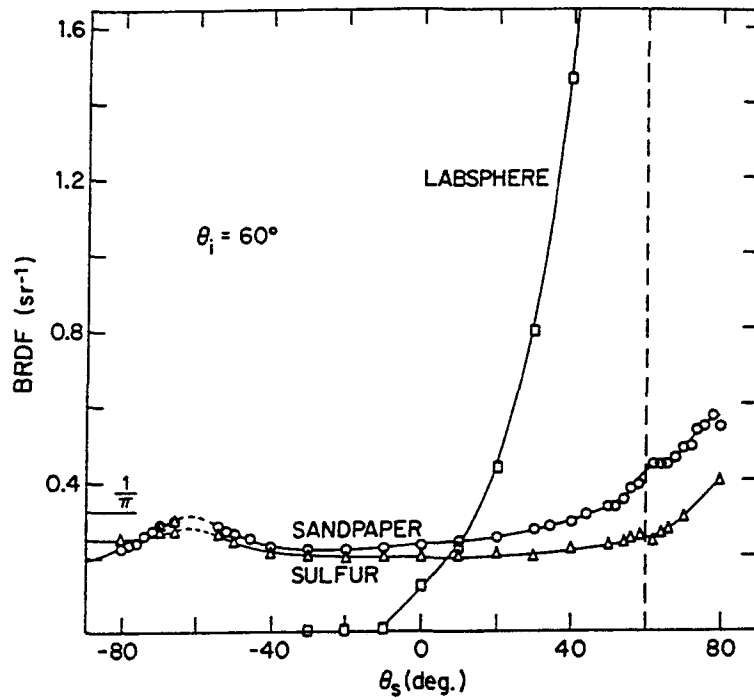


Fig. 4. BRDF of three samples at $10.6 \mu\text{m}$ for an angle of incidence of 60° .

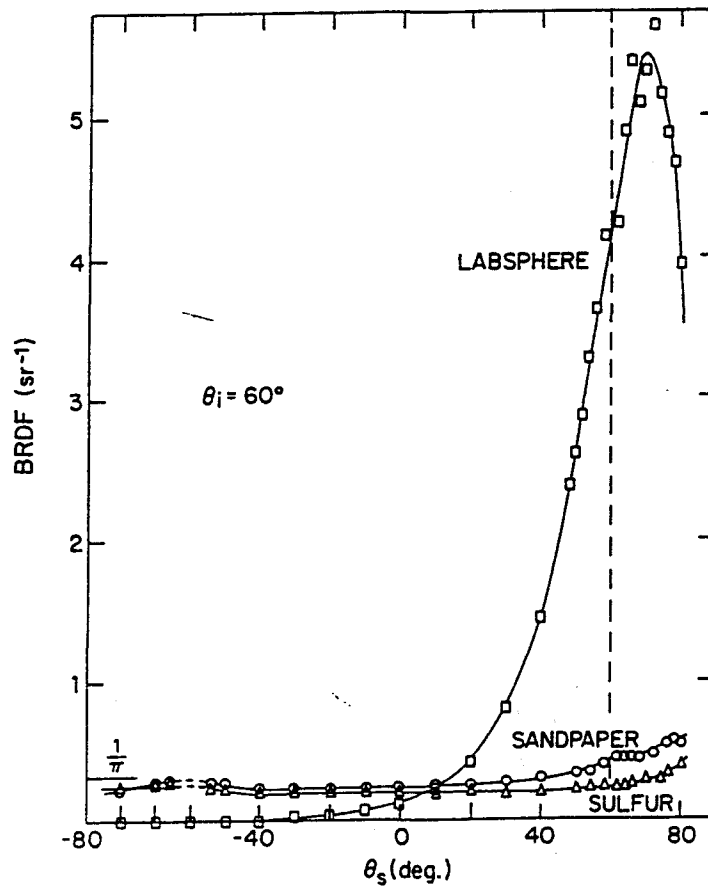


Fig. 5. BRDF of three samples at $10.6 \mu\text{m}$ for an angle of incidence of 60° (note change of scale of ordinate axis).

Figures 1-5 show the results of BRDF for the samples of sulfur, gold-coated sandpaper and Labsphere gold. In each figure the angle of incidence is constant and the scattering angle is shown on the abscissa, with positive angles signifying forward scattering and negative angles indicating backscattering. The angle of incidence is noted as θ_i in each figure. Several results are evident:

1. While sulfur and sandpaper are fairly constant as a function of θ_s , showing a fairly Lambertian behavior, the Labsphere sample has a distinct peak around the specular reflection angle. A closer look at this peak shows that at higher angles of incidence (45 and 60°) the peaks occur at scattering angles which are higher than the specular angle. This is the well-known off-specular effect, which was described and calculated by Torrance and Sparrow.⁽⁷⁾
2. A small peak is observed at the "retro" angle which is equal to the negative of the specular angle and is due to radiation which is backscattered in the direction of the incident beam. This effect was studied recently by Gu *et al.*⁽⁸⁾

Subsequent to these measurements several samples of flame sprayed aluminum (FSA) were obtained, both gold-coated and bare. The results of BRDF measurements for these samples are shown in Fig. 6, where the two types of FSA are compared at angles of incidence of 10 and 45°. It is seen that the bare FSA has a BRDF which is approx. 85% of the BRDF of its gold-coated counterpart. Since the gold-coated FSA was fairly Lambertian and highly reflective its BRDF was more fully measured at angles of incidence of 10, 30, 45 and 60° (see Fig. 7). It is seen that as the angle of incidence increases, this surface deviates from Lambertian behavior. While there is never a specular peak, the BRDF curves at $\theta_i = 45$ and 60° rise at the higher scattering angles (above $\theta_s = 50^\circ$), showing "glint". This is apparently unavoidable in these metal surfaces. At the same time there is also a pronounced "retro" hump in these curves.

A comparison between the results for sandpaper in Figs 1-5 with those on gold-coated FSA in Fig. 7 showed that at small angles of incidence the two surfaces gave BRDF's which were equal to within $\pm 5\%$. However, the performance of the FSA was superior at the higher angles of incidence, when the curves for sandpaper showed higher "glint" at the forward scattering angles than the curves for FSA, while falling below the FSA curves at the extreme negative values of θ_s .

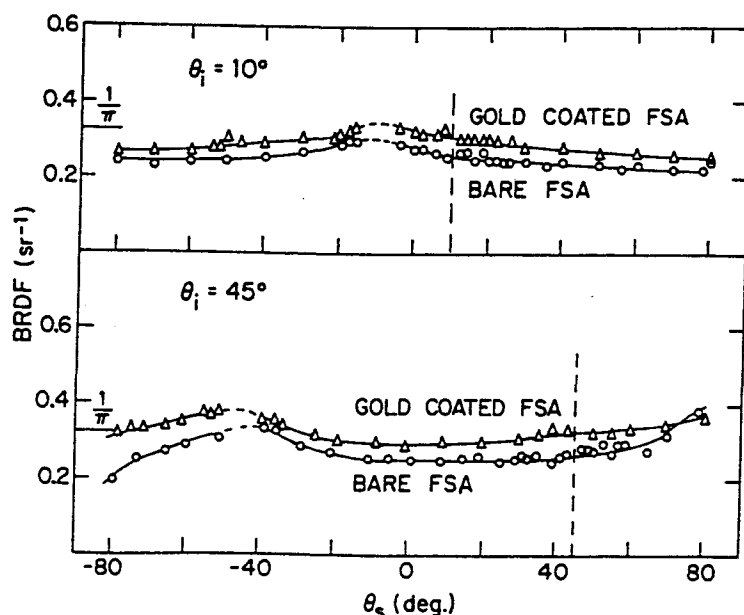


Fig. 6. BRDF of bare and gold-coated flame sprayed aluminum samples at 10 and 45° angles of incidence.

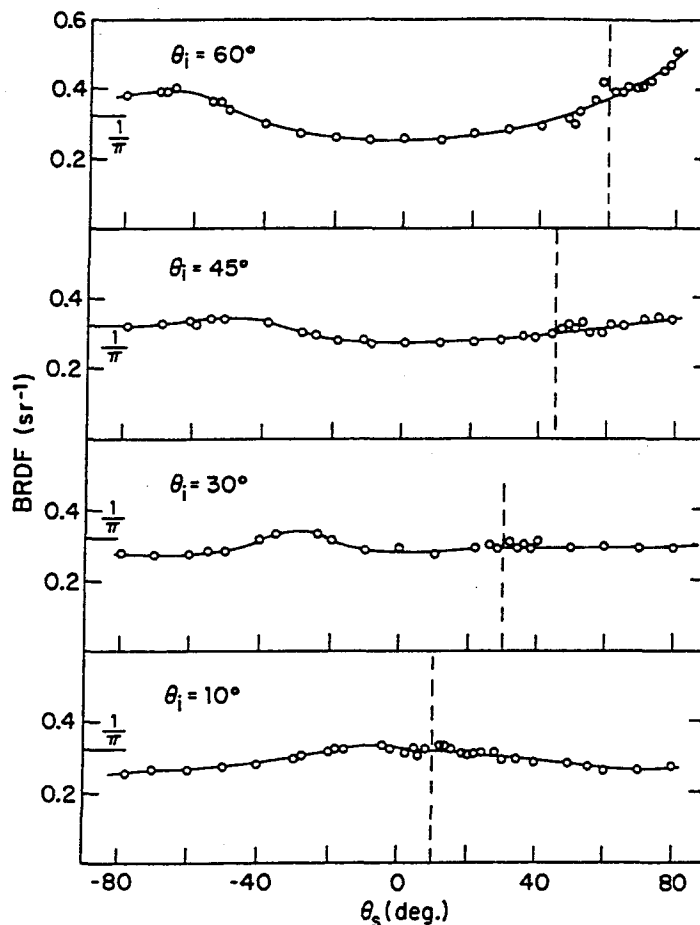


Fig. 7. BRDF of a sample of gold-plated flame sprayed aluminum for angles of incidence of 10, 30, 45 and 60°.

The overall accuracy of the present BRDF results is estimated at $\pm 5\%$, while the relative accuracy from angle to angle is about $\pm 1\%$.

CONCLUSIONS

The present study has shown that sulfur is probably the most Lambertian surface available in the IR, with a BRDF of $0.23 \pm 0.01 \text{ sr}^{-1}$. A more durable and sturdy standard is FSA. It should be pointed out that FSA can be manufactured in large areas of many square feet. In order to increase its reflectance a gold coating may be applied with the desired result of high and flat BRDF curves. These surfaces are believed to be good candidates for use as standard reference materials for diffusely reflecting surfaces in the IR.

Acknowledgements—The authors are grateful to Jim Palmer of the Optical Sciences Center for many fruitful discussions and help in providing several of the samples used in this study. Thanks are also due to Eustace Dereniak for providing the bare FSA samples. The technical help of V. Sinclair was greatly appreciated.

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HUGHES**ELECTRO-OPTICAL & DATA SYSTEMS GROUP**

23 October, 1991

Dr. Robert P. Breault
4601 East First Street
Tuscon, AZ 85711

Dear Bob:

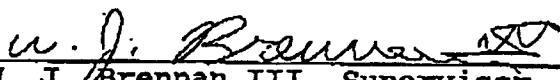
Here is the information you requested from me at the MODIL Scatter Workshop in Bozeman, MT concerning a vendor for Flamesprayed Aluminum. There is actually a Mil Standard (MIL-STD-869) which specifies the fabrication method; I understand that one use of this material is as a non-slip surface. The wire used in the fabrication should conform to MIL-W-6712B, Table II. This material contains some Silicon which may enhance the Lambertian properties at $10.6\mu\text{m}$, however we have not experimented with any other wire alloys. The diameter of the wire can be varied; this controls the macroscopic roughness of the finished surface. We have used #11 gauge wire exclusively; 0.125" diameter is another commonly used size. Our vendor for this material has been:

Plasma Coatings, Inc.
15331 S. Avalon Blvd.
Gardena, CA 90248
(213)-532-3064

Please note that this area code may be changing to 310 in the near future. My contact at Plasma Coatings is a Mr. Dominick Fillipis.

Our reference samples of this material have not varied in Total Hemispheric Reflectance (THR) over periods up to 5 years, within the accuracy of the measuring equipment. Figure 1 shows the BRDF of a typical sample at three wavelengths. Figure 2 compares several different samples, while Figure 3 compares different areas of one sample. In all cases the samples were spinning at approximately 200 RPM to average the effects of laser speckle. Please contact me if you have any questions concerning the data at (213)-616-7867.

Sincerely,


W. J. Brennan III, Supervisor,
Scattering Measurements Laboratory

2000 East El Segundo Boulevard
PO Box 902, El Segundo, CA 90245
(213) 616-1375

Files: BR911010.002, 911018.006, 911022.002. FLAMESPRAYED ALUMINUM #5 AT 3 WAVELENGTHS. 5 AOI.

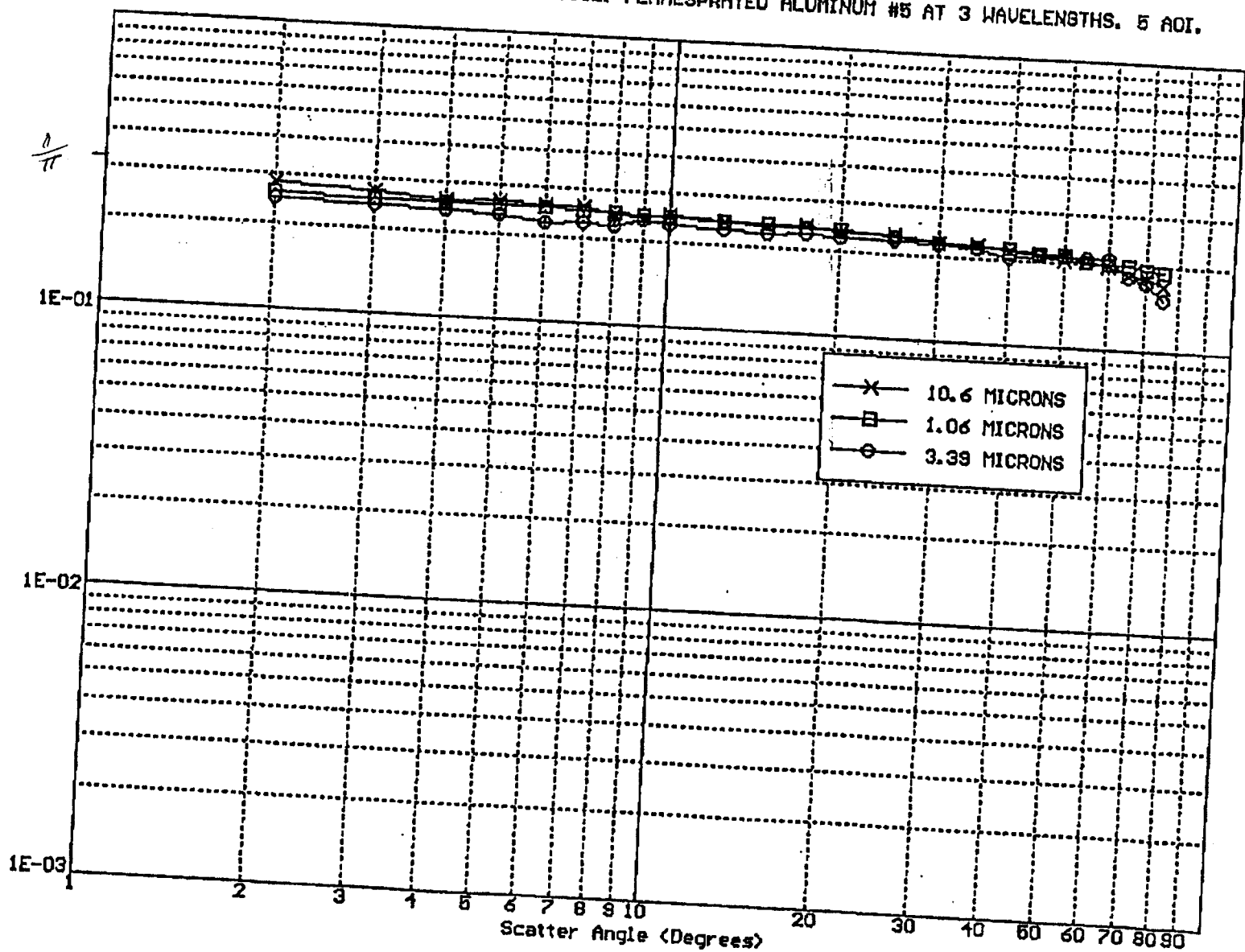


FIGURE 1

BRDF

Scatter Angle (Degrees)

Files: BR900302.009 THRU .015. COMPARISON OF SEVEN FLAMESPRAYED ALUMINUM SAMPLES. 10.6um, 6 AOI

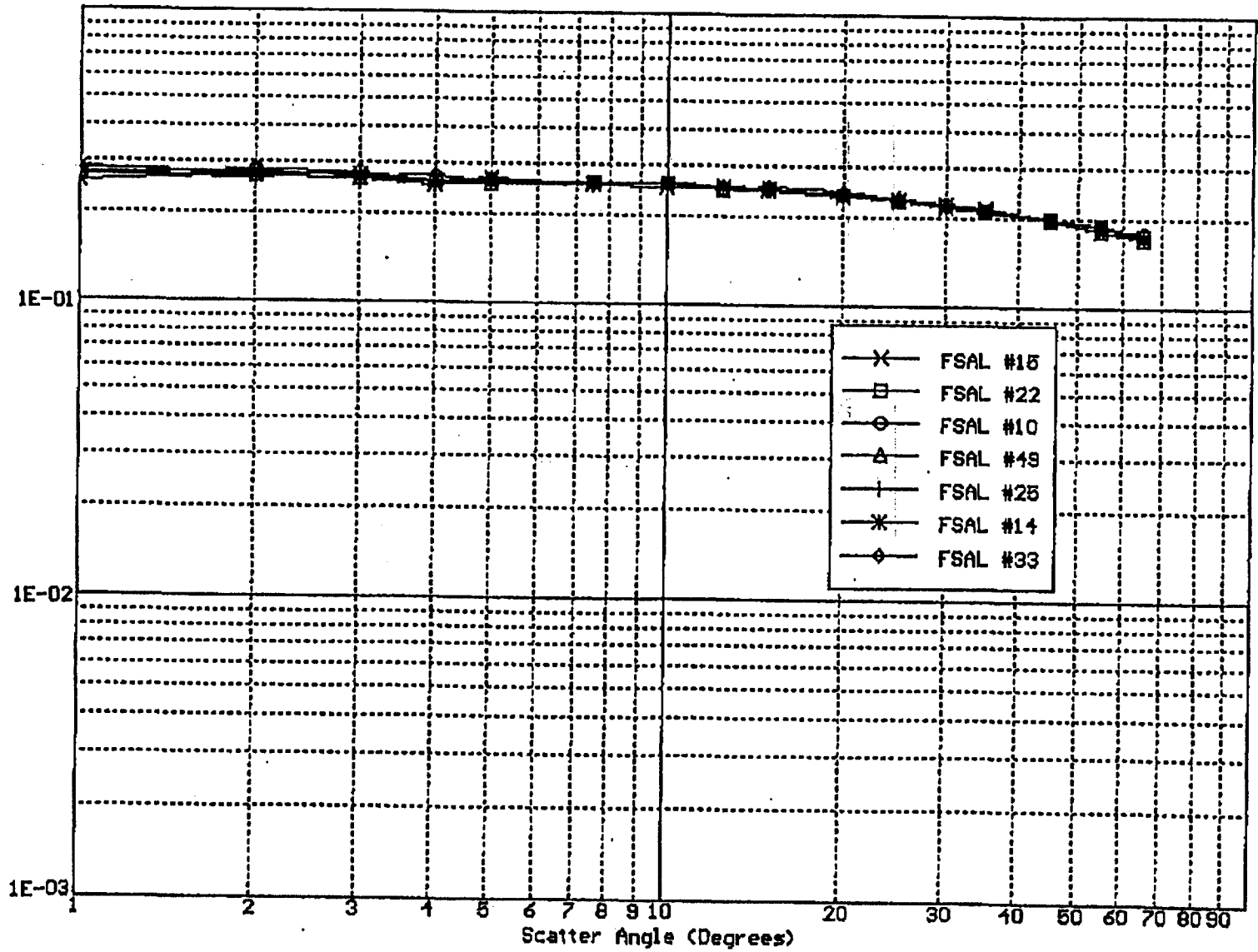


FIGURE 2

BRDF, 1/sr

Scatter Angle (Degrees)

Files: BR911010.011 THRU BR911010.013. FSAL #10 AT 3 LOCATIONS. 10.6um, 8 AOI.

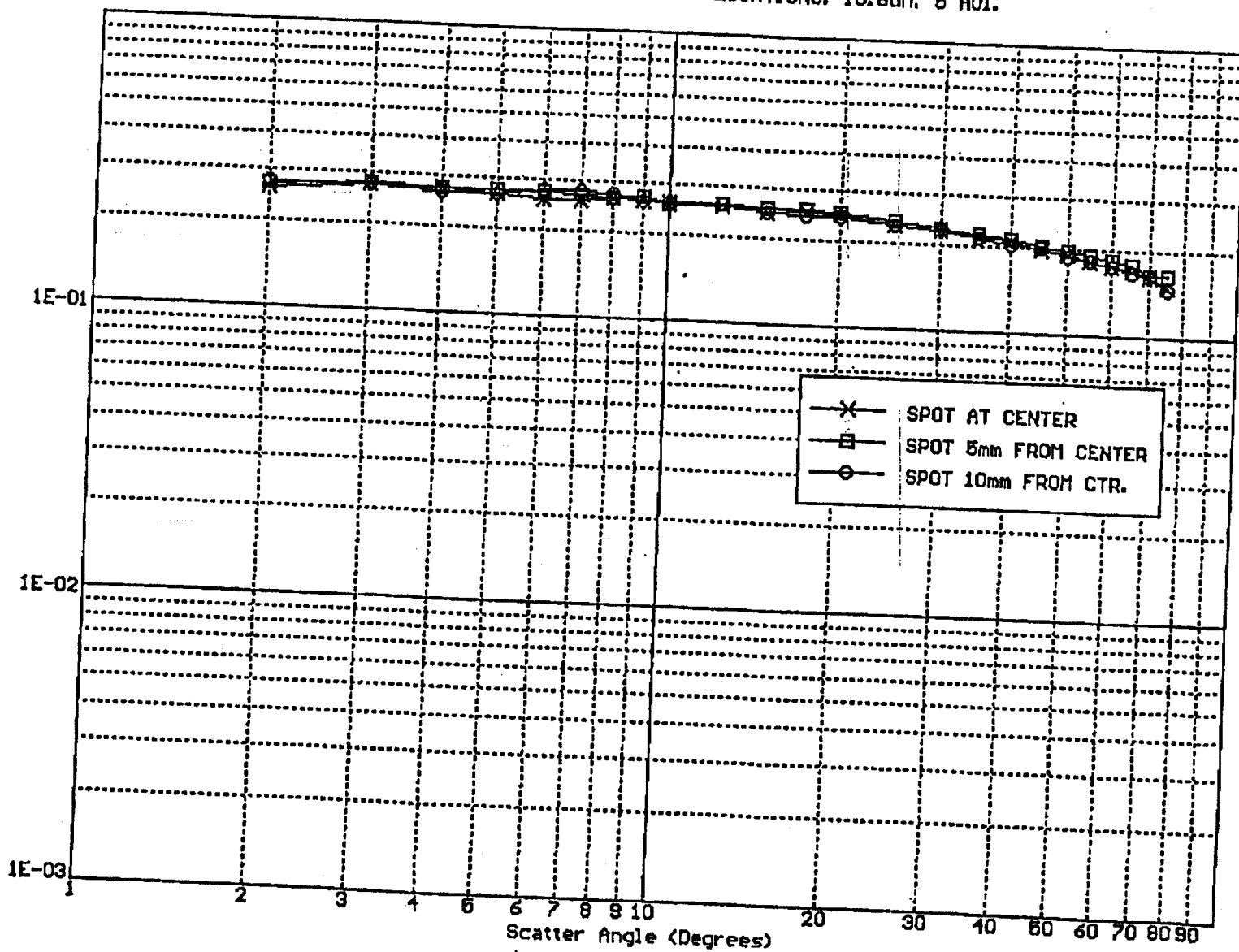


FIGURE 3

BRDF, 1/SF

TOTAL P.05