

Received 10 May 1971

Transverse Vibrations of Double-Tapered Cantilever Beams

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The differential equation is developed from the Bernoulli-Euler equation for the free vibrations of a double-tapered cantilever beam. The beam tapers linearly in the horizontal and in the vertical planes simultaneously. From a computer solution of this equation, a table has been developed from which the fundamental frequency, second, third, fourth, and fifth harmonic can easily be obtained for various taper ratios. Charts are plotted for selected taper ratios in the vertical plane to show the effect of taper ratios on frequency.

INTRODUCTION

This analysis is a continuation of the work^{1,2} started by the authors on the vibration of tapered cantilever beams. This type of beam linearly tapered in either the horizontal or the vertical plane finds wide application for electrical contacts and for springs in electromechanical devices. Occasionally, however, it is advantageous to use a beam which tapers linearly in the horizontal and in the vertical plane simultaneously, and it is necessary to be able to determine easily and accurately the natural frequencies of vibration.

The vibration of tapered cantilever beams has been considered by many investigators; however, only a few have studied the case of the beam tapered linearly in the horizontal and the vertical plane simultaneously. Martin³ in 1956, working on the vibration of turbine blades, developed relations expressing the ratios between frequencies of tapered and untapered beams. Housner and Keightley⁴ in 1962 analyzed a rectangular double-tapered cantilever beam and a conical-shell beam. Lindberg⁵ in 1963 developed a method of deriving a dynamic stiffness matrix for determining the frequency of any nonuniform beam and applied this method in general terms to a linearly double-tapered cantilever beam. Rao⁶ in 1964 developed a method for determining the ratio of the frequency of a tapered beam to the frequency of a uniform beam.

The methods mentioned above are interesting but are in general too complicated for average design use. Because of this, the present investigation was undertaken to provide a quick and accurate method of determining the first five frequencies of double-tapered cantilever beams over a wide range of taper ratios.

The differential equation of motion for a vibrating double-tapered cantilever beam has been derived and, because of its complexity, has been solved on a computer. From this solution, tables have been developed from which the fundamental frequency, second, third, fourth, and fifth harmonic can be obtained for various

taper ratios. Charts have also been plotted to show the effect of taper ratios. In this analysis of transverse vibrations, shear and rotary inertia were neglected.

I. BEAM OF LINEARLY VARIABLE THICKNESS AND OF LINEARLY VARIABLE WIDTH

For the beam shown in a plan and elevated view of Fig. 1, the thickness h at a distance x from the free end is given by

$$h = h_1 + (h_0 - h_1)(x/l) \quad (1)$$

and the width b is given by

$$b = b_1 + (b_0 - b_1)(x/l). \quad (2)$$

The Bernoulli-Euler equation of a vibrating beam is

$$\frac{\partial^2}{\partial x^2} \left(EI \frac{\partial^2 y}{\partial x^2} \right) = - \left(\frac{\rho A}{g} \right) \frac{\partial^2 y}{\partial t^2}, \quad (3)$$

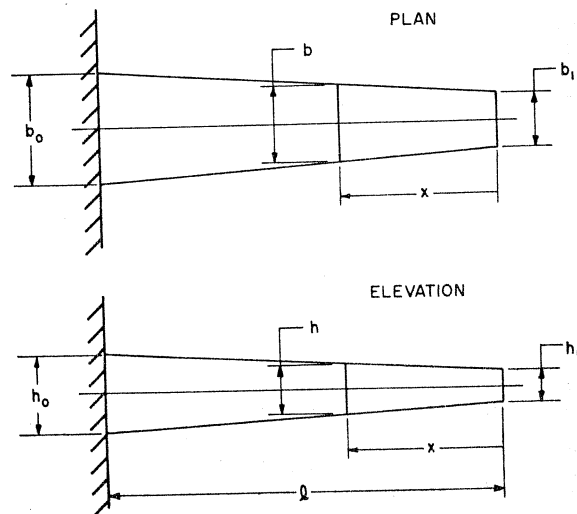


FIG. 1. Cantilever beam tapered linearly in horizontal and in vertical planes simultaneously.

TABLE I. Factor $(lk)^2 \alpha = h_0/h_1$, $\beta = b_0/b_1$.

where $\rho A/g$ is density, A cross-section (constant), E the modulus of inertia. If one has a frequency ω

For the beam

$$I = \frac{1}{12}bh^3 =$$

and

$$A = bh =$$

Substituting $u = x/l$ (so that $u = 0$ for numerical

$$\begin{aligned} & \frac{d^4z}{du^4} + \frac{2d^3z}{du^3} \left\{ \frac{1}{[h_1 + (h_0 - h_1)u]} \right. \\ & \quad + \frac{6d^2z}{du^2} \left\{ \frac{1}{[b_1 + (b_0 - b_1)u]} \right. \\ & \quad \left. \left. + \frac{(h_0 - h_1)}{[h_1 + (h_0 - h_1)u]} \right\} \right\} \end{aligned}$$

To simplify

Equation 7 th

$$\frac{d^4z}{du^4} + \frac{2d^3z}{du^3} \left[\frac{3}{1 - \frac{1}{2} \frac{d^2z}{du^2}} \right] + \frac{6d^2z}{du^2} \left\{ \frac{1}{1 - \frac{1}{2} \frac{d^2z}{du^2}} \right\}$$

Equation 9
tapered beam
to give value
following bou

at x

at x

where $\rho A/g$ is the mass per unit length (ρ weight density, A cross-sectional area, g gravitational constant), E the modulus of elasticity, and I the moment of inertia. If one assumes a sustained free vibration at frequency ω of $y(x,t) = z(x) \sin \omega t$, Eq. 3 becomes

$$\frac{d^2}{dx^2} \left(EI \frac{d^2 z}{dx^2} \right) = \left(\frac{\rho A}{g} \right) \omega^2 z. \quad (4)$$

For the beam tapered in two directions as shown,

$$I = \frac{1}{12} b h^3 = \frac{1}{12} \left[b_1 + (b_0 - b_1) \frac{x}{l} \right] \left[h_1 + (h_0 - h_1) \frac{x}{l} \right]^3 \quad (5)$$

and

$$A = b h = \left[b_1 + (b_0 - b_1) \frac{x}{l} \right] \left[h_1 + (h_0 - h_1) \frac{x}{l} \right]. \quad (6)$$

Substituting Eqs. 5 and 6 into Eq. 4 and letting $u = x/l$ (so that u will always be in the interval $[0,1]$ for numerical integration) gives the following equation:

$$\begin{aligned} \frac{d^4 z}{du^4} + \frac{2d^3 z}{du^3} \left\{ \frac{3(h_0 - h_1)}{[h_1 + (h_0 - h_1)u]} + \frac{(b_0 - b_1)}{[b_1 + (b_0 - b_1)u]} \right\} \\ + \frac{6d^2 z}{du^2} \left\{ \frac{(b_0 - b_1)(h_0 - h_1)}{[b_1 + (b_0 - b_1)u][h_1 + (h_0 - h_1)u]} \right. \\ \left. + \frac{(h_0 - h_1)^2}{[h_1 + (h_0 - h_1)u]^2} \right\} \\ = \frac{12l^4 \rho \omega^2}{Eg} \left[\frac{1}{h_1 + (h_0 - h_1)u} \right]^2 z. \quad (7) \end{aligned}$$

To simplify Eq. 7, let $\alpha = h_0/h_1$, $\beta = b_0/b_1$, and

$$k^4 = 12\rho\omega^2/Egh_1^2. \quad (8)$$

Equation 7 then becomes

$$\begin{aligned} \frac{d^4 z}{du^4} + \frac{2d^3 z}{du^3} \left[\frac{3(\alpha - 1)}{1 + (\alpha - 1)u} + \frac{\beta - 1}{1 + (\beta - 1)u} \right] \\ + \frac{6d^2 z}{du^2} \left\{ \frac{(\beta - 1)(\alpha - 1)}{[1 + (\beta - 1)u][1 + (\alpha - 1)u]} \right. \\ \left. + \frac{(\alpha - 1)^2}{[1 + (\alpha - 1)u]^2} \right\} = \frac{(lk)^4 z}{[1 + (\alpha - 1)u]^2}. \quad (9) \end{aligned}$$

Equation 9 is the equation of motion for a double-tapered beam. It was solved by numerical integration to give values of (lk) for various taper ratios for the following boundary conditions:

at $x=0$ or $u=0$, $d^2 z/du^2 = 0$ and $z=0$,
at $x=l$ or $u=1$, $dz/du = 0$ and $z=0$.

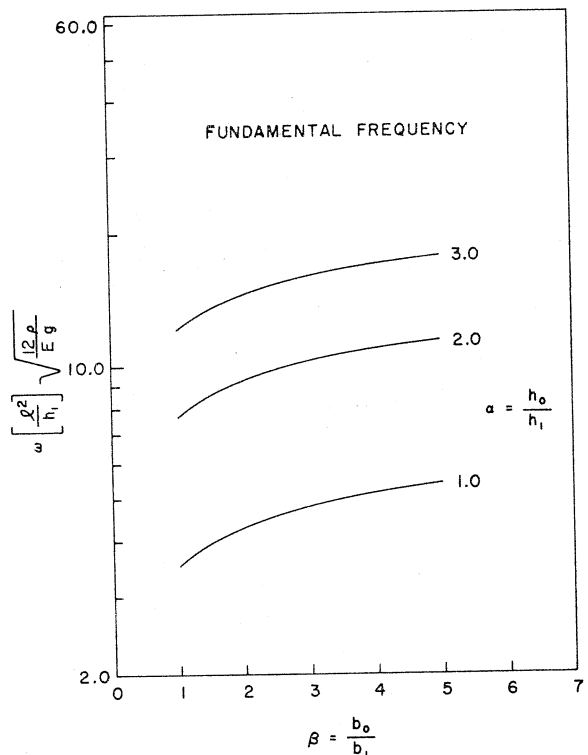


FIG. 2. Effect of taper ratios on fundamental frequency.

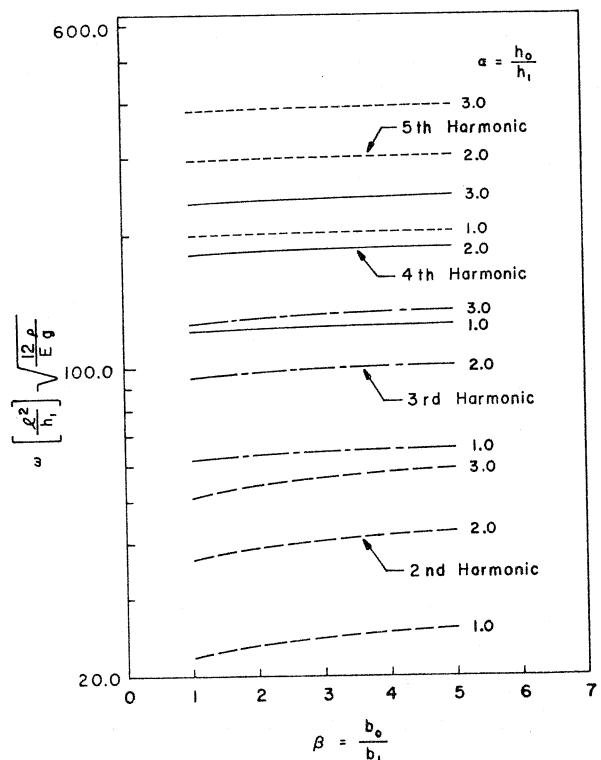


FIG. 3. Effect of taper ratios on second, third, fourth, and fifth harmonic frequencies.

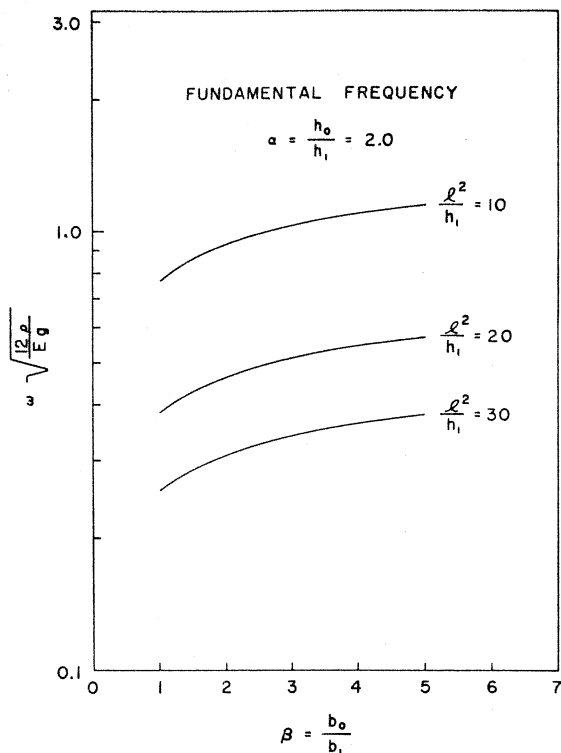


FIG. 4. Effect of factor l^2/h_1 on fundamental frequency.

With these boundary conditions, the solution becomes that of a double-tapered cantilever beam which is truncated and tapers from the fixed end only.

An equation for ω can be written in terms of (lk) from Eq. 8 as

$$\omega = \frac{(lk)^2}{l^2/h_1} \left(\frac{Eg}{12\rho} \right)^{\frac{1}{2}} \quad (10)$$

Table I was developed from the solution of Eq. 9 to give values of $(lk)^2$ for various values of α and β for the fundamental, second, third, fourth, and fifth harmonics. With the values from this table, frequencies can easily be calculated using Eq. 10. It should be mentioned that this table can also be used to calculate the frequencies of a uniform beam where $\alpha = \beta = 1.0$ or for a beam tapered in only one plane.

From Table I, curves were plotted of $\omega(l^2/h_1) \times (12\rho/Eg)^{\frac{1}{2}}$ vs β for $\alpha = 1.0, 2.0$, and 3.0 for the five harmonics as shown in Figs. 2 and 3. Figure 4 shows curves for the fundamental frequency of $\omega(l^2/h_1) \times (12\rho/Eg)^{\frac{1}{2}}$ vs β for $l^2/h_1 = 10, 20$, and 30 with $\alpha = 2.0$.

From Figs. 2 and 3 it can be seen that the taper ratio β in the horizontal plane has much less effect upon frequency than does the ratio α in the vertical plane. Also, the effect of β decreases as the harmonic increases so that at the fourth and fifth harmonics β has practically no effect on frequency. It is interesting to note

from Fig. 3 that there is an overlapping of the curves for the third and fourth harmonics depending upon the value of α . This is also true for the fourth and fifth harmonics.

Figure 4 was plotted to show the effect of the factor l^2/h_1 upon frequency. Several values of l^2/h_1 are given which might be encountered in an actual design. As can be seen, frequencies decrease as this factor increases.

A. Example

Determine the fundamental frequency of vibration of a beryllium copper cantilever beam with the following dimensions:

$$\begin{aligned} h_1 &= 0.010 \text{ in.}, & b_1 &= 0.020 \text{ in.}, \\ h_0 &= 0.020 \text{ in.}, & b_0 &= 0.060 \text{ in.}, \\ l &= 0.500 \text{ in.}, \\ \alpha &= h_0/h_1 = 2, & \beta &= b_0/b_1 = 3. \end{aligned}$$

From Table I, for $\alpha = 2$ and $\beta = 3$, $(lk)^2 = 10.236$. Also, $l^2/h_1 = 0.500^2/0.010 = 25$. For beryllium copper,

$$\begin{aligned} \rho &= 0.297 \text{ lb/in.}^3, \\ E &= 20 \times 10^6 \text{ psi (heat treated).} \end{aligned}$$

From Eq. 10,

$$\begin{aligned} \omega &= \frac{(lk)^2}{l^2/h_1} \left(\frac{Eg}{12\rho} \right)^{\frac{1}{2}} \\ &= \frac{10.236}{25} \left(\frac{20 \times 10^6 \times 386}{12 \times 0.297} \right)^{\frac{1}{2}} \\ &= 19\,100 \text{ rad/sec,} \\ f &= 19\,100/2\pi = 3040 \text{ cps.} \end{aligned}$$

ACKNOWLEDGMENT

The authors express their appreciation to Calvin G. Gray, Jr., for his work in setting up the computer program.

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