LIGO LASER INTERFEROMETER GRAVITATIONAL WAVE OBSERVATORY

LIGO Laboratory / LIGO Scientific Collaboration

LIGO- T040027-03-R *LIGO* 12th April 2008

Conceptual Design of Beamsplitter Suspension for Advanced LIGO

Norna A Robertson and Mark Barton

Distribution of this document: LIGO Science Collaboration

This is an internal working note of the LIGO Project.

California Institute of Technology LIGO Project – MS 18-34 1200 E. California Blvd. Pasadena, CA 91125

Phone (626) 395-2129
Fax (626) 304-9834

E-mail: info@ligo.caltech.edu

P.O. Box 1970
Mail Stop S9-02
Richland WA 99352
Phone 509-372-8106

Fax 509-372-8137

Massachusetts Institute of Technology LIGO Project – NW17-161 175 Albany St Cambridge, MA 02139 Phone (617) 253-4824 Fax (617) 253-7014

E-mail: info@ligo.mit.edu

P.O. Box 940
Livingston, LA 70754
Phone 225-686-3100
Fax 225-686-7189

http://www.ligo.caltech.edu/

1 Introduction

This is Rev-03 of the beamsplitter conceptual design document.

History

Rev-00: 9th February 2004

This version, entitled "Design of Beamsplitter Suspension for Advanced LIGO" presented the case for making the beamsplitter suspension a triple pendulum rather than a quadruple pendulum as used for the ETMs and ITMs. A conceptual design based on the size of beamsplitter at that time (350 mm diameter by 60 mm thick) was presented and curves for predicted seismic isolation performance and thermal noise were given. It was shown that these met the noise requirements for the beamsplitter. The thermal noise curve was produced assuming that the beamsplitter was suspended by four silica fibres of circular cross-section.

Rev-01: 19th November 2007

This version reflected the fact that several factors had changed since the original document was written.

- a) The beamsplitter (BS) size has been increased to 370 mm diameter x 60 mm thick. Currently it is expected to have a wedge angle of 0.9°. This diameter has been shown to have sufficient free aperture to give an acceptable level of optical loss with or without flats on the side— see G070471-00-E for information on losses with flats.
- b) A reassessment of the need for silica fibres has taken place. The baseline is now to use steel wires.
- c) The decision has been taken that the design of the BS and folding mirror (FM) suspensions should be the same.

Rev-02: 16 January 2008

The document has been modified to include transfer functions (from the symmetric MATLAB model) and thermal noise curves (from the Mathematica model) produced using the *same* parameter set (the current set at the time of writing) rather than slightly differing sets which had evolved over the previous few months. The thermal noise curves are presented with more easily read axes. The detailed listing of the Mathematica parameters has been replaced with a reference. The section on choice of parameters has been reduced with details moved into Appendix C. The thermal noise section has been edited.

Rev-03: April 2008

Section 4.2 added – discussion of phi value used for steel wire.

Appendix E added – diagrams and descriptions giving identification of parameters used in the MATLAB model, as listed in appendix A.

Current prototype design rendering has been included (figure 7).

Section 7 added re requirement for a reaction chain. Conclusions section renumbered as 8.

2 Beamsplitter Requirements

Currently the noise requirement at 10 Hz from the sum (suitably added) of BS optics axis motion and vertical motion is 2e-17 m/rt Hz (ref. the cavity optics noise requirements document T010007-02 and ref P Fritschel, e-mail 27 Jan 2004). See Rev-00 of the conceptual design document for fuller discussion. Technical noise sources should be 1/10 of the fundamental noise requirement.

3 Choice of Parameters

The original working design which was investigated was of a triple suspension with approximately equal masses (12.7 kg for the original size of BS) and equal wire lengths of 60 cm at each stage. The choice of equal masses and equal wire lengths as a baseline has come from experience with previous designs and leads to good coupling of modes. In addition using three equal lengths gives the best isolation for a given overall length. For various reasons (available length, change in size of the optic, consequences of changing from silica fibres to steel wires) this original design has been modified. The current parameter list is given in appendix A, and details on the history and reasons for changes are given in Appendix C.

4 Suspension Thermal Noise

4.1 Thermal noise estimate using steel wire and wedged optic

In the 2004 design it was shown that a final stage of the suspension consisting of 4 silica fibres of circular cross-section, 140 micron radius (stress ~500 MPa) and 60 cm length comfortably met the noise requirement (see rev-00 for more details). Silica was chosen as the baseline design. However this decision has since been revisited. There are compelling reasons to use steel wire if it gives acceptable performance: its use gives a significant reduction in complexity of design and construction. One key issue is how much vertical thermal noise is coupled into the longitudinal direction. It is found that with the use of steel wires and a coupling factor of 0.001 from vertical to longitudinal motion, the thermal noise estimate meets the noise requirement at 10 Hz. See figures 1 and 2 below, which are produced assuming an unwedged optic. These estimates have been carried out using Mark Barton's Mathematica model of the beamsplitter, see Appendix B for more details.

The baseline has now been changed to the use of steel wires in conjunction with a coupling factor of \leq 0.001 to be met by the suspension design and the optical layout. From consideration of the optical layout, the practical consequence of requiring the coupling to be 0.001 or less is that the orientation of the wedge on the beamsplitter should be horizontal (rather than vertical) and this is now the baseline configuration for the optic. RODA M070120-02 captures the salient details of the current baseline design: beamsplitter optic size, geometry, wedge orientation and suspension wire material.

Using a horizontal wedge makes very little difference to the overall thermal noise performance expect for the introduction of some coupling of roll into vertical motion, introducing an extra peak at the highest roll mode of 25 Hz. This peak can be removed by moving the spot on the optic from the centre to the sweet spot (the new centre of mass) which for the current parameter set is 2.1 mm from the centre. See figure 3. However even without moving the spot, the roll peak is hardly visible in the sum of longitudinal and vertical with 0.001 coupling, see figure 4.

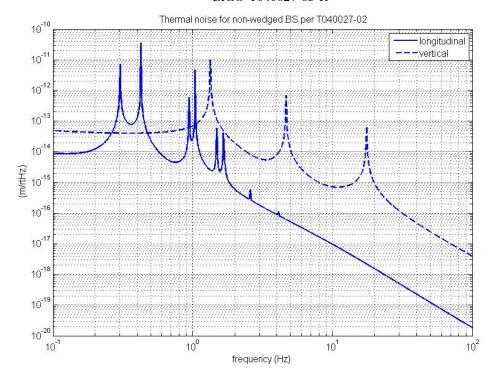


Figure 1. Thermal noise for BS on steel wires, parameters as referenced in Appendix B (unwedged case). Longitudinal and vertical noise estimates are shown separately. Note that at frequencies above the highest vertical mode at 17.5 Hz the vertical thermal noise curve lies more than a factor of 100 above the longitudinal curve.

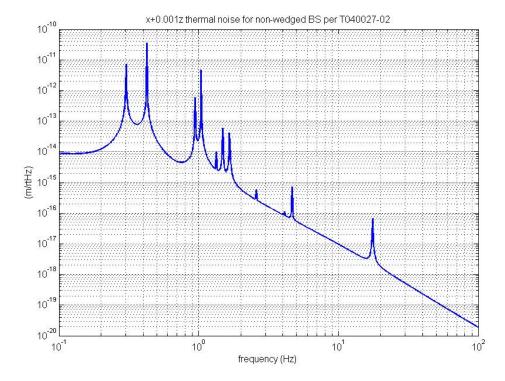


Figure 2. Thermal noise for BS on steel wires (unwedged case) with longitudinal and 0.001 times vertical summed quadratically. The noise is dominated by the longitudinal contribution except at the vertical mode peaks including the highest vertical mode at 17.5 Hz. The level at 10 Hz is 10^-17 m/rt Hz, below the requirement of 2 x 10^-17 m/rt Hz.

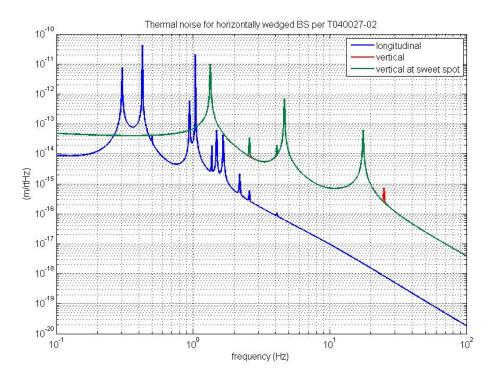


Figure 3. Thermal noise for BS on steel wires, parameters as referenced in Appendix B (horizontal wedged case). Longitudinal and vertical noise estimates are shown separately. The highest roll mode peak can be seen in the vertical plot at 25 Hz. The peak disappears when the beam spot is moved from the centre of the mirror face to the centre of mass (sweet spot),

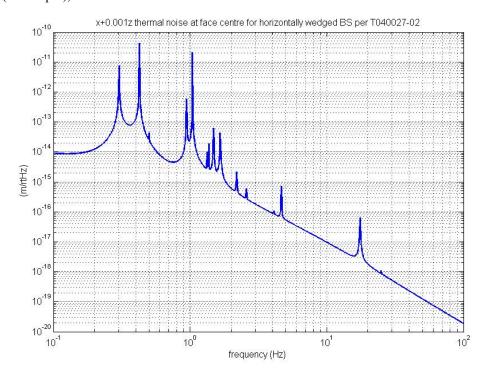


Figure 4. Thermal noise for BS on steel wires (horizontally wedged case) with longitudinal and 0.001 times vertical summed quadratically, for the beam hitting centre of face of optic. The residual roll peak can just be seen at 25 Hz.

4.2 Value of phi for steel wire suspension

The value of the intrinsic loss (phi) in the wire assumed for these curves is $2x10^{-4}$ (ref G Cagnoli et al Phys Lett A 255, p 230, 1999). Recent work by Penn, Harry, Evans, Weiss et al has shown that whereas the intrinsic phi for steel music wire may be even better than this at $\sim 6x10^{-5}$ (see G080108-00-Z), the design of break-off bars at the mirror as used in LIGO1 gives higher loss and variability. Evidence for this comes both from studying the violin mode Qs seen in LIGO 1 and from bench-top experiments. More repeatable results have been obtained using a grooved prism design of break-off, or a clamp. This is an ongoing area of R and D, and any lessons learnt on good design will be taken into account. It is of interest to note that a value of intrinsic phi \sim four times larger than that assumed for the figures above, i.e. a value of 8 x 10^{-4} , would raise the noise level at 10 Hz by a factor of 2, still just meeting the noise requirement. Note that the thermoelastic peak is well above 10 Hz (\sim 460 Hz for the current wire diameter) and so does not significantly affect this argument.

5 Seismic Isolation, Mode Frequencies and Damping

The longitudinal and vertical transfer functions derived from the MATLAB model of the beamsplitter for the parameter set given in appendix A are shown in figures 5 and 6. The mode frequencies are also given in the appendix. The longitudinal transfer function using eddy current damping (with damping time ~10 secs) has a magnitude at 10 Hz of ~1.6 x 10^{-6} . This, combined with active platform noise level of 2 x 10^{-13} m/ $\sqrt{}$ Hz at 10 Hz, gives a noise level at the optic of ~3.2 x 10^{-19} m/ $\sqrt{}$ Hz at 10 Hz. The vertical transfer function is ~8.3 x 10^{-3} at 10 Hz (with damping time ~5 secs), giving vertical noise level at optic of ~1.7 x 10^{-15} m/ $\sqrt{}$ Hz. Including a 10^{-3} coupling factor gives a residual noise level in the horizontal due to vertical motion of 1.7 x 10^{-18} m/ $\sqrt{}$ Hz. Taking the quadratic sum of these numbers yields a total essentially the same as the noise due to vertical alone, and lying well below the requirement.

6 Other Noise Sources

Using the MATLAB model we can also estimate the magnitude of pitch and yaw contributions. The larger of these transfer functions at 10 Hz is for yaw, at \sim 7 x 10⁻⁶. Assuming an angular input at the platform of around 2 x 10⁻¹³ rad/ $\sqrt{\text{Hz}}$ and a 1mm beam offset we find a horizontal noise level of \sim 1.4 x 10⁻²¹ m/ $\sqrt{\text{Hz}}$ at 10 Hz, negligible compared to the requirement.

A further consideration is that of noise introduced by local control. A combination of steep electronic filtering and some eddy current damping (ECD) should yield a workable solution. In fact ECD could comfortably be used without any active control for some modes, and ECD is being incorporated into the design. It has been checked that the thermal noise associated with using ECD is below the noise requirement for the beamsplitter – see Appendix C.

7 Consideration of Requirement for a Reaction Chain

It was originally assumed that the beamsplitter and folding mirror suspensions would require a reaction chain down to the level of the penultimate mass (also called the intermediate mass in a triple pendulum) to allow low-noise feedback. However if the reaction chain is not needed there is obvious saving on design effort. Ken Strain has carried out estimates of the noise introduced by the motion of the actuators assumed attached rigidly to the active platform; see T060157-01-K. The actuator motion is coupled into force noise acting on the intermediate mass and hence into

displacement of the optic. It is shown that using actuators consisting of LIGO1 style coils with double-length magnets (2 mm diam x 6 mm long), which would give 10mN rms force, that there is a safety margin of at least 120. If a larger actuation force is required the Birmingham design of actuator could be used. For 40mN rms force and assuming a larger offset from the sweet spot, the coupling is 4 times smaller than the allowed maximum. Further details can be found in T060157-01-K. In conclusion it appears that a reaction chain is not required and the baseline design does not include one.

8 Conclusions

We have investigated the use of a triple pendulum suspension for the beamsplitter and conclude that it appears to satisfy the noise requirements. The use of steel wires instead of silica fibres has been studied with respect to suspension thermal noise considerations and it is concluded that using steel wires in conjunction with a vertical to longitudinal coupling of 0.001 gives acceptable performance.

The latest parameter set at the time of finishing revision (02) of the document (January 2008) is given in Appendix A. However it should be noted that this is still a conceptual design. Detailed design is currently being carried out and these parameters should only be taken as a guide to the likely final set.

A solidworks rendering of the design of the triple pendulum within its support structure as currently being developed at the Rutherford Appleton Laboratory is shown in figure 7 (courtesy of Joe O'Dell). This depicts an all-metal prototype. The yellow struts are stiffeners to increase the resonant frequencies of the support structure. The magenta piece supports the OSEMs for global alignment and control at the intermediate mass.

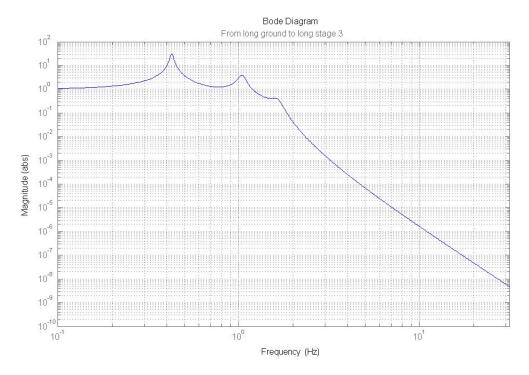


Figure 5. Horizontal (longitudinal) transfer function for beamsplitter triple suspension (with eddy current damping).

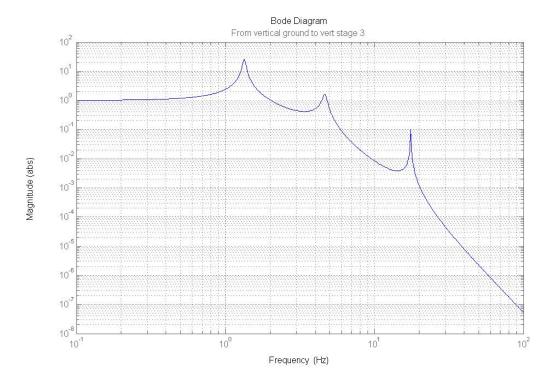


Figure 6. Vertical transfer function for beamsplitter triple suspension (with eddy current damping).

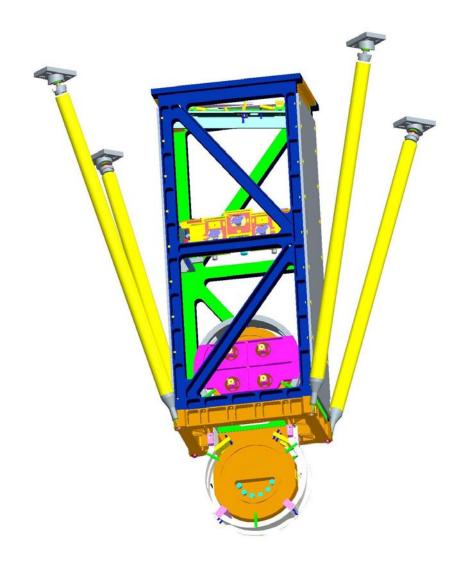


Figure 7. Solidworks rendering of beamsplitter prototype triple suspension.

Appendix A

A.1 Summary of parameters used in the MATLAB code to generate figures 5 and 6.

All numbers are in SI units.

See Appendix E for full explanation of parameter names

```
m1: 1.2627e+001
material1: 'steel'
   I1x: 1.6350e-001
   I1y: 2.4230e-002
   I1z: 1.6190e-001
    m2: 1.3517e+001
    ix: 5.7090e-002
    ir: 1.8500e-001
   I2x: 2.3130e-001
   I2y: 1.1932e-001
   I2z: 1.1932e-001
    m3: 1.3517e+001
material3: 'silica'
    tx: 5.7090e-002
    tr: 1.8500e-001
   I3x: 2.3130e-001
   I3y: 1.1932e-001
   I3z: 1.1932e-001
    11: 6.1200e-001
    12: 6.1000e-001
    13: 5.0000e-001
   nw1:2
   nw2: 4
   nw3: 4
    r1: 3.1500e-004
    r2: 1.8700e-004
    r3: 1.2300e-004
    Y1: 2.1190e+011
    Y2: 2.1190e+011
    Y3: 2.1190e+011
   11b: 2.5000e-001
   a1b: 6.5000e-002
   h1b: 2.4000e-003
   ufc1: 2.8087e+000
   st1: 7.7938e+008
intmode 1: 1.4457e+002
   12b: 1.4000e-001
   a2b: 2.7556e-002
```

```
h2b: 1.6000e-003
             ufc2: 3.2469e+000
              st2: 7.8948e+008
          intmode 2: 3.0733e+002
              su: 0
              si: 1.5000e-002
              sl: 5.0000e-003
              n0: 7.7000e-002
              n1: 1.3000e-001
              n2: 6.0000e-002
              n3: 1.9150e-001
              n4: 1.8650e-001
              n5: 1.8650e-001
            stage2: 1
              d0: -1.8859e-003
              d1: -6.9059e-004
              d2: -6.9059e-004
              d3: -7.1965e-005
              d4: -7.1965e-005
              tl1: 6.0781e-001
              tl2: 5.9428e-001
              tl3: 4.9986e-001
            1 cofm: 1.7019e+000
           1 total: 1.8869e+000
            ribbon: 0
              db: 0
               g: 9.8100e+000
              kc1: 1.9663e+003
              kc2: 2.8127e+003
1 suspoint to centreofoptic: 1.7019e+000
1 suspoint to bottomofoptic: 1.8869e+000
             flex1: 2.8859e-003
             flex2: 1.6906e-003
             flex3: 1.0720e-003
            flex3tr: 1.0720e-003
          longpitch1: [3.0258e-001 4.2554e-001 9.4183e-001]
          longpitch2: [1.0410e+000 1.4855e+000 1.6624e+000]
              yaw: [5.0122e-001 1.3672e+000 2.1890e+000]
          transroll1: [4.2414e-001 1.0394e+000 1.5825e+000]
          transroll2: [2.5848e+000 4.1031e+000 2.4986e+001]
           vertical: [1.3368e+000 4.6494e+000 1.7531e+001]
```

These frequencies can be compared to those in the Mathematica model given in appendix B. The agreement is good to 4 sig. figs.

Notes

- 1) The "d" values shown above are the actual positions of the break-off points to get an "effective" "d" value of 1 mm, taking into account the flexure lengths of the wires. The longitudinal compliance of the blades has not been included in this calculation. FEA models will be required to estimate the compliance and hence the revised positioning of the wire break-off points to compensate for that compliance.
- 2) There has been discussion of increasing the first pitch mode (currently 0.3 Hz) to decrease the overall rms motion of the optic. This can be done by increasing some of the "d" values and correspondingly decreasing wire lengths (to keep the overall height the same). This is ongoing work and is not reflected in the above parameter set.
- 3) The blade parameters and vertical compliance are modeled using the opt.m routine. These should only be taken as representative and not the final design.
- 4) The top mass is represented by the parameters from the Solidworks design for this mass which has been put together by the RAL team. Ref e-mail from Joe O'Dell 17 Oct 2007.
- 5) The penultimate mass and optic are represented in the MATLAB model as being identical in mass, size and moments of inertia. In addition they are represented as symmetric cylinders. In practice
- i) the penultimate mass will be made of metal
- ii) the optics will be horizontally wedged as described in Appendix C.1 below. It is proposed that the penultimate mass is similarly wedged, and hung such that the wedge is oppositely oriented so that the overall loading on the blades above is symmetric.

Appendix B

Mark Barton's Mathematica models used to generate the thermal noise curves can be found at http://www.ligo.caltech.edu/%7ee2e/SUSmodels/

under the sidebar –follow the link to the Triple Xtra-Lite model page.

For the unwedged case the resulting mode frequencies in Hz are

"N	f type		
1	0.3025809648892492 pitch	3 pitch2	
2	0.424137285679903 y3	y2	
3	0.42553476024458786	x 3	x2
4	0.5012213586831055 yaw3	yaw2	
5	0.9417991878837583 pitch	2 pitch3	
6	1.0393741072077616 y2	y3	
7	1.0409607596847752 x2	x 3	
8	1.3367465497338584 z3	z2	

```
9
      1.3672350227875962 yaw1
10
      1.4855813775299265 pitch1
      1.5825732377865354 v1
11
                                 y2
12
      1.6624534003792206 x1
                                 x2
13
      2.188980912315431 yaw2
                                 yaw3
14
      2.5848004448313247 roll1
                                 roll3
15
      4.1031558811539375 roll1
                                 roll3
16
      4.649514275151824 z1
17
      17.530495104930278 z2
                                 z3
18
      24.985132724495962 roll2
                                 roll3"
```

For the horizontally wedged beamsplitter with optic parameters as described in Appendix C.1 the resulting mode frequencies in Hz are

```
N
      f
                            type
                           pitch3 pitch2
1
    0.3025713527967201
2
    0.4241368839053088
                           y3
                                y2
3
    0.4255231004518705
                           x3
                                x2
4
    0.5010788376001445
                           yaw3 yaw2
5
    0.9417746185630207
                           pitch2 pitch3
6
    1.0393734894344189
                           y2
                                y3
7
     1.041104703254956
                          x^2
                                x3
8
     1.3367460041342083
                           z3
                                z2
     1.367205502017105
                          yaw1
10
     1.4855740334487244
                            pitch1
                           y1
11
     1.5825364214780144
                                 y2
12
     1.6623505850841167
                            x1
                                 x2
13
     2.189237845349943
                           yaw2 yaw3
14
     2.5846090953208374
                           roll1 roll3
15
     4.102872628595596
                           roll1 roll3
16
     4.649513116237302
                           z1
17
     17.533422108920426
                            z2
                                 z3
18
     24.985170357403824
                            roll2 roll3
```

The wedged case assumes a symmetric horizontal wedge on the BS, an opposite wedge on the penultimate mass, adjustments to the wires such that everything hangs straight, rotation of the structure by half the wedge angle so that the HR face is realigned with +x, and in the case of the noise at the sweet spot, alignment of the beam on the sweet spot in front of the COM.

Note that the coordinates ("type") in the listing come from a crude mode ID function that ranks the coefficients in the eigenvector in descending order and prints coefficient names until half the total squared amplitude in the mode has been accounted for.

Appendix C: History of modifications to parameter set.

C.1 Design of beamsplitter mass

The details of diameter, thickness and wedge for the beamsplitter have evolved since the original conceptual design document was written. At the time of finalising Rev-01 (19th November 2007) RODA M070120-02 has been produced giving the design as follows: 370 mm diameter, horizontal symmetrical wedge with full wedge angle 0.9 degrees, thick end of wedge 60 mm thick, giving a mass of 13.5 kg. The mass is represented in the MATLAB model by assuming a thickness of the beamsplitter which is the average of the thin end and thick end of the wedge (Note that the MATLAB model assumes symmetry in the mass shapes). The penultimate mass has been modeled to be identical in mass and size, but will be made of metal with suitable holes to give the correct mass

C.2 Violin Mode Frequencies and Length of Wires

The SUS group was asked by Peter F to consider shortening the length of the final stage of the suspension so that its violin mode frequency is higher than what would be obtained with the 600 mm length originally proposed. By shortening to 500 mm and allowing a stress level of \sim 710 MPa (slightly more than the working value assumed for other Adv LIGO wire suspensions of 670 MPa) the frequency is raised from \sim 240 Hz to 300 Hz. Note that the use of steel rather than silica has reduced the expected violin mode frequency due to steel's higher density.

C.3 Overall Length of Suspension

The original overall length was chosen to satisfy the available length for a beamsplitter suspension in a BSC (noting that this was at that time expected to be 70 mm longer than for an ETM) *prior* to considerations to reduce the overall length of BSC suspension structures as summarized in T040028-00. Since then the recommendations on length in T040028 have been adopted, and the decision to make the FM the same design as the BS has been taken. Since the FM must necessarily be very close to the same length as an ITM (they are adjacent to each other and the laser beam is close to horizontal), this implies that for a common BS/FM design, the choice for the length of the BS or FM is now such that the BS, FM and ITM mirror centres are the same distance from the optics table. Note that this doesnt imply that the suspension lengths will necessarily be the same. The distance between the top suspension point and the optics table above need not be the same.

Ian W at RAL has indicated that a longer pendulum length for the beamsplitter or folding mirror could be incorporated within the same overall structure length by changing the way the top blade assembly is fixed within the structure compared to how this is done in the quad. The overall length of the pendulum could be increased by 66mm. Since this in principal gives a little more isolation, it has been used in the latest parameter set. The details on length are as follows

As per the following document, the optic table to optic CL (CL = centre line) for the ETM quad suspension is 1742 mm

http://www.eng-external.rl.ac.uk/advligo/Reviews/PDR3/documents/overview/t060142-00-k.pdf.

For the quad the length from tip of top blade to centre of optics is 1636 mm. Thus this allows 1742-1636 = 106 mm as space to fit in the blade supports and mount to the table in the quad. For the beamsplitter Ian is proposing that we can mount the blade tips closer to the table by 66 mm, so that they are now only 40 mm from the table. This means that we can make the overall length of the splitter from blade tip to centre of optic be 1636 + 66 = 1702 mm.

Appendix D: Use of Eddy Current Damping

In the current detailed design for the top mass ECD units similar to those being used in the ETM/ITM noise prototype are being incorporated. These units are arranged in clusters of 4 magnets (nominally 10 mm diam x 10 mm thick) with 4 such clusters acting in each of the longitudinal and vertical directions, arranged so that they also provide pitch, roll and yaw damping. Four clusters of four such magnets will give a damping constant of $b \sim 27$ kg/s (ref P060013-00-R). The transfer functions shown in figures 5 and 6 assume this value of b in longitudinal and vertical directions. The decay time to 1/e = 10.1 secs for longitudinal and 4.7 secs for vertical (we may choose to reduce magnet strength for this direction).

We can estimate the thermal noise due to this damping and show that it is acceptable. The noise force at the top mass where the damping is applied is given by $F^2 = 4k$ Tb, where k = Boltzmann's constant and T = temperature(K).

For b = 27 kg/s, $F = 6.7 \times 10^{-10} \text{ N/rt Hz}$.

From the MATLAB model we find the following:

a) Longitudinal TF at 10 Hz for force at top mass to displacement of mirror, TF(long) = 8.7×10^{-10} m/N.

Hence longitudinal motion due to thermal noise = $F \times TF(long) = 5.8 \times 10^{-19} \text{ m/rtHz}$.

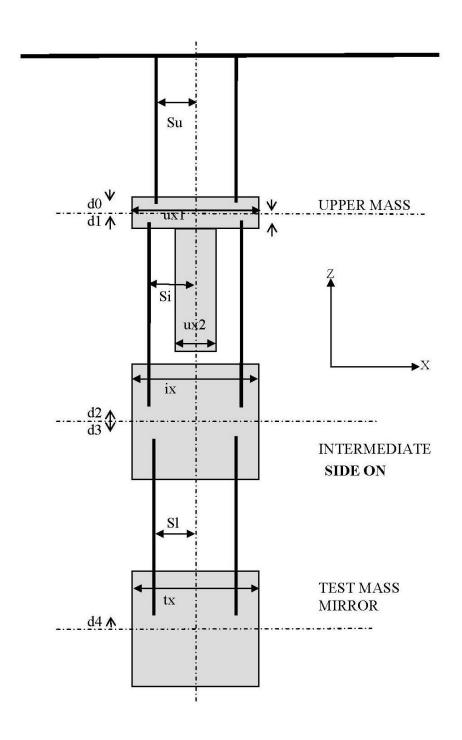
b) Vertical TF at 10 Hz for force at top mass to displacement of mirror, TF(vert) = $1.9 \times 10^{\circ}-6 \text{ m/N}$.

Hence vertical motion due to thermal noise = $F \times TF(vert) = 1.3 \times 10^{-15} \text{ m/rtHz}$. Assuming coupling of 0.1%, this gives longitudinal motion of 1.2 x 10^-18 m/rt Hz.

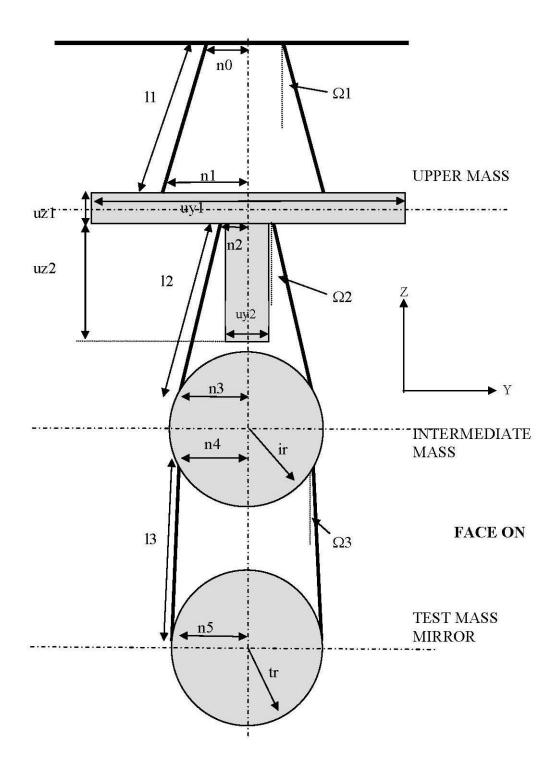
These values of longitudinal motion should be compared to the *technical* noise requirement for the beamsplitter of 2×10^{-18} m/rtHz at 10 Hz. They are both below that value.

Appendix E

E.1 The parameters of a triple pendulum (side on view)



E.2 The parameters for a triple pendulum (face on view)



E.3 Other parameters listed in appendix A.

m1, m2, m3: masses from top to bottom

Iix, Iiy, Iiz where i = 1,2,3 from top to bottom mass = moments of inertia as follows

Iix: moment of inertia (transverse roll)

Iiy: moment of inertia (longitudinal pitch)

Iiz: moment of inertia (yaw)

nwi = number of suspension wires at each stage from top to bottom

ri = wire radius from top to bottom

Yi = Young's modulus of wire/fibre from top to bottom

11b, a1b, h1b: length, width at root, thickness of top blades

ufc1, st1, intmode_1: uncoupled frequency of top blade with mass immediately below it, stress in blade and estiamted first internal mode frequency (all data returned from opt.m m-file routine)

12b etc – same as above for lower blades

stage 2 = 1

If pend.stage2 is defined and non-zero, d0-d4 are interpreted as raw values, i.e., as actual wire breakoff vertical positions

tl1, tl2, tl3: centre to centre vertical separations at each stage - from top suspension point to centre of top mass, centre of top mass to centre of intermediate mass, and centre of intermediate mass to centre of beamsplitter optic respectively

ribbon = 0: round wires/fibres are used (i.e not ribbons)

db = 0: no natural damping included

g: accel. due to gravity

kc1, kc2: blade stiffness (top and bottom respectively)

1 suspoint to centreofoptic: length from top suspension point to centre of topic = tl1+tl2+tl3

1 suspoint to bottomofoptic: length from top suspension point to bottom of optic

flex1, flex2, flex3: flexure length for wire (top to bottom respectively)

flextr – flexure length for ribbon in transverse/roll direction (same as flex3 if round fibre used)